

From pion formfactor to exclusive Drell-Yan processes

XXII Baldin ISHEPP, *Session in Honor of the 80th Anniversary of Professor Anatoly Efremov*

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Outline

- AV and pion formfactor
- From pion FF to Exclusive DY
- GPDxDA vs GPDxGPD
- Factorization vs dispersion relations
- Similarity to twist-3 semiinclusive SSA
- Gluons GPDs case and quarkonia/Higgs CEP
- Conclusions

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V. Efremov, A.V. Radyushkin

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Fig. 2

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ОБЪЕДИНЕННЫЙ
ИНСТИТУТ
ЯДЕРНЫХ
ИССЛЕДОВАНИЙ
ДУБНА



E2 - 9717

A.V.Efremov, A.V.Radyushkin

ASYMPTOTICS OF PION ELECTROMAGNETIC
FORM FACTOR IN SCALE INVARIANT
QUARK MODEL

1976

Factorization and Asymptotical Behavior of Pion Form-Factor in QCD (In Russian)

A.V. Efremov, A.V. Radyushkin (Dubna, JINR). Nov 1979. 12 pp.

Published in **Phys.Lett. B94 (1980) 245-250**

JINR-P2-12900

DOI: [10.1016/0370-2693\(80\)90869-2](https://doi.org/10.1016/0370-2693(80)90869-2)

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■ QCD

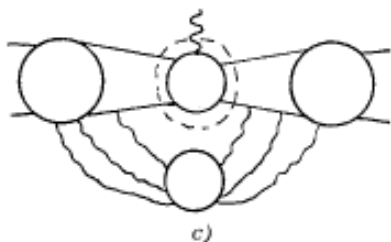
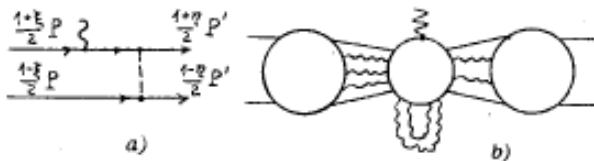


Fig. 5

- Note symmetric kinematics – most convenient for GPDs



E2 - 11983

A.V.Efremov, A.V.Radyushkin

ASYMPTOTICAL BEHAVIOUR
OF PION ELECTROMAGNETIC FORM FACTOR
IN QCD

1978

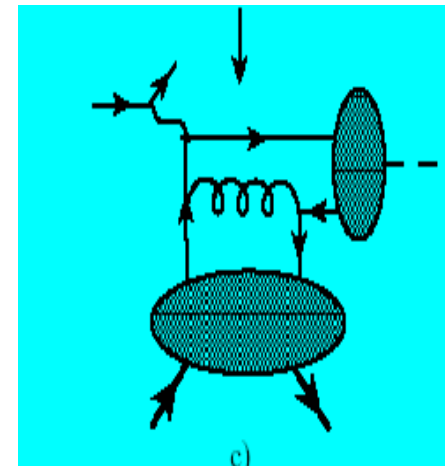
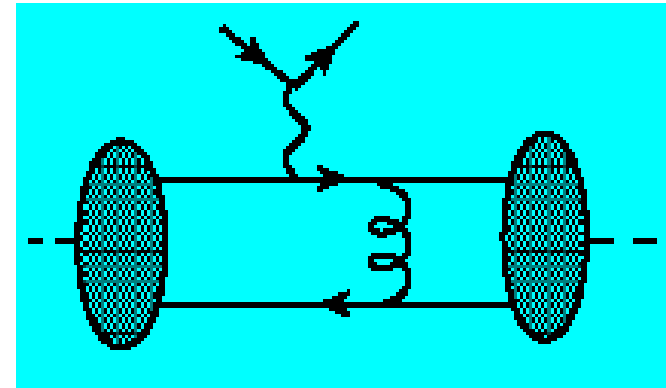
Way to (semi) exclusive DY

- Starting from (Pion) form factor- 2 DA's

$$F \propto \left(\int dx \frac{\phi(x)}{1-x} \right)^2$$

- 1 DA -> GPD : Exclusive mesons production (Frankfurt, Strikman) / DY ("classical" mechanism – Pire, Schymanowski, Wagner)

$$M \propto \int dx \frac{\phi(x)}{1-x} \int dx \frac{H(x, \xi)}{x - \xi + i\epsilon}$$



