

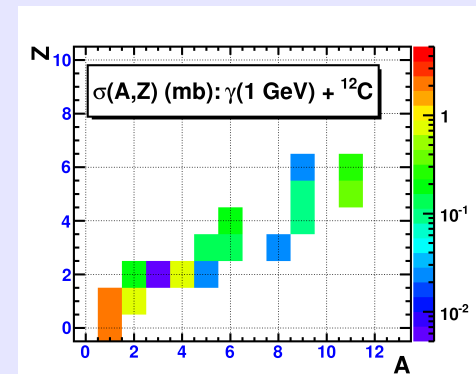
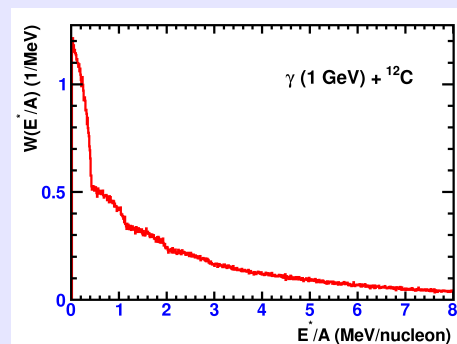
Multifragment break-up of ^{12}C in photonuclear reactions: a theorist's point of view

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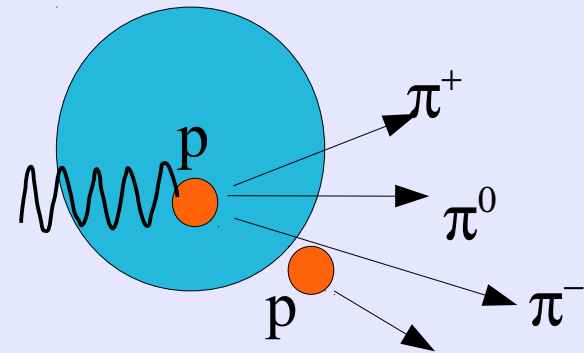
XXII International Baldin Seminar on
High-Energy Physics Problems
Dubna, Russia, September 15-20, 2014



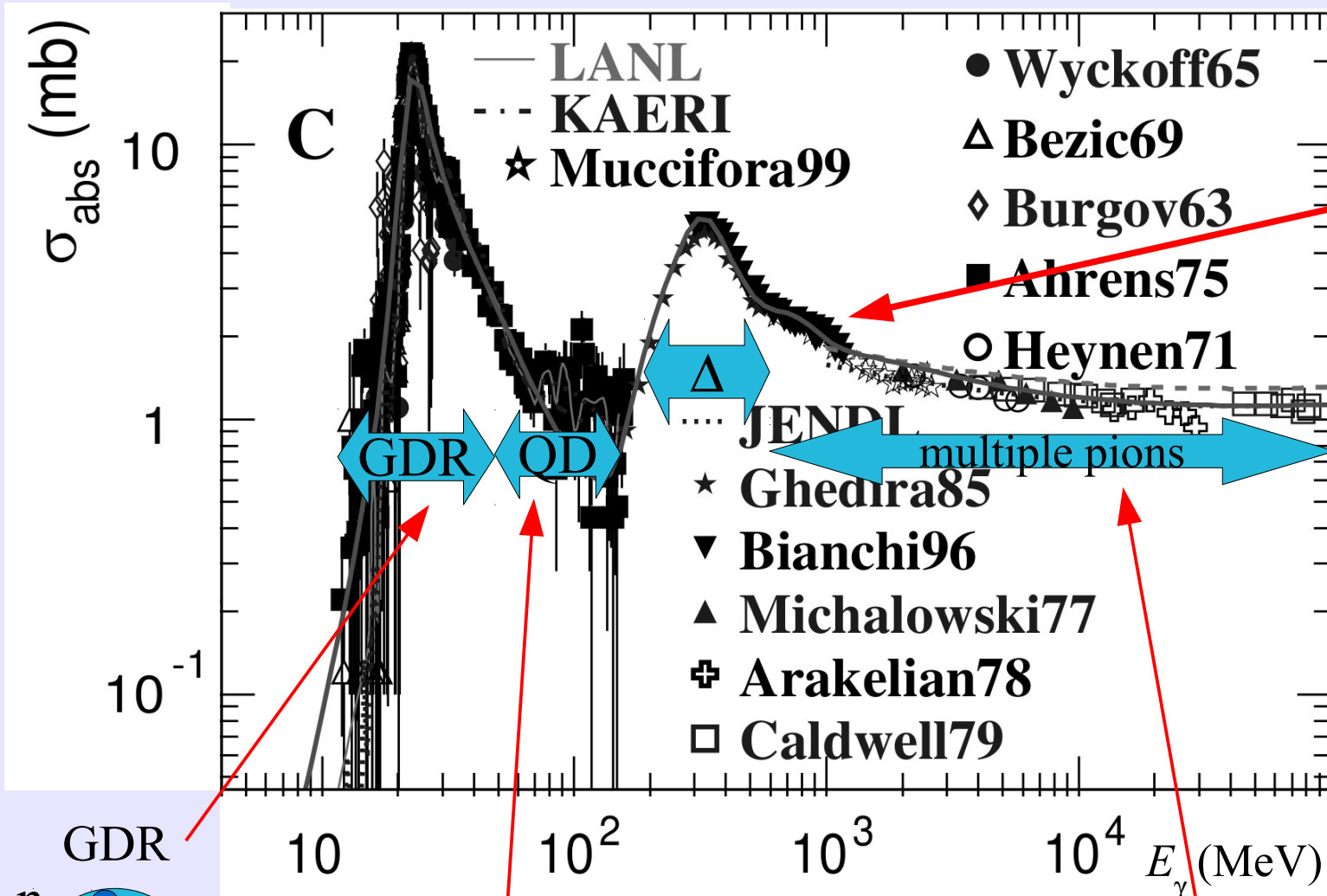
Content

- Absorption of ~ 1 GeV photons by ^{12}C – an experiment at the GRAAL facility to study multifragment decays of ^{12}C
 - Details of this experiment in Prof. Nedorezov's talk
 - Here a theorist's point of view is presented:
 - a model to describe the break-up of ^{12}C ;
 - comparison with photonuclear reactions on heavy nuclei and induced by other projectiles;
 - relations to electromagnetic dissociation of light nuclei
 - Comparison with the rates of multifragment decays measured at GRAAL
 - Fragmentation reactions in carbon-ion therapy
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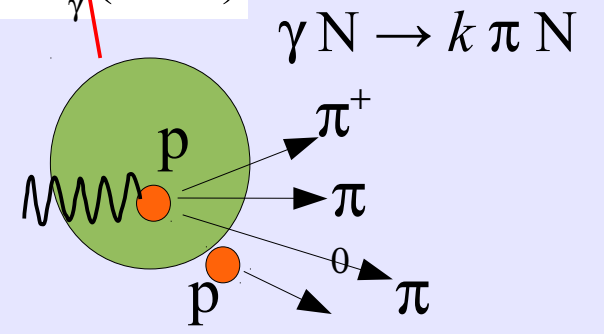
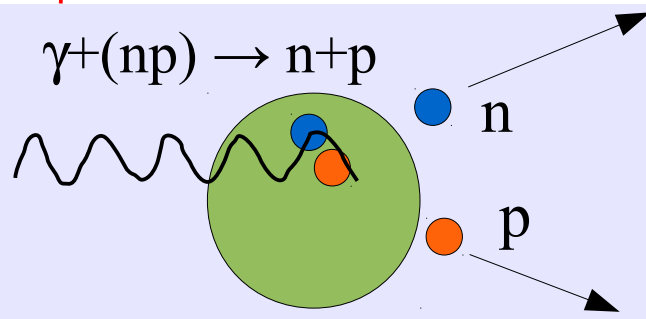
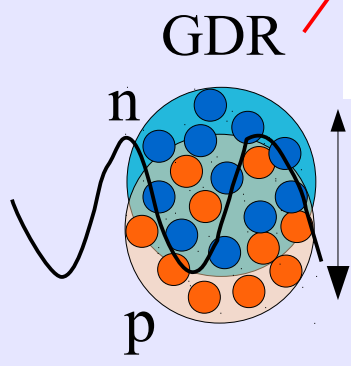
I. Modelling of photonuclear reactions



Photoabsorption on carbon: a variety of processes

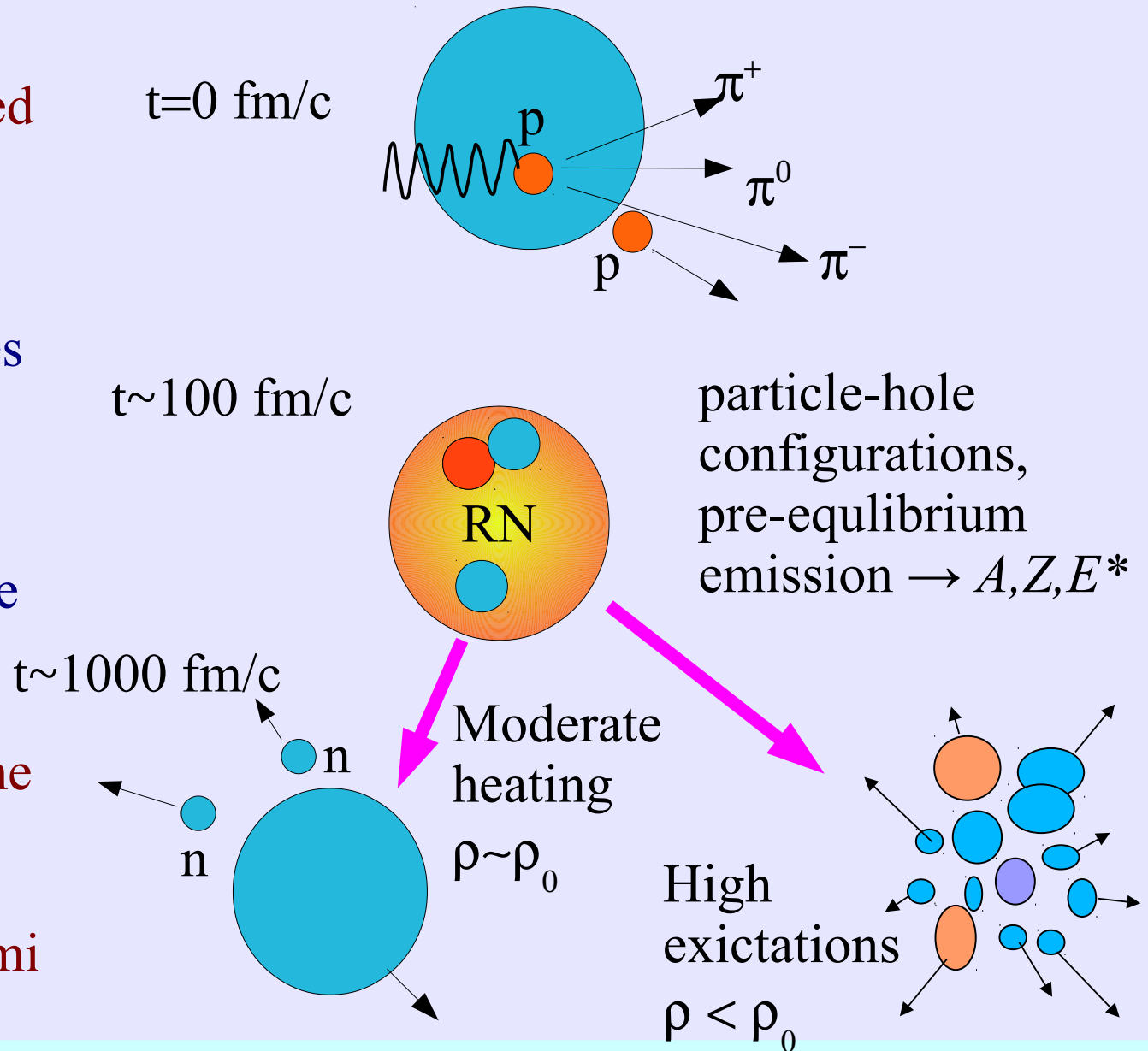


This experiment at GRAAL:
 700-1500 MeV



Basics of the photonuclear reaction model. It is a part of the Relativistic Electromagnetic DIssociation (RELDIS) model

