# Lorentz boosts in interacting systems 

Eugene Stefanovich

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## OUTLINE

1. Wigner-Dirac theory of relativistic interactions
2. Currie-Jordan-Sudarshan "no interaction" theorem
3. Example: 2-particle interacting system
4. Example: $\mu-\tau$ oscillating neutrino pair
5. Example: fast moving unstable particle
6. Conclusions

## Basics of relativistic quantum theory

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quantum mechanics + principle of relativity $=$
Unitary representation of the Poincare group in the Hilbert space of the system

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| Generator of | Physical observable of | Operator |
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| time translations | Total energy = Hamiltonian | H |
| space translations | Total momentum | $\mathbf{P}$ |
| space rotations | Total angular momentum | J |
| boosts |  | K |
|  | Center-of-mass position | $\mathbf{R}=-\mathbf{c}^{2} / 2\left(\mathbf{K H}^{-1}+\mathbf{H}^{-1} \mathbf{K}\right)+$ spin terms |

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Commutation relations of the Poincare Lie algebra:

$$
\begin{gathered}
{[P, H]=0, \quad\left[J_{i}, P_{j}\right]=i \hbar \varepsilon_{i j k} P_{k}, \quad\left[K_{i}, P_{j}\right]=-\left(i \hbar / c^{2}\right) H \delta_{i j}, \quad[K, H]=-i \hbar P,} \\
{\left[R_{i}, P_{j}\right]=i \hbar \delta_{i j}, \cdots}
\end{gathered}
$$

## System of two non-interacting particles in 1D

$$
\begin{aligned}
& \mathrm{H}_{0}=\mathrm{h}_{1}+\mathrm{h}_{2} \\
& \mathrm{P}_{0}=\mathrm{p}_{1}+\mathrm{p}_{2} \\
& \mathrm{~K}_{0}=\mathrm{k}_{1}+\mathrm{k}_{2}
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\end{array} \quad \mathrm{x}_{\mathrm{j}}(\mathrm{t})=\mathrm{e}^{\mathrm{i} \mathrm{H}_{0} \mathrm{t}} \mathrm{x}_{\mathrm{j}} \mathrm{e}^{-\mathrm{iHot}}=\mathrm{x}_{\mathrm{j}}(0)+\mathrm{v}_{\mathrm{j}} \mathrm{t}
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Lorentz transformations for times and positions of events

$$
\begin{gathered}
x^{\prime}=x \cosh \theta-c t \sinh \theta \\
\mathrm{t}^{\prime}=-(\mathrm{x} / \mathrm{c}) \sinh \theta+\mathrm{t} \cosh \theta
\end{gathered}
$$

## Forms of Relativistic Dynamics

1．A．M．DIRAC


## Forms of Relativistic Dynamics

1. A. M. DinAC

2. Generators of the Poincare Lie algebra should be modified in the presence of interactions (Dirac's instant form)

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& \mathrm{H}=\mathrm{H}_{0}+\mathrm{V} \\
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1. A. M. Dirac <br> S. Jossrs Codece, Canbridge, Englind
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2. Interactions $V$ and $\mathbf{Z}$ must be chosen so that all Poincare commutators are preserved, e.g.,

$$
\left.\left[K_{i 0}+Z_{i}\right) P_{j}\right]=-\left(i \hbar / c^{2}\right)\left(H_{0}+V\right) \delta_{i j}
$$

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Lorentz transformations are not valid

$$
\begin{gathered}
x^{\prime} \neq x \cosh \theta-c t \sinh \theta \\
t^{\prime} \neq-(x / c) \sinh \theta+t \cosh \theta
\end{gathered}
$$

## Relativistic Invariance and Hamiltonian Theories of Interacting Particles*

D. G. Currifit T. F. Jordan, and E. C. G. Sudarshaf

Department of Physics and Aslromomy, University of Rochester, Rochester, New York

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"...the combined assumptions of Lorentz symmetry and Lorentz transformation of particle positions rule out any interaction."

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"...the combined assumptions of Lorentz symmetry and Lorentz transformation of particle positions rule out any interaction."

How to resolve this paradox?

- reject Lorentz (Poincare) symmetry? No.
- reject Hamiltonian dynamics? No.
- reject existence of particle positions? No.
- reject existence of particles? No.
- reject Lorentz transformations for positions (and momenta) of interacting particles? Yes!

Two particles with a fully relativistic interaction in 1D
Shields, Morris, Ware, Su, Stefanovich, Grobe, Phys.Rev. A 82 (2010) 052116

Potential: $\mathbf{U}(\mathbf{r})=\mathrm{U}_{0} /\left[\mathrm{r}^{2}+\mathrm{a}^{2}\right]^{1 / 2}$ introduced via Bakamjian-Thomas prescription

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$\square$
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constant t
This argument is not valid fo....
Cause

A fully relativistic model of neutrino propagation and oscillation
Stefanovich 2012

# A fully relativistic model of neutrino propagation and oscillation 

Stefanovich 2012


For a general relativistic Hamiltonian in the Hilbert space of the $v_{\mu}-v_{\tau}$
system, the diagonal element $\Omega_{\mu}(p)$ is subject to only one condition:

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\begin{aligned}
E_{2}(p) & <\Omega_{\mu}(p)<E_{3}(p) \\
H & =\left[\begin{array}{ll}
\Omega_{\mu}(p) & f(p) \\
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Trajectory of the $\mu$-neutrino component

$$
L(t)=c t+\Delta L(E, t)
$$

$$
\Delta L(E, t)=\left(d \Omega_{\mu}(p) / d p-c\right) g(E, t)
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$\mathbf{g}(\mathbf{E}, \mathbf{t})$ is an oscillating function of $\mathbf{t}$ assumption: $\mathrm{d} \Omega_{\mu}(\mathrm{p}) / \mathrm{dp}=1.0000237 \mathrm{c}$

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Decay law of a moving unstable particle: a relativistic model
Stefanovich 1996, Khalfin 1997, Shirokov 2004

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| at rest | $\omega(0, t) \approx \exp (-\Gamma t)$ |
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| in the moving frame (Einstein's theory) | $\omega(\theta, t) \approx \exp (-\Gamma t / \cosh \theta)=\exp (-\chi)$ |
| in the moving frame (Wigner-Dirac theory) | $\omega(\theta, t) \approx \exp (-\chi)+\Delta(\chi, \theta)$ |

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Properties of the imaginary unstable particle a:
mass $\mathrm{m}_{\mathrm{a}}=1000 \mathrm{MeV} / \mathrm{c}^{2}$
width $\Gamma=20 \mathrm{MeV} / \mathrm{c}^{2}$

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The expected effect for muon $\left(\approx 10^{-18}\right)$ is too small to be observed. $: *$

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3. See, however, arXiv:1211.2913!