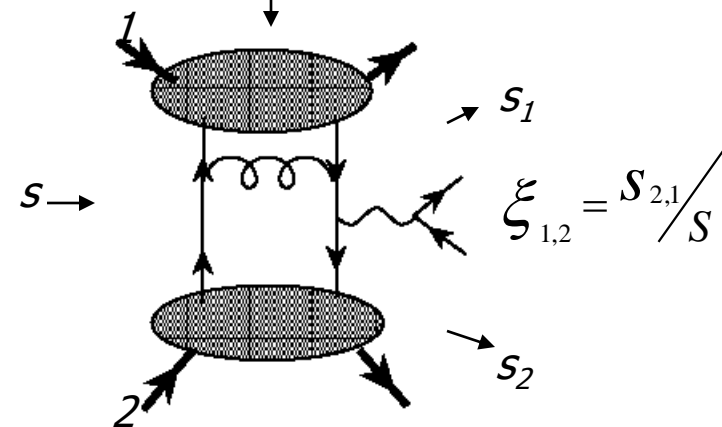
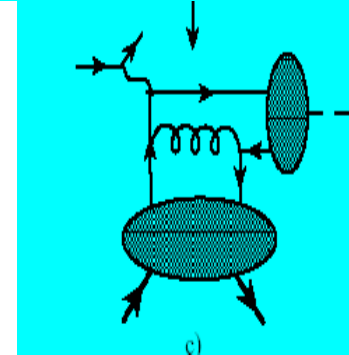
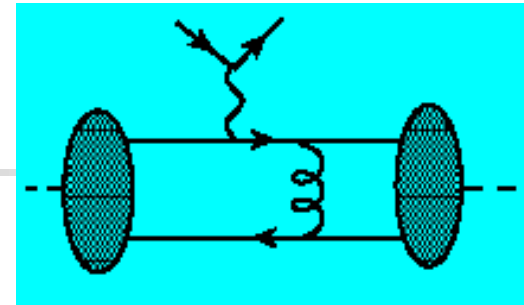
Next step: 2 DA's -> 2 GPD's

- Exclusive double diffractive DY process
- Analytic continuation:

$$M \propto \int dx \frac{H(x, \xi_1)}{x - \xi_1 \pm i\epsilon} \int dy \frac{H(y, \xi_2)}{y - \xi_2 \mp i\epsilon}$$

- DIFFERS from direct calculation – NO factorization in physical region

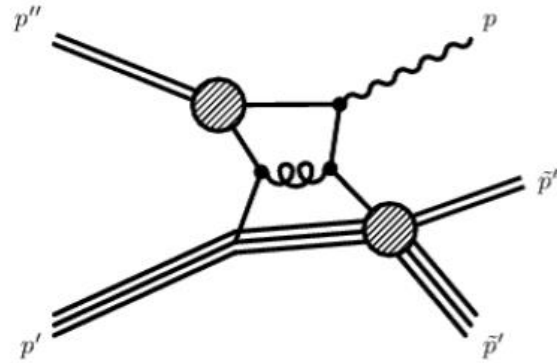
$$M \propto \iint dx dy \frac{H(x, \xi_1) H(y, \xi_2)}{(x - \xi_1)(y - \xi_2) + i\epsilon}$$



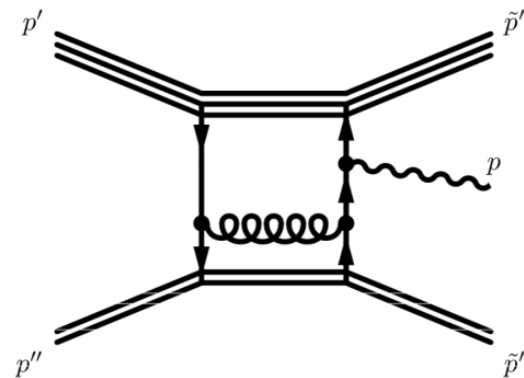
Comparison of different "handbag" contributions for pion-nucleon diffractive DY (AP,OT, in progress)

- Exclusive process $p\pi^0 \rightarrow p\pi^0\gamma^*$:

- Classical



- Double diffractive





GPDxGPD

- Probability amplitude:

$$S_{p\pi 1} = i \frac{(2\pi)^4}{2N_c^2} \frac{eg_{(s)}^2}{(2V)^{\frac{5}{2}} \sqrt{\epsilon' \epsilon'' \tilde{\epsilon}' \tilde{\epsilon}'' \epsilon}} \cdot e_\lambda^* \cdot \frac{1}{(\bar{P}', \bar{P}'')} \cdot \left(\frac{\bar{P}'^\lambda}{\xi_2} - \frac{\bar{P}''^\lambda}{\xi_1} \right) \cdot I_1 \cdot \\ \cdot \delta(p' + p'' - \tilde{p}' - \tilde{p}'' - q)$$

$$\xi_1, \xi_2 \in [0, 1] \quad \bar{P}' = \frac{p' + \tilde{p}'}{2} \quad \bar{P}'' = \frac{p'' + \tilde{p}''}{2}$$

- Longitudinal polarisation



Compton FFs

$$I_1 = \int_{-1}^1 \frac{e_u H_p^{(u)}(x) H_\pi^{(u)}(y)}{(x + \xi_1)(y - \xi_2) - i\varepsilon_g} dx dy - \int_{-1}^1 \frac{e_u H_p^{(u)}(x) H_\pi^{(u)}(y)}{(x - \xi_1)(y + \xi_2) - i\varepsilon_g} dx dy +$$
$$+ \int_{-1}^1 \frac{e_d H_p^{(d)}(x) H_\pi^{(d)}(y)}{(x + \xi_1)(y - \xi_2) - i\varepsilon_g} dx dy - \int_{-1}^1 \frac{e_d H_p^{(d)}(x) H_\pi^{(d)}(y)}{(x - \xi_1)(y + \xi_2) - i\varepsilon_g} dx dy$$

$$H_p^{(q)}(x) = \begin{cases} H_p^q(x) & x > 0 \\ -H_p^{\bar{q}}(-x) & x < 0 \end{cases}$$



Classical mechanism

- Probability amplitude:

$$S_{p\pi 2} = i \frac{(2\pi)^4}{2N_c^2} \frac{ee_q g^2(s)}{(2V)^{\frac{5}{2}} \sqrt{\varepsilon' \varepsilon'' \tilde{\varepsilon}' \tilde{\varepsilon}'' \varepsilon}} \cdot e_\lambda^* \cdot \frac{1}{(\bar{P}, p'')} \cdot \left(2\bar{P}^\lambda - \frac{p''^\lambda}{\xi} \right) \cdot I_2 \cdot \delta(p' + p'' - \tilde{p}' - \tilde{p}'' - q)$$

$$\xi \in [0, 1] \quad \bar{P} = \frac{p' + \tilde{p}' + \tilde{p}''}{2}$$

- Also longitudinal polarisation



Compton FF

$$I_2 = \int_{-1}^1 dx \int_0^1 \frac{H_{p\pi}(x)\Phi(y)}{(x+\xi)y+i\varepsilon_g} dy + \int_{-1}^1 dx \int_0^1 \frac{H_{p\pi}(x)\Phi(y)}{(x-\xi)y-i\varepsilon_g} dy$$

$$H_{p\pi}(x) = \frac{1}{\sqrt{2}} \left[e_u H_{p\pi}^{(u)}(x) - e_d H_{p\pi}^{(d)}(x) \right]$$

$\Phi(y)$ – distribution amplitude of pion

Pion decay constant cancel in DA and

TDA (natural normalization in

accordance with soft-pion theorems)



Comparison of the processes

- Ratio of the cross-sections (prefactor due to EM GI):

$$\frac{A_1}{A_2} = \frac{2(1 + \xi_1)(1 + \xi_2)}{2(1 + \xi)\xi_1\xi_2} \cdot \frac{|I_1|^2}{|I_2|^2} = \frac{s(s_1 + s_2 - s)}{s_1 s_2} \cdot \frac{|I_1|^2}{|I_2|^2}$$

$$s = (p' + p'')^2 \quad s_1 = (\tilde{p}' + q)^2 \quad s_2 = (\tilde{p}'' + q)^2$$