- Firstly, an Intranuclear cascade process initiated by a projectile is simulated
- Secondly, the properties of a residual nucleus (A, Z, E^*), which is left after the completion of the intranuclear cascade are calculated
- Thirdly, the decay of the thermalized residue is modelled by evaporation/SMM/Fermi break-up models



Let's estimate the number of partitions to split a nucleus

Famous Euler's problem: find the number of ways, $P(A_0)$, an integer A_0 can be represented as a sum of integers, $A_0 = A_1 + A_2 + A_3 + \dots$

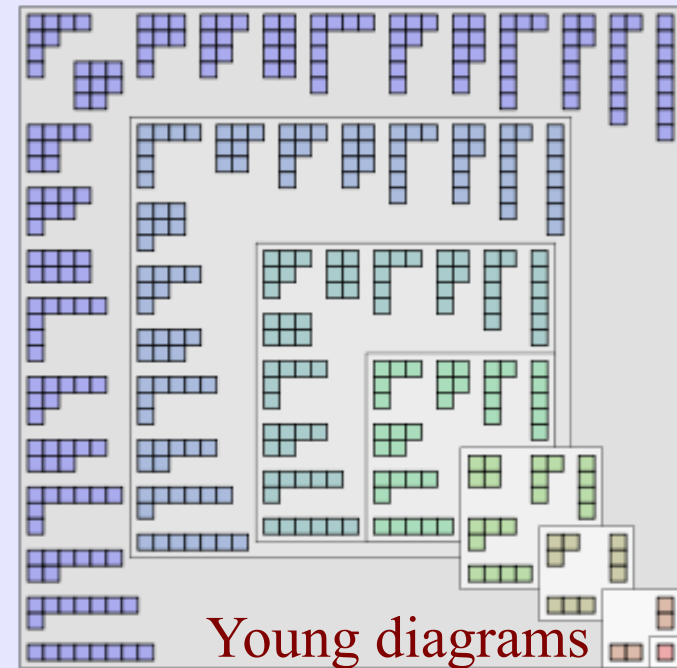
The total number of partitions is rapidly increasing with A_0 :

$$P(10)=42, \quad P(12)=77, \quad P(16)=231,$$

Large enough already for decays of excited carbon and oxygen nuclei!

$$P(50)=2.17 \cdot 10^5, \quad P(100)=2 \cdot 10^8$$

See A.S. Botvina, A.D. Jackson,
I.N. Mishustin, Phys. Rev. E62 (2000) R64



Young diagrams
for partitions

At large A_0 the result is well approximated by asymptotic

Hardy-Ramanujan formula, with the average partition multiplicity:

$$P(A_0) = \frac{1}{4\sqrt{3}A_0} \exp\left(\pi\sqrt{\frac{2A_0}{3}}\right) \quad \langle M \rangle = \frac{1}{\pi} \sqrt{\frac{3A_0}{2}} \ln\left(\frac{6A_0}{b\pi^2}\right) \quad b = 0.3150$$

Specific to light nuclei ($A < 17$): Fermi break-up model

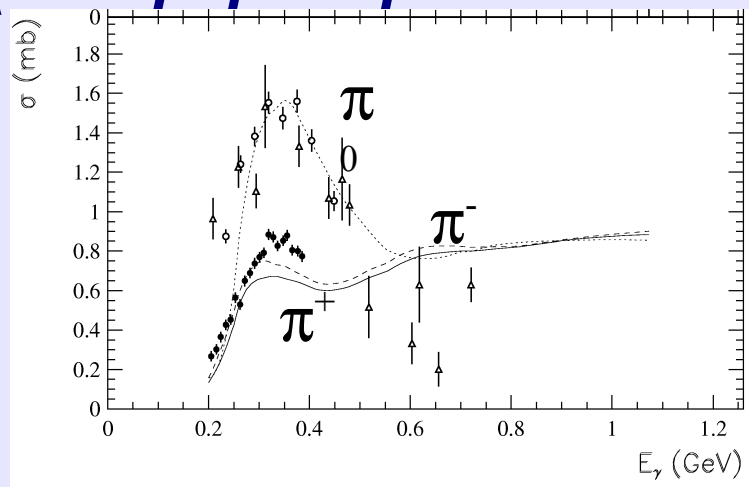
- Excitation energies of light nuclei can be comparable to their total binding energies - explosive decays.
- Fermi break-up model implemented by A.S. Botvina and co-authors
- Fragments are considered in their ground states and also in low-energy excited states stable to nucleon emission.
- The probability of a given decay channel is defined by its statistical weight.
- The list of possible decay channels is quite long, e.g. ~ 200 channels for ^{12}C , ~ 1000 for ^{16}O



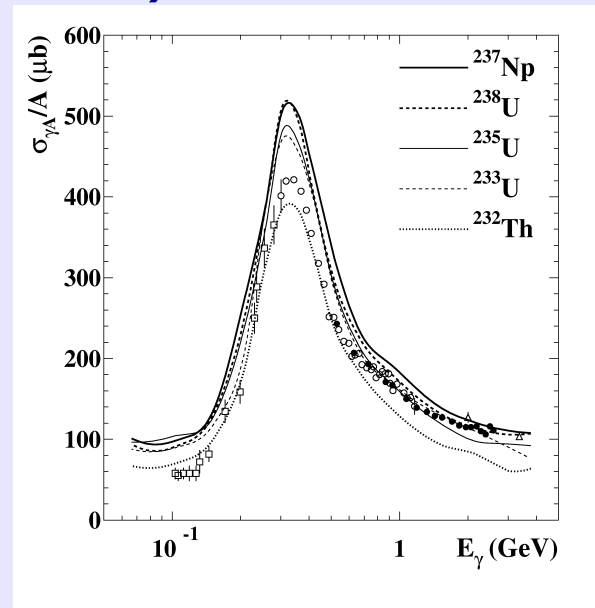
*Wassily Kandinsky
"Several circles", 1926
Guggenheim
Museum, NY*

RELDIS model developed at INR, NBI, GSI, FIAS

(25+ papers published since 1995)

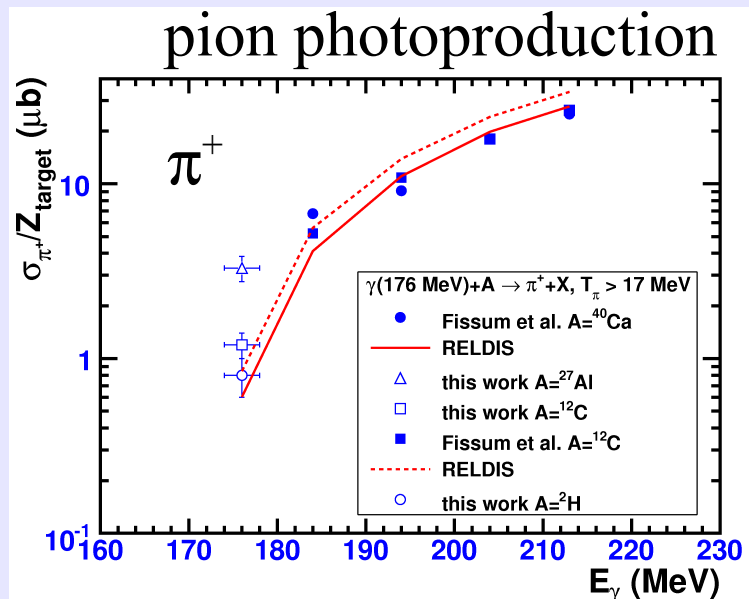


A.S. Iljinov, ..., I.P. et al.,
Nucl. Phys. A 616(1997)575

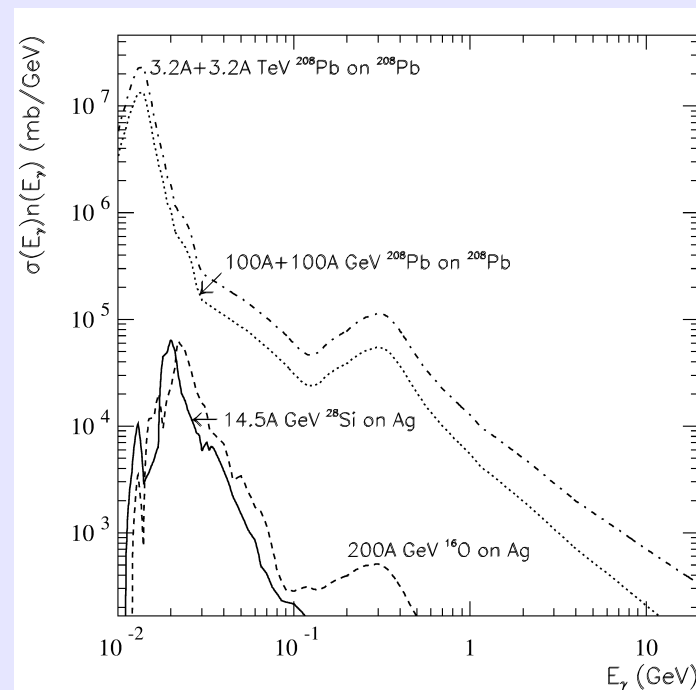


I.P. et al.,
Eur. J. Phys. A
24(2005)69

fission of
heavy nuclei by
photons



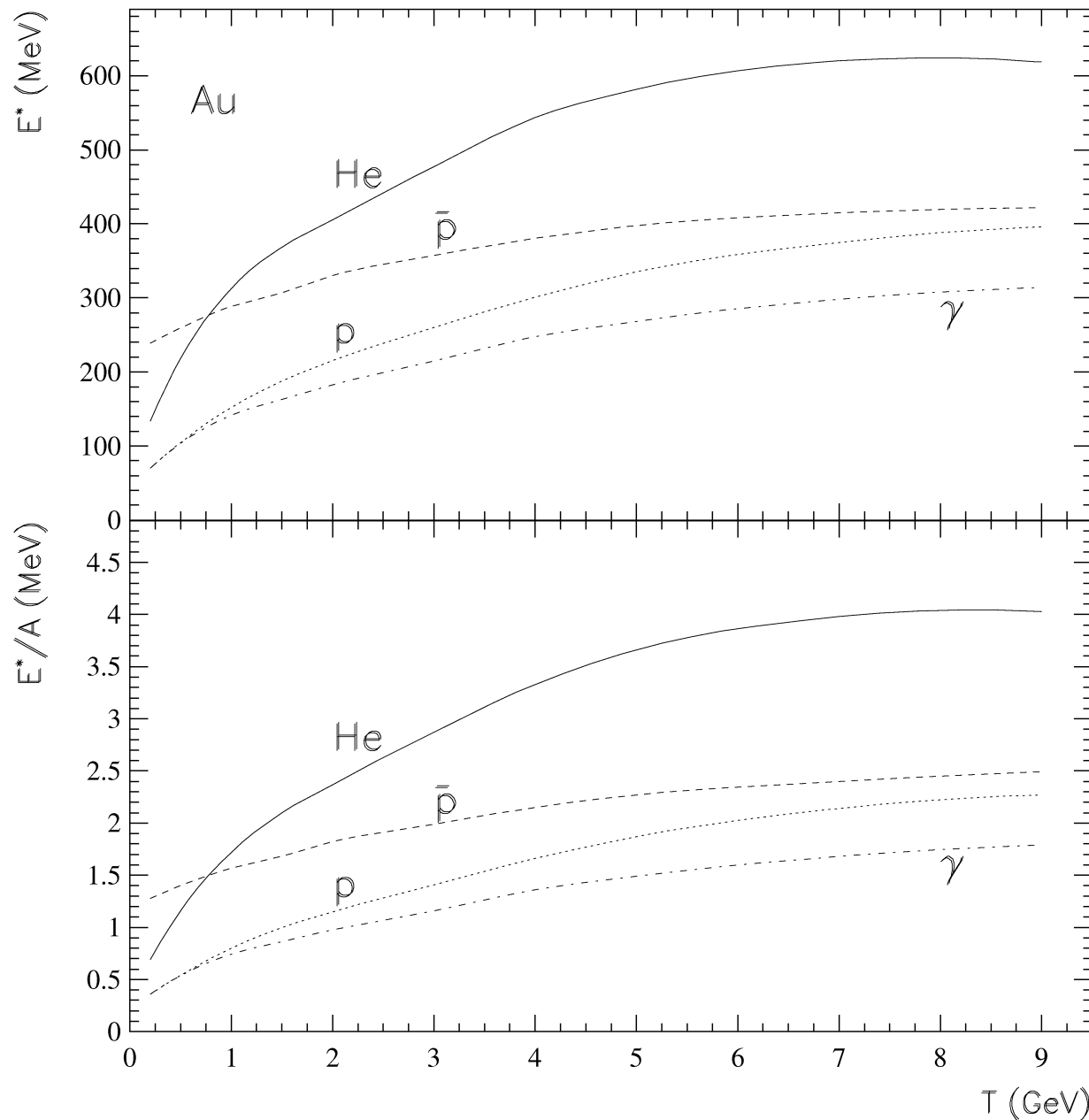
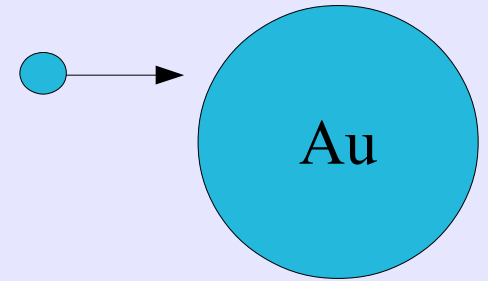
P.Golubev, ..., I.P. et al.,
Nucl. Phys. A 806(2008)216