- GPDxGPD – dominant in forward region



Interference

- Cross section:

$$\frac{d\sigma}{d^3\tilde{p}'d^3\tilde{p}''} = \frac{\alpha_{(em)}\alpha_{(s)}^2}{2^6 N_c^4 \pi^2 \tilde{\epsilon}'\tilde{\epsilon}''} \cdot \frac{1}{(p', p'')^2} \cdot \left[\frac{8s^2}{s_1 s_2} |I_1|^2 + \frac{8s}{s_1 + s_2 - s} |I_2|^2 + 4s \left(\frac{1}{s_2} + \frac{s}{s_1(s_1 + s_2 - s)} \right) (I_1^* I_2 + I_2^* I_1) \right] \cdot \delta(q^2 - m_\gamma^2)$$

- Interference term – nonsymmetric wrt $s_{1,2}$ interchange



Compton FFs: analytic continuation and IR regularization

- Pole prescription unclear – cuts in s and $s_{1,2}$ produce different signs
- Similar to pion dissociation to dijet (D. Ivanov et al)
- $s_{1,2}/s$ the same when both positive or both negative – cancellation of cuts
- Similar to cancellation of cuts in s and Q^2 for semi-inclusive annihilation
- Suggestion for time-like DVCS: Q^2/s – cancel; Q^2/u due to Q^2 cut: different duality properties with – resonances in vector rather than s - channel (BG-type ones)
- Direct calculation: IR regularization (of imaginary part) like in GK model

Similarity to twist 3 Single Spin Asymmetries

- Another direction of AV activity at about the same time

AV

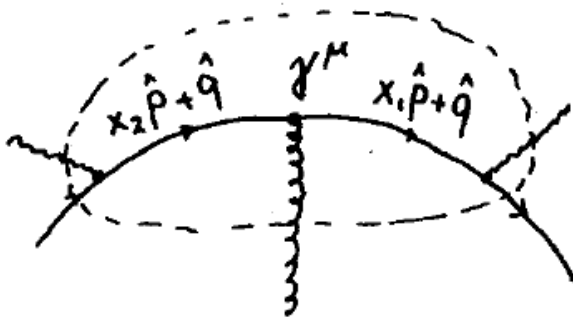


Fig. 3

QCD Asymmetry and Polarized Hadron Structure Functions (In Russian)

A.V. Efremov, O.V. Teryaev (Dubna, JINR). Aug

1984. 6 pp.

Published in **Phys.Lett. B150 (1985) 383**

JINR-P2-84-603

DOI: [10.1016/0370-2693\(85\)90999-2](https://doi.org/10.1016/0370-2693(85)90999-2)

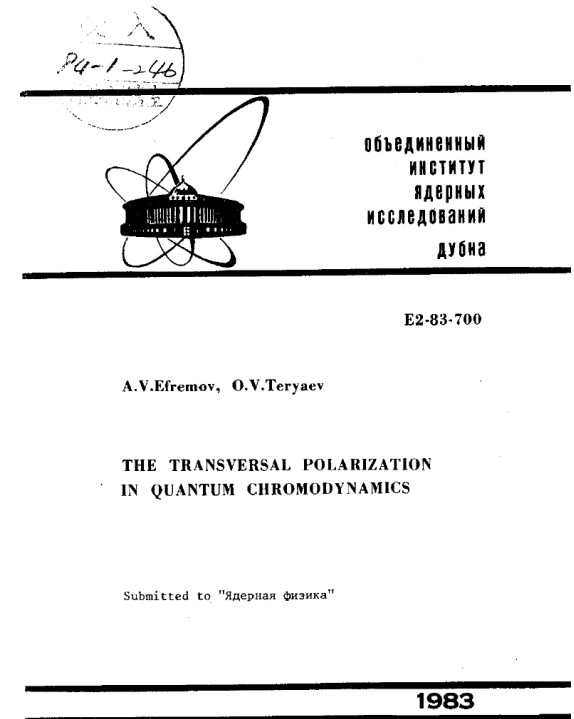
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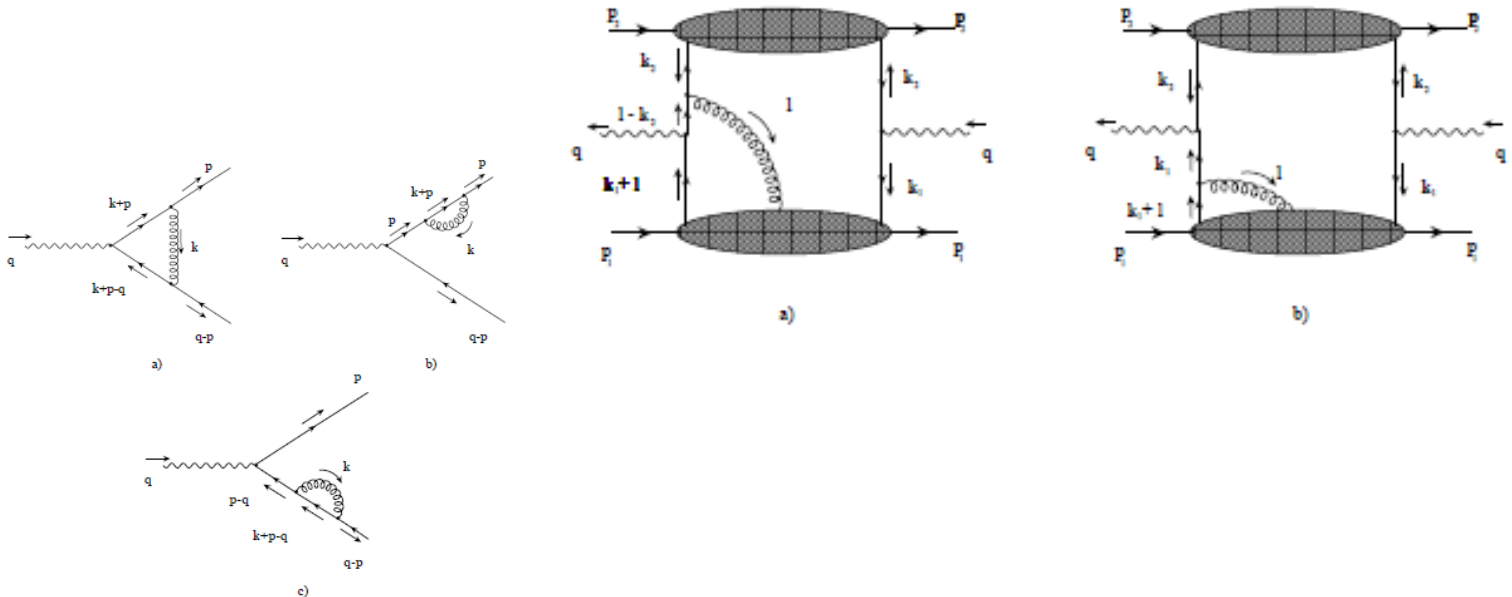
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Electromagnetic Gauge Invariance in DY process (IA,OT)

- Extra diagram – factor 2 in transverse (TM integrated) asymmetry
- Absent in pQCD



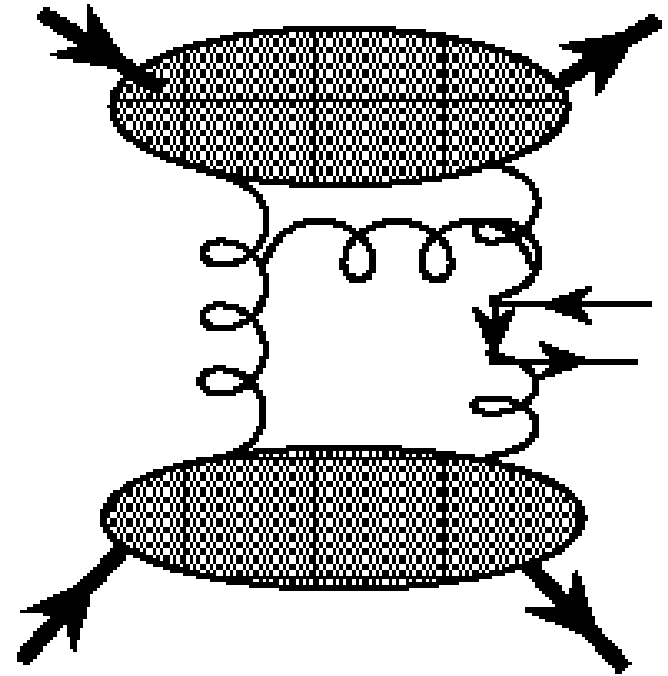


Comparison

- DY: imaginary part of EOM poles
 $\langle \dots G(x) \dots \rangle / x \sim \langle \dots A(x) \dots \rangle$
- Gluon DVMP – also!
- For free quarks target (recent discussion with D.Ivanov) – from gluon propagator: GI but much simpler in axial gauge with correct phase prescription
- One should expect similar situation with cuts for twist 3 SSA (for free quark? Onium??)
- Pion to 2 jets – poles with different prescriptions and claimed factorization breaking similar to that for Sivers function?
-