I.P. et al.,
Phys. Rev. C
57(1998)1920

Electromagnetic
dissociation at
CERN SPS
SIS (GSI),
AGS (BNL)
(see below)

Average excitation energy of a residual nucleus at the end of the intranuclear cascade.



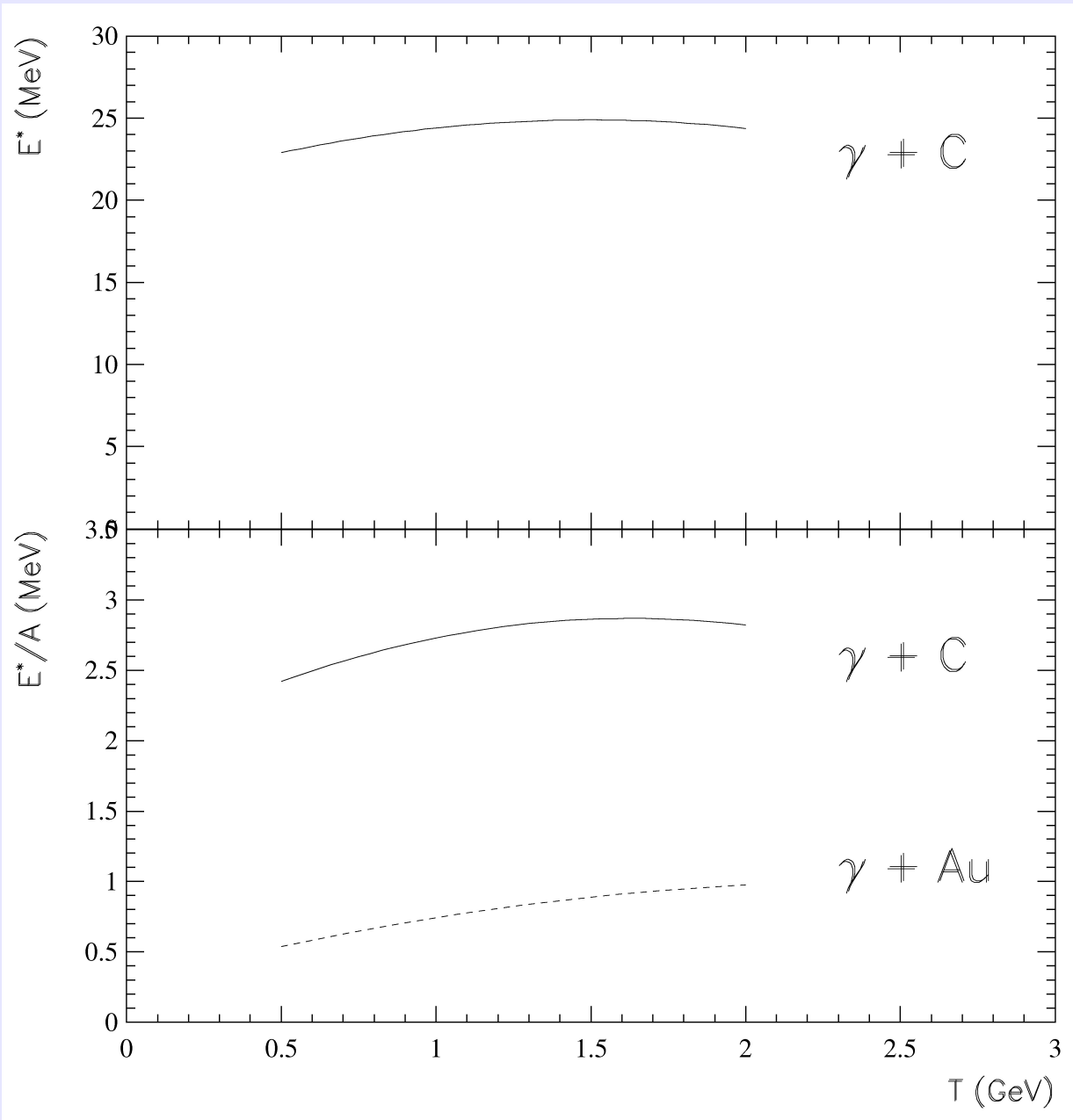
To be compared:
proton,
antiproton,
photon, and ^4He projectiles

Photons heat the nucleus
less effectively compared
to other projectiles

Antiprotons work good
at low energy due to their
annihilation

Ions (e.g. ^4He) is the best
option to heat up the nucleus

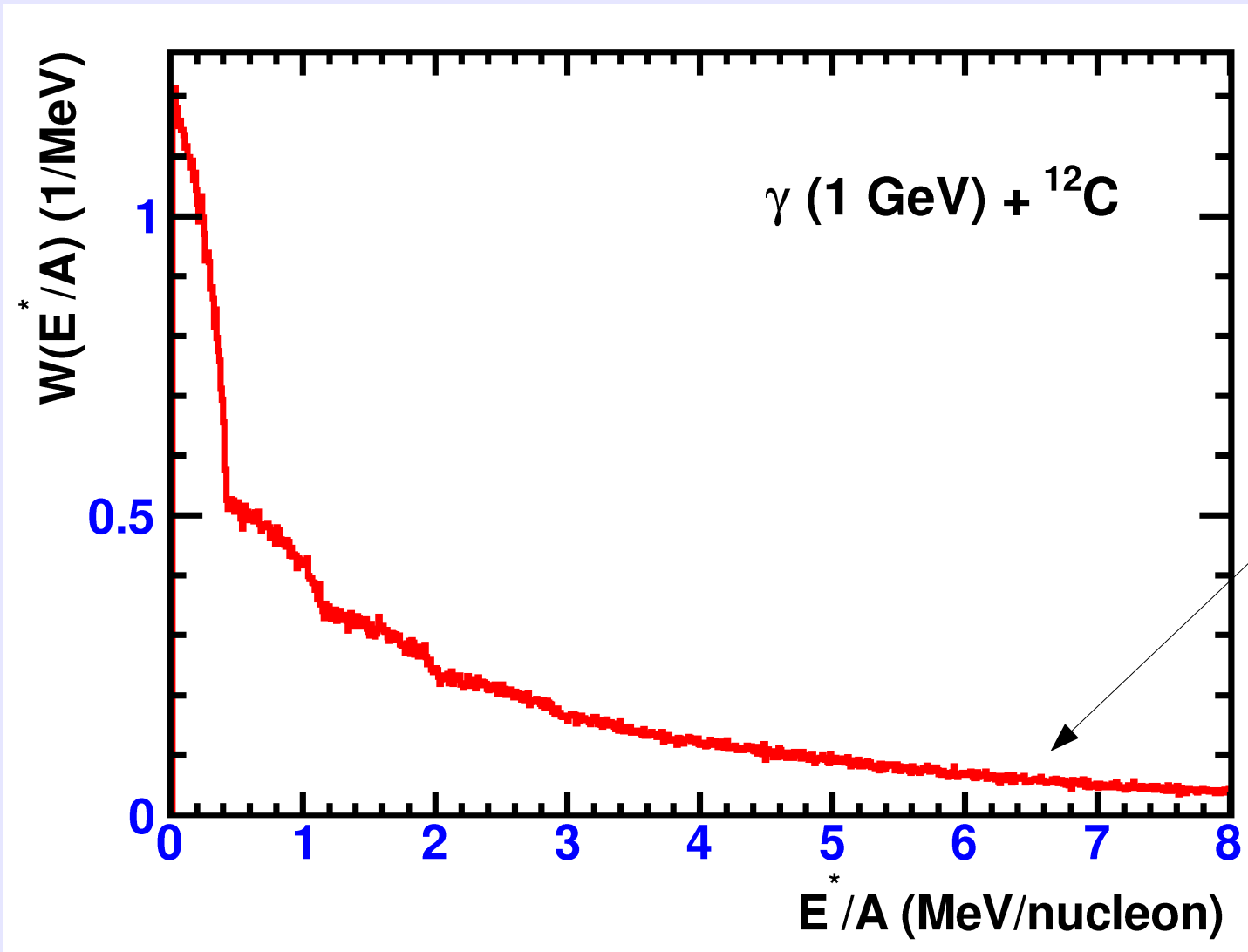
Comparison of photon absorption by ^{12}C and ^{197}Au



Due to a small number of nucleons in ^{12}C the average excitation energy calculated per nucleon of the residual nucleus is much higher compared to the photoabsorption on ^{197}Au .