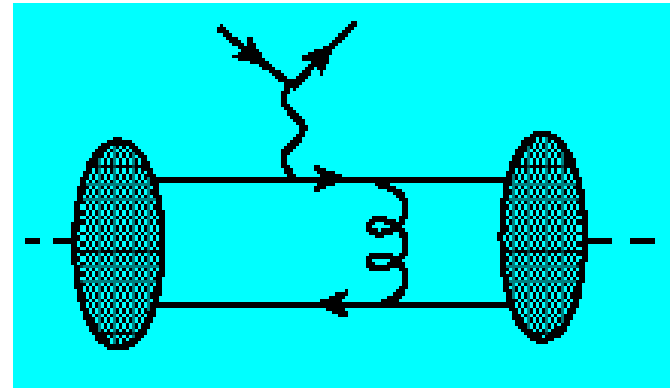
Double Diffraction: gluons

- One or both GPDs may be gluonic
- Complementary description of LHC DD (Higgs, Quarkonia, dijets)



Comparing collinear and (benchmark) KMR mechanisms

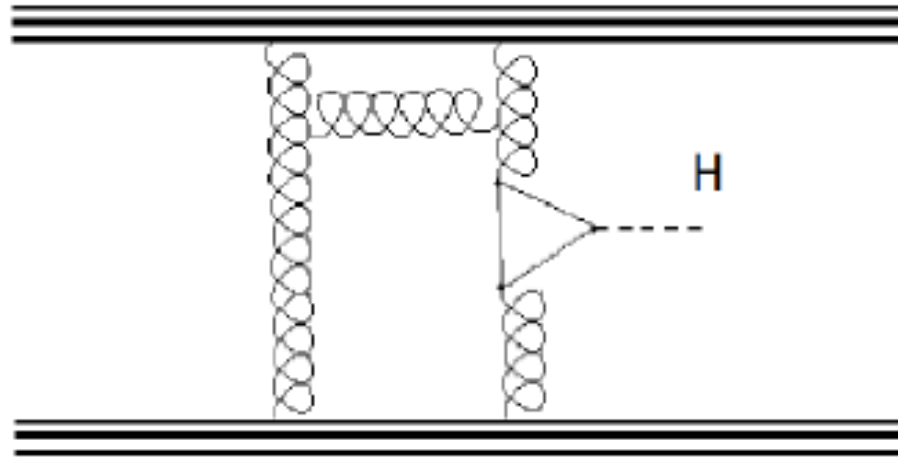
- Recall starting point – pion FF



- KMR analog for pion FF: Feynman mechanism
- Collinear – dominant for large Q^2
- Is something remained from this property after transition to GPDxGPD (wf overlap \rightarrow screening gluon)?!
- Collinear – no gap survival probability? Substituted by hard gluon?

Higgs boson production (AP,OT, in progress)

- Collinear factorization – hard gluon exchange





Gauge

- Axial gauge:

$$D_{\mu\nu}(k) = \frac{1}{k^2} \left[g_{\mu\nu} + (n^2 + \alpha k^2) \frac{k_\mu k_\nu}{(k, n)^2} - \frac{k_\mu n_\nu + k_\nu n_\mu}{(k, n)} \right]$$

- Partly fixing gauge:

$$n = p' + ap''$$



Probability amplitude

- Approximation of small skewnesses:

$$S_{fi} = -i \frac{10\pi^2 g_{(s)}^4}{(2V)^{\frac{5}{2}} v \sqrt{p'^0 p''^0 \tilde{p}'^0 \tilde{p}''^0 p^0}} \cdot \int_0^1 f(x) f(y) dx dy \cdot \delta(p' + p'' - \tilde{p}' - \tilde{p}'' - p)$$

- Gauge dependence cancelled



Conclusions

- AV ideas work, develop and lead to new relations between their consequences
- Imaginary parts in twist 3 transverse spin-dependent DY and exclusive DY amplitude have similar properties
- Non-universality and factorization breaking for inclusive and exclusive processes – similar?