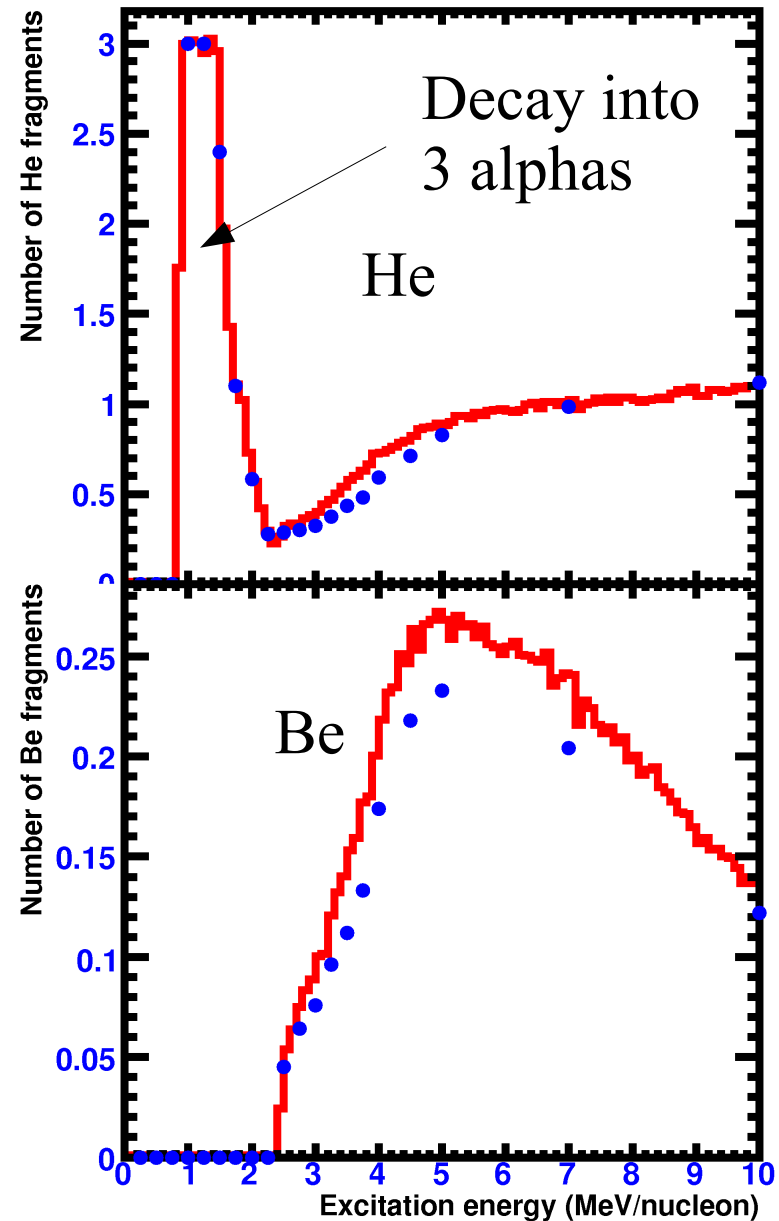
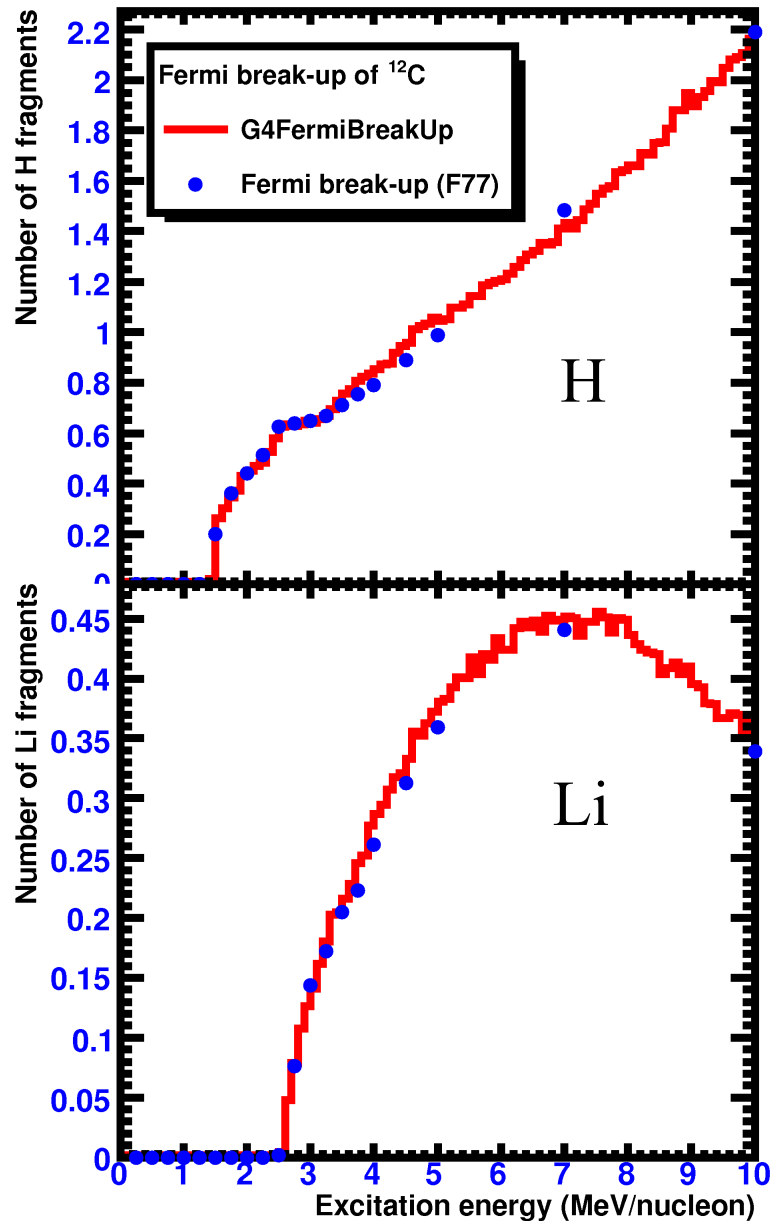
The onset of the explosive break-up is expected at ~ 3 MeV/nucleon

The average values are not very informative: the distribution of excitation energy is very wide

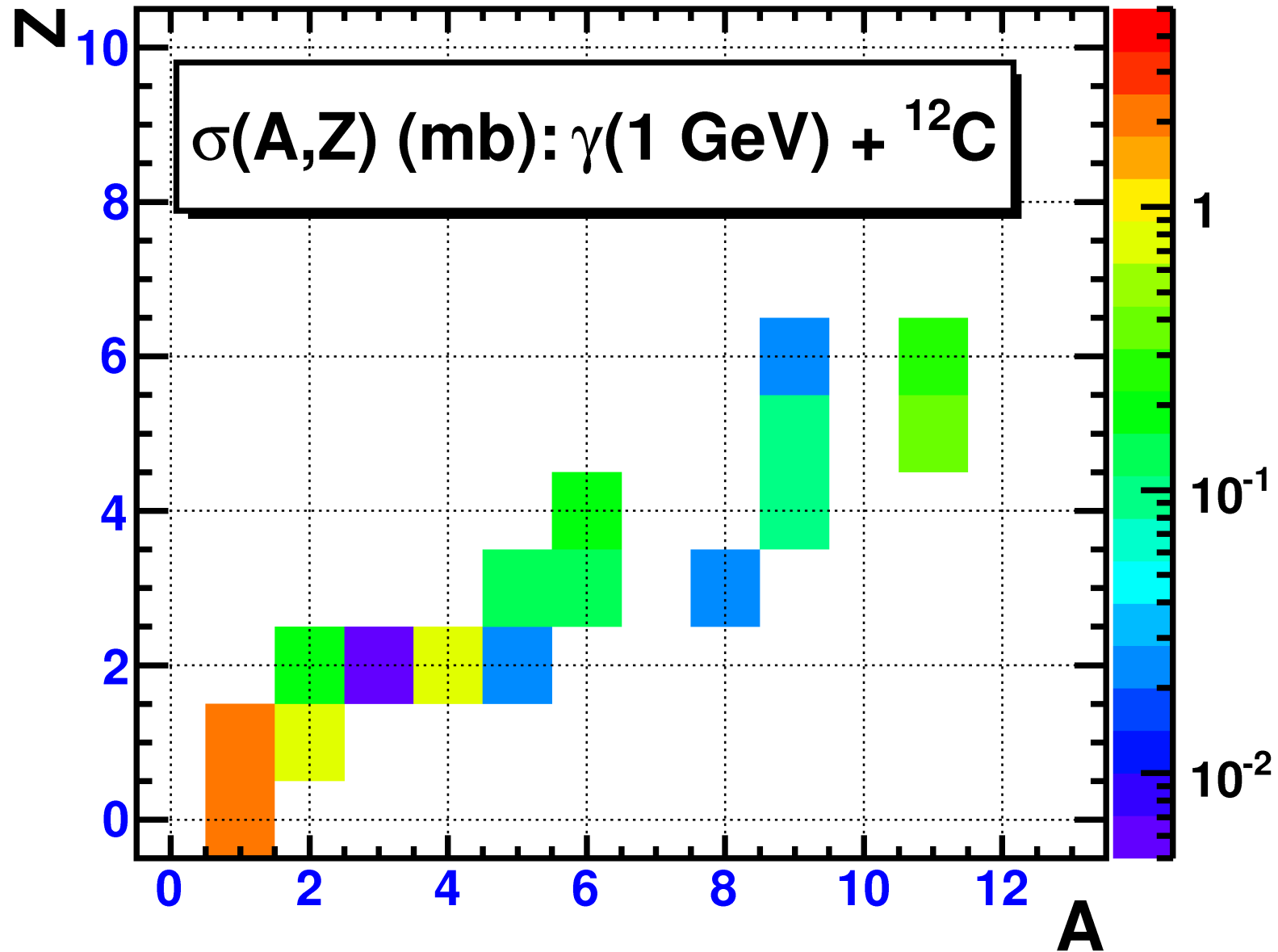


Break-up
into
individual
nucleons

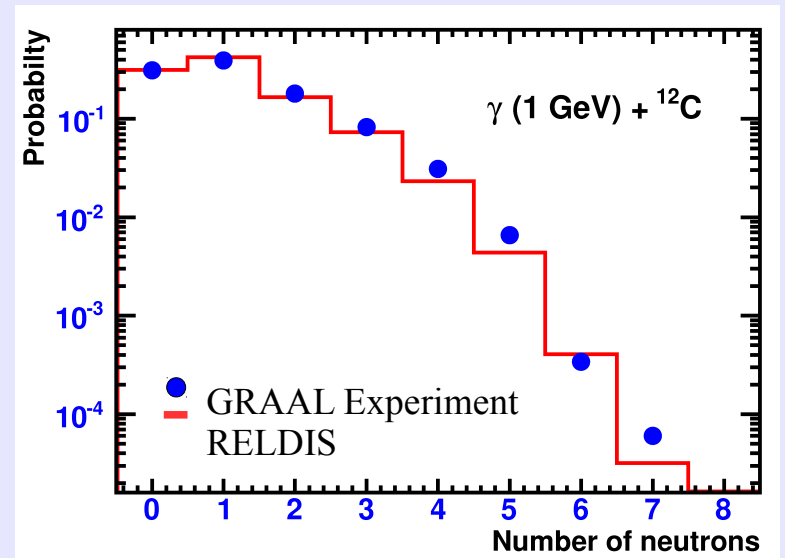
Average number of fragments of ^{12}C of a given element: Fermi Break-up model by A.S. Botvina et al. vs Geant4 implementation



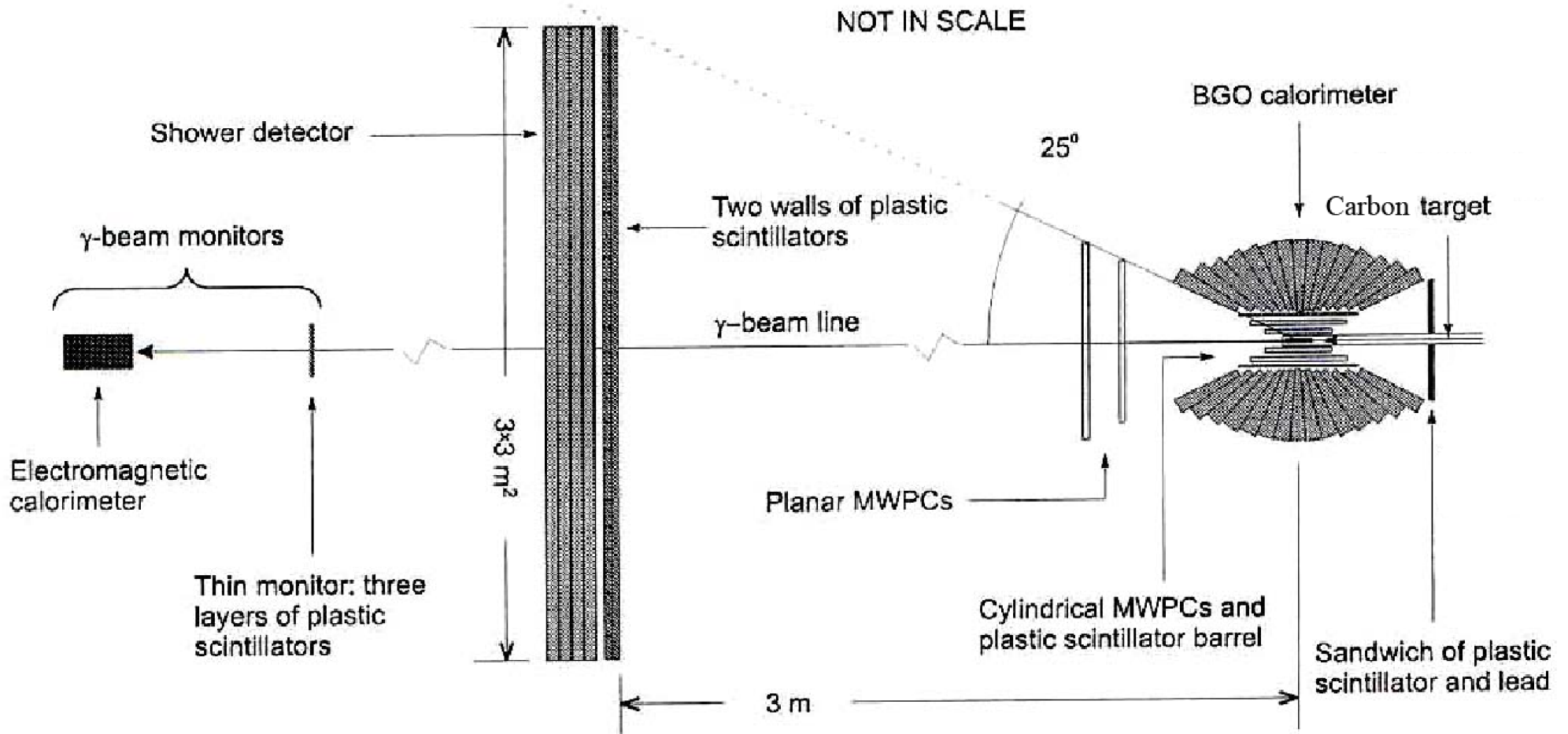
Calculated map of nuclear fragments



II. Comparison with the GRAAL data

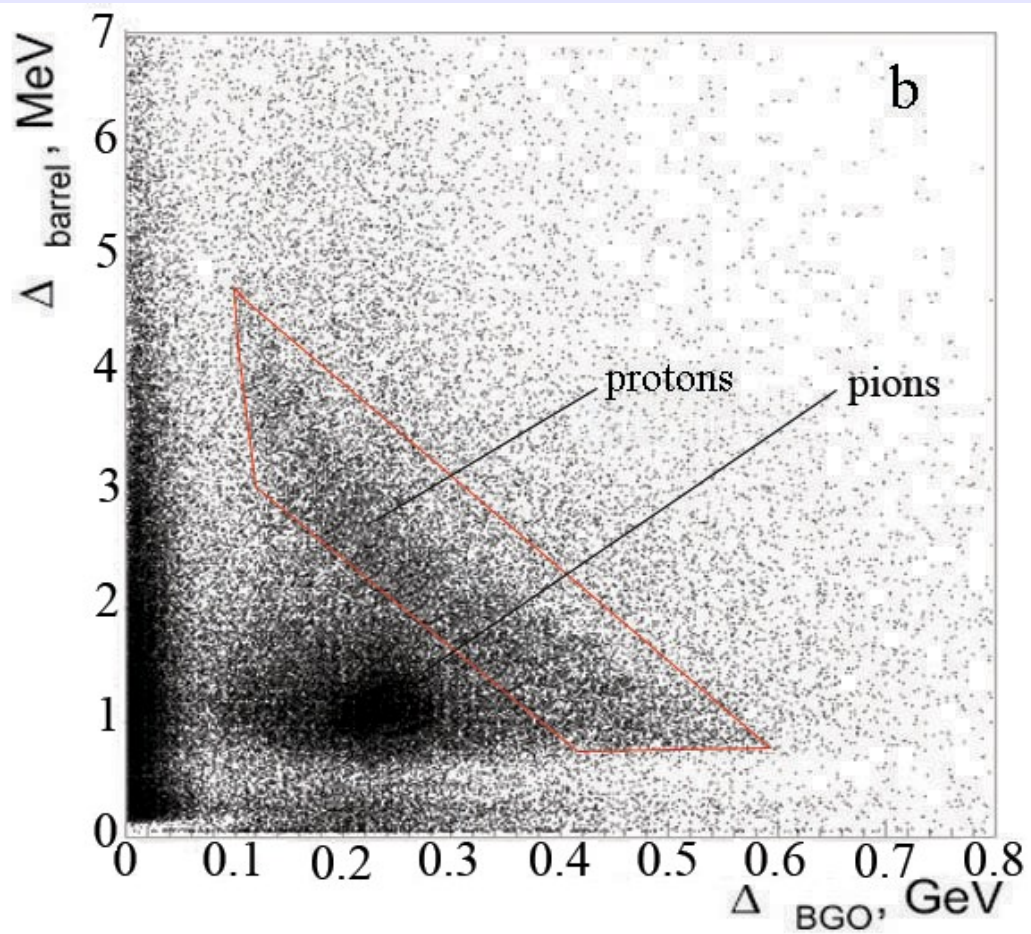
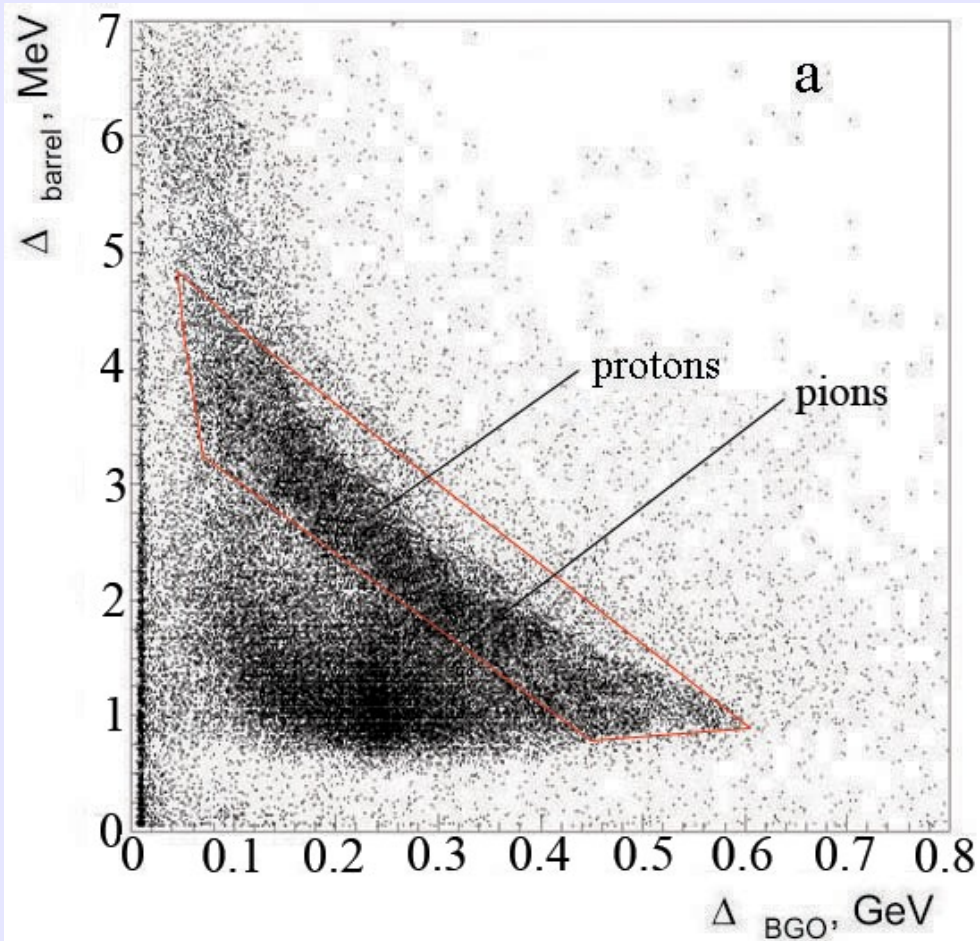


General layout of the GRAAL setup

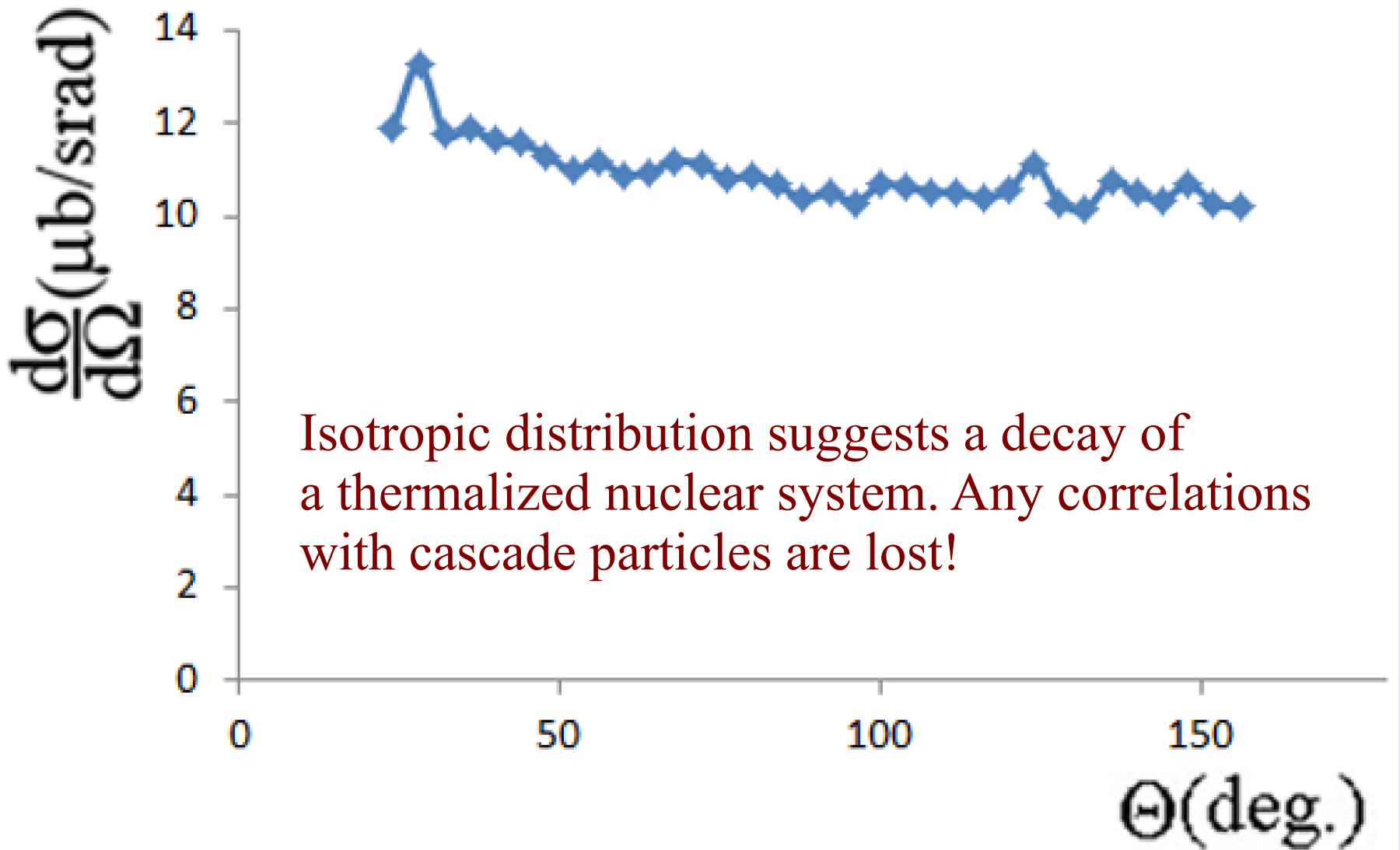


4π detector with a good efficiency of registration of protons and neutrons. This is important for studies of decays of highly excited nuclei.

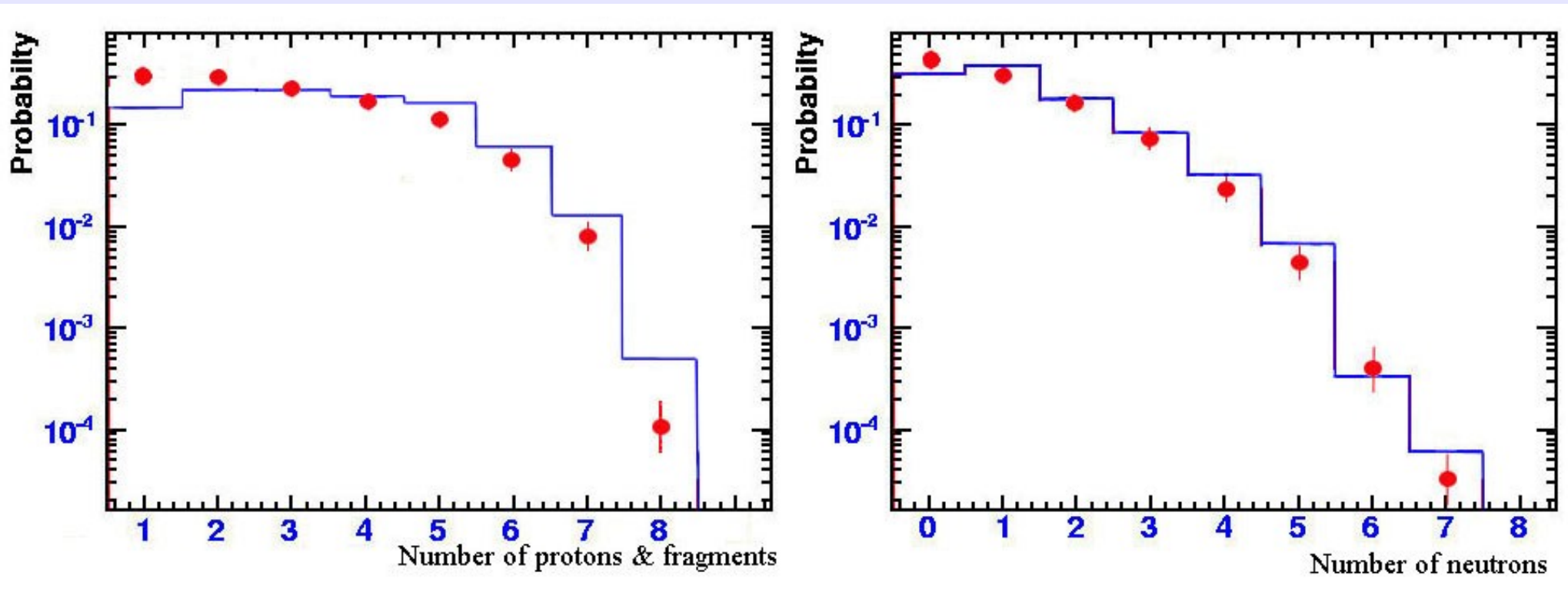
Distinguishing protons from pions in the BGO calorimeter. Simulated (a) and detected (b) events are shown



Measured angular distribution of nucleons produced in photodisintegration of ^{12}C in events with more than 7 nucleons.



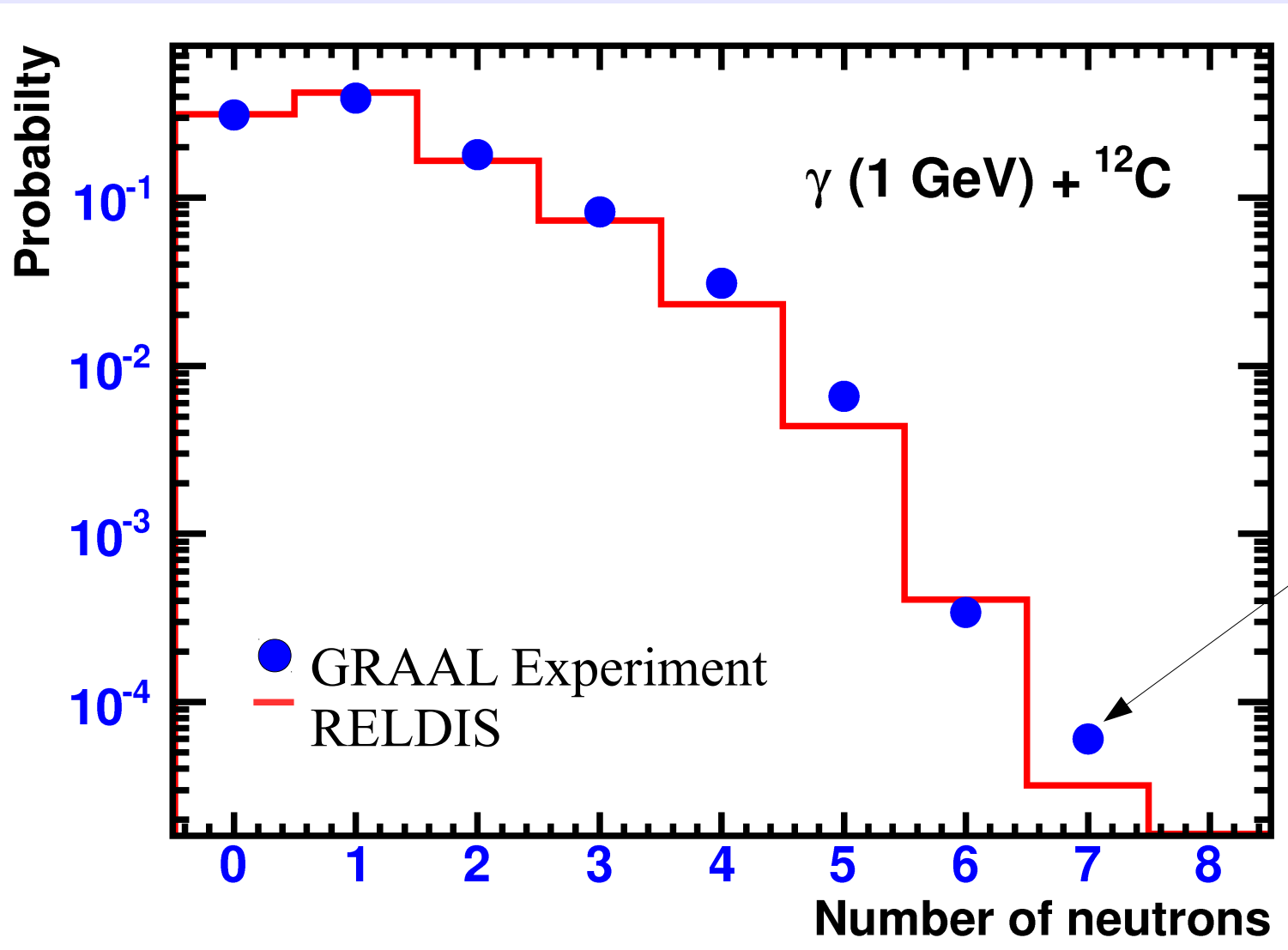
Distributions in the number of protons and neutrons



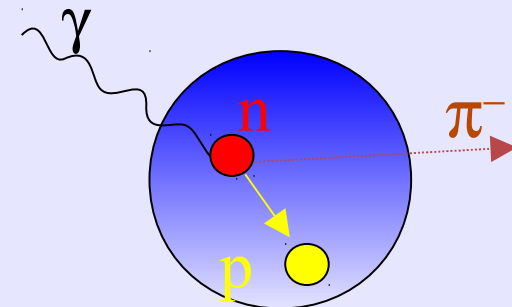
Measured (points) and calculated (histograms) probabilities of photodisintegration of ^{12}C at 0.7-1.5 GeV with given numbers of protons (left) and neutrons (right). Only statistical errors are shown.

Neutron multiplicity distributions

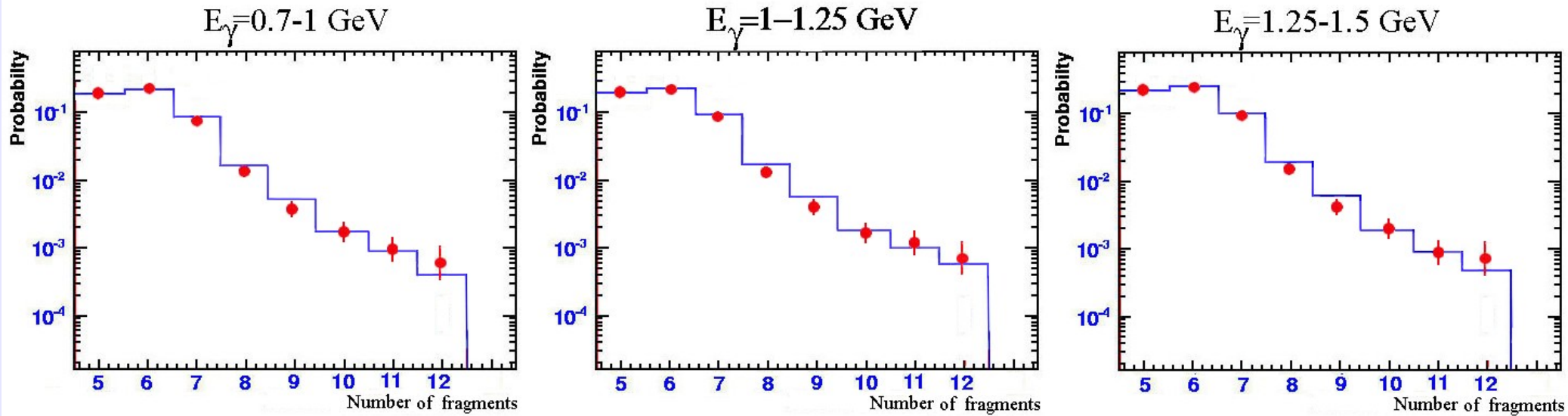
Demonstrate the possibility of a complete disintegration of ^{12}C into individual nucleons



Very rare charge-exchange events



Measured and calculated probabilities of ^{12}C photodisintegration with a given number of fragments.

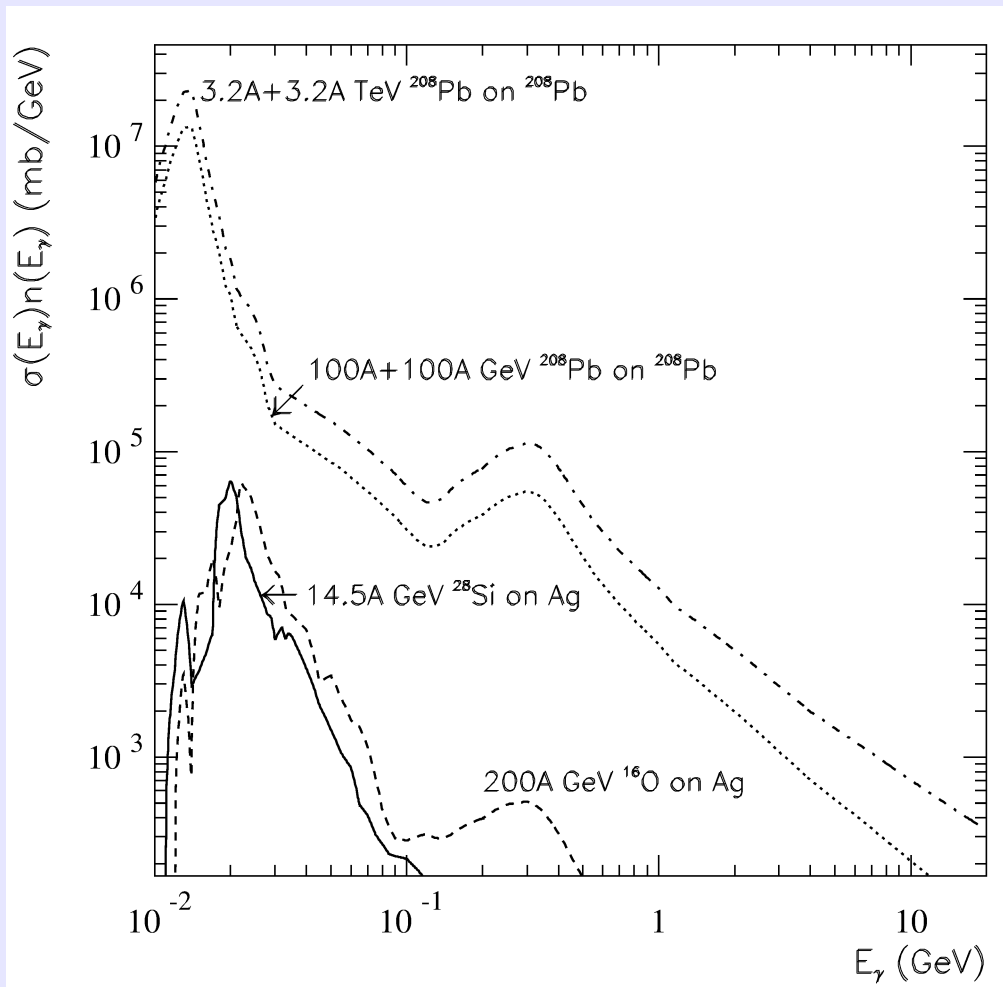


Very good agreement with calculated distributions in all three intervals of photon energy!

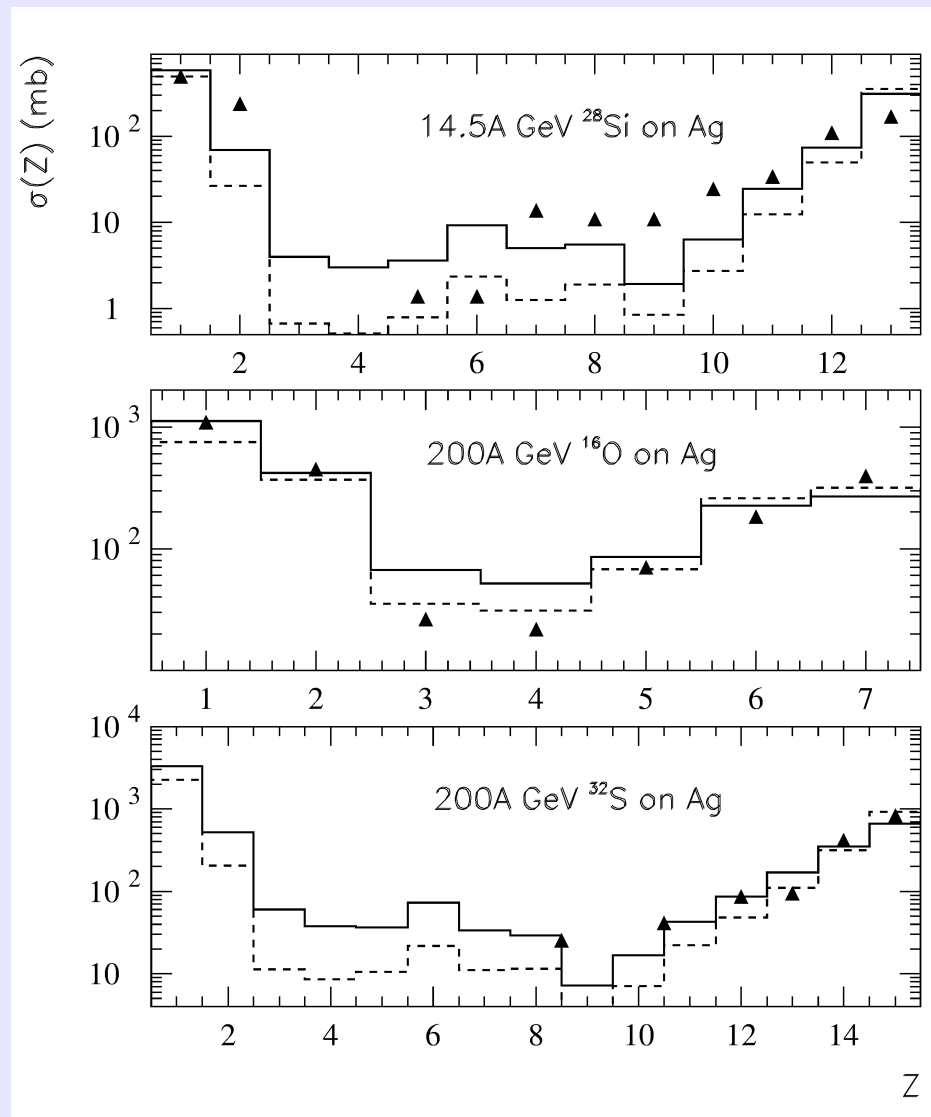
On average, 8 fragments are produced once per ~ 100 events, while 12 fragments (complete disintegration) once per 2000 events.

Electromagnetic dissociation of light nuclei in nuclear emulsion

Wide distributions of virtual photon energies



Multifragment decays are also seen

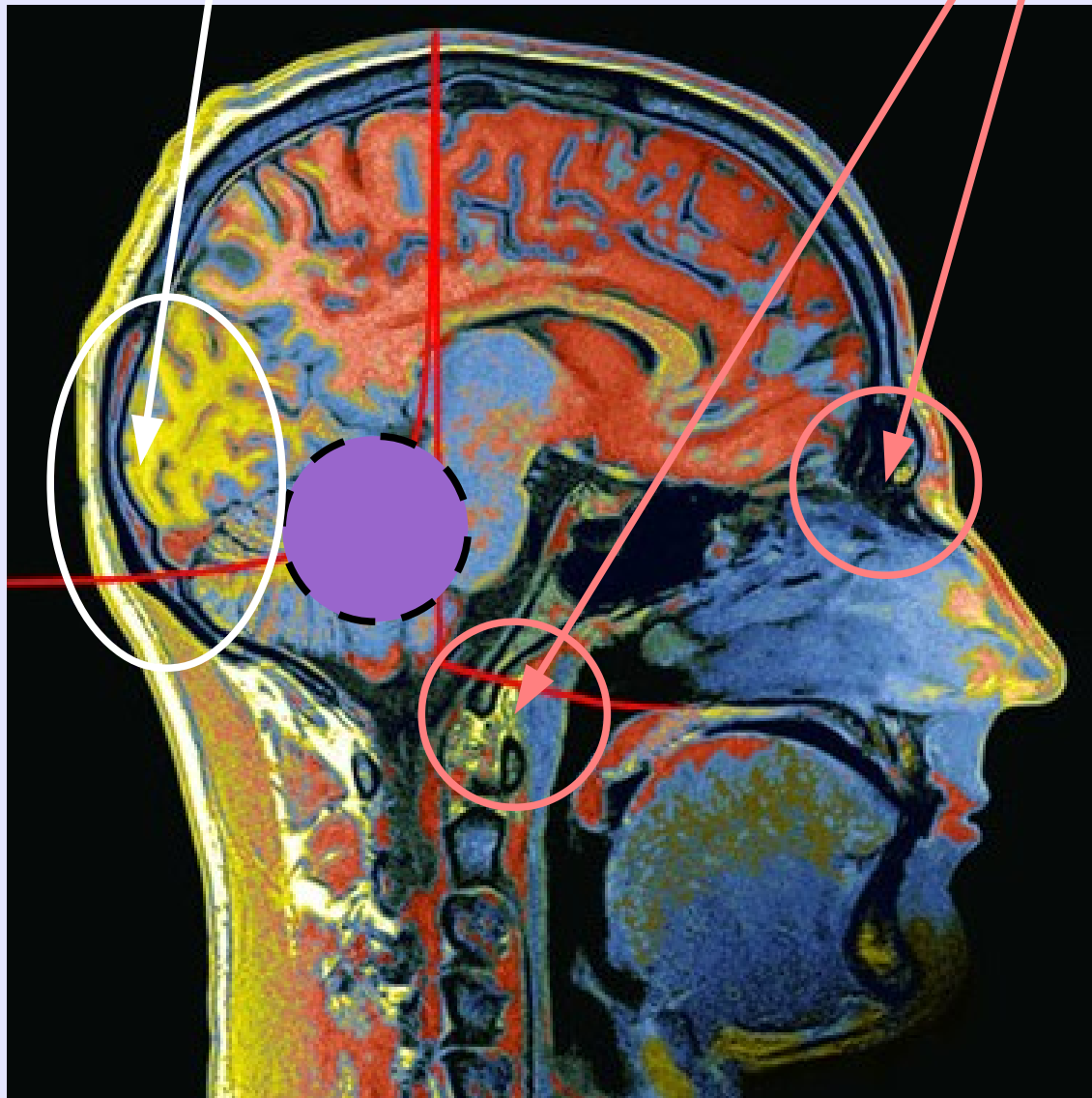


Described w/o multiple photon absorption. I.P. et al. Phys. Rev. C **57** (1998) 1920

III. Why it is important to study the fragmentation of ^{12}C ?



Carbon-ion therapy of cancer: nuclear beams are focused on tumor thus sparing healthy tissues and organs at risk

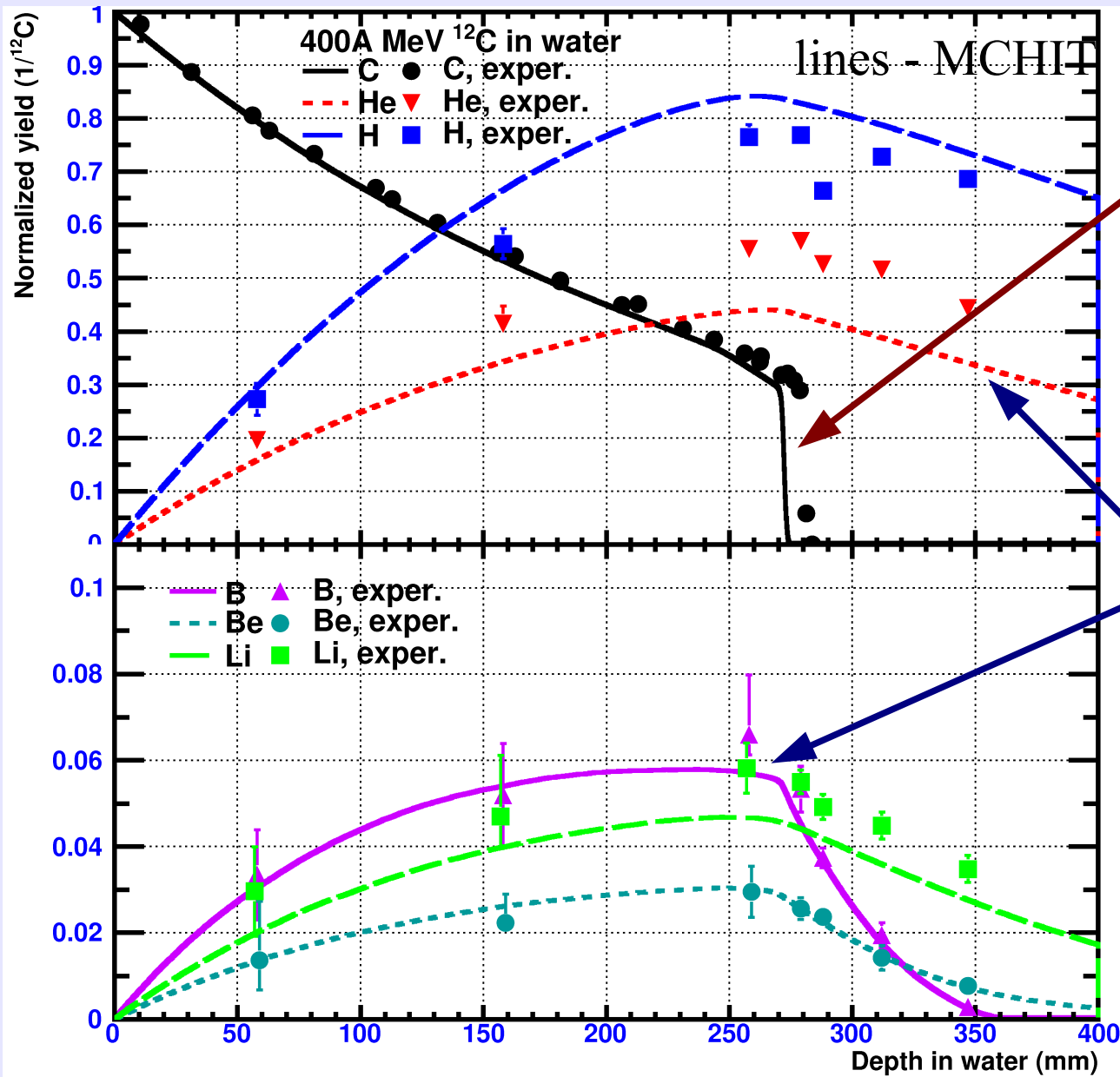


^{12}C beam

Some 10 000 patients treated worldwide so far.

About 100 000 – with protons

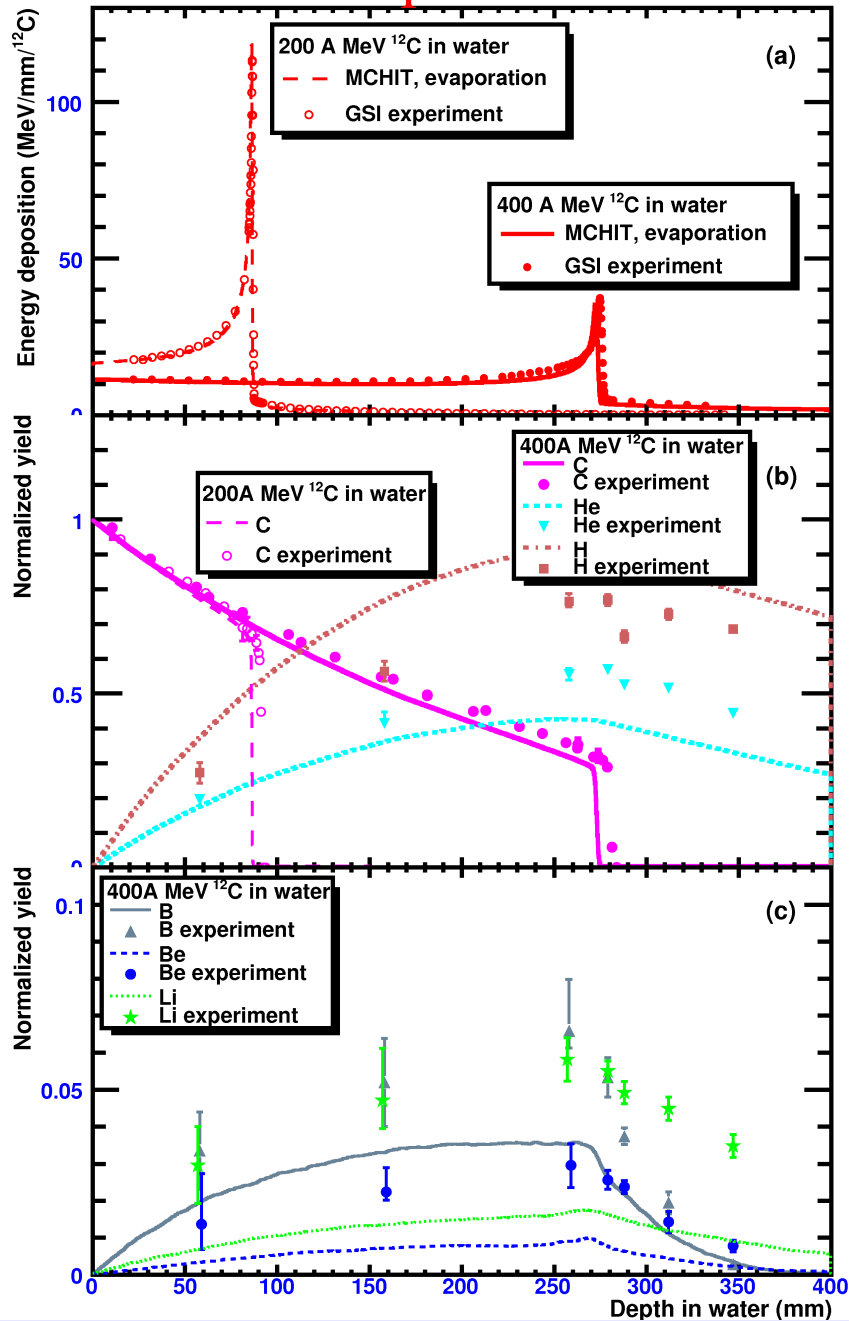
Fragmentation of 400 A MeV ^{12}C beam



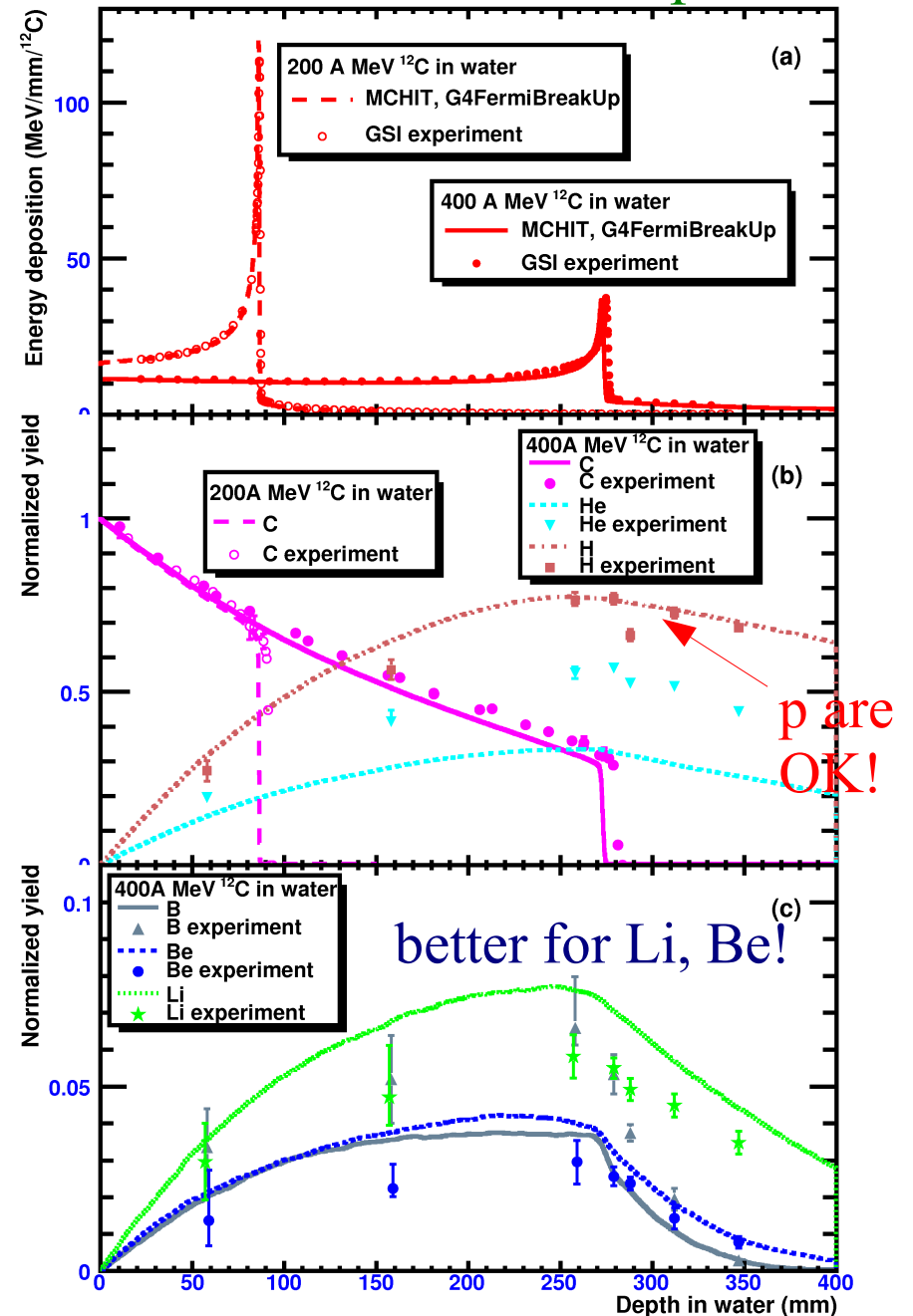
- Up to 70% of beam nuclei are fragmented
- Secondary fragments are created, from protons till Boron with various radiobiological properties
- A lot of work for nuclear fragmentation models!

De-excitations in therapy simulations

Evaporation



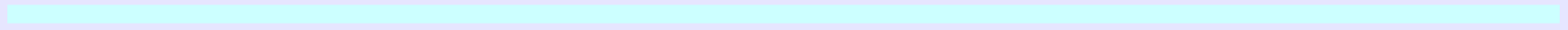
Fermi break-up



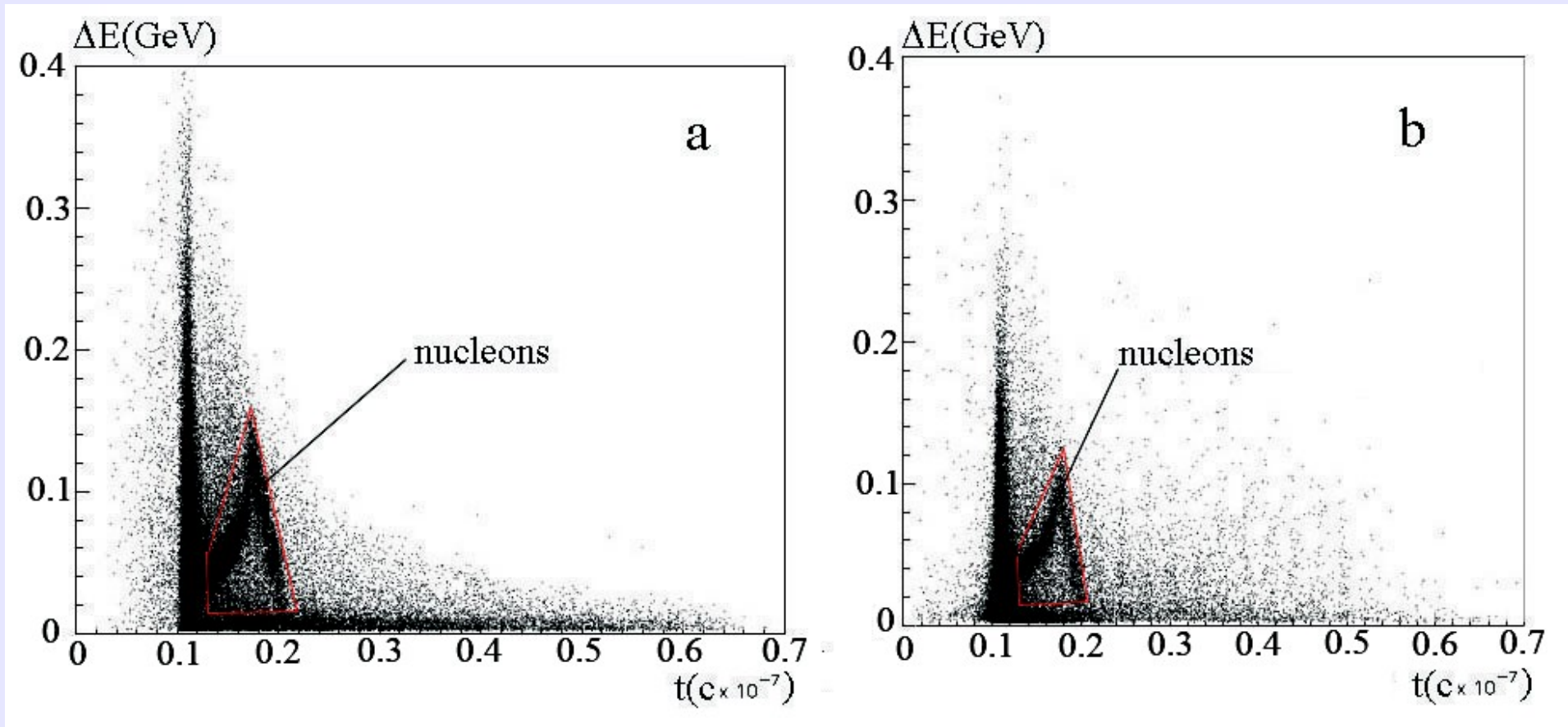
Conclusions

- The RELDIS model predicts a wide distribution of excitation energies of nuclear residues which are created in the photoabsorption of ~ 1 GeV photons by ^{12}C nuclei.
- The most probable photodisintegration events are characterized by emission of 1 or 2 nucleons.
- However, a complete disintegration of ^{12}C into individual nucleons is also seen in a small ($\sim 0.05\%$) fraction of photoabsorption events.
- Isotropic distribution of nucleon emission in high (>7) multiplicity events suggests that they are emitted by a hot thermalized nuclear residue rather than in a cascade process.
- The model describes the fragment multiplicity distributions very well.
- The present study helps to understand nuclear reactions taking place in human tissues during carbon-ion therapy of cancer.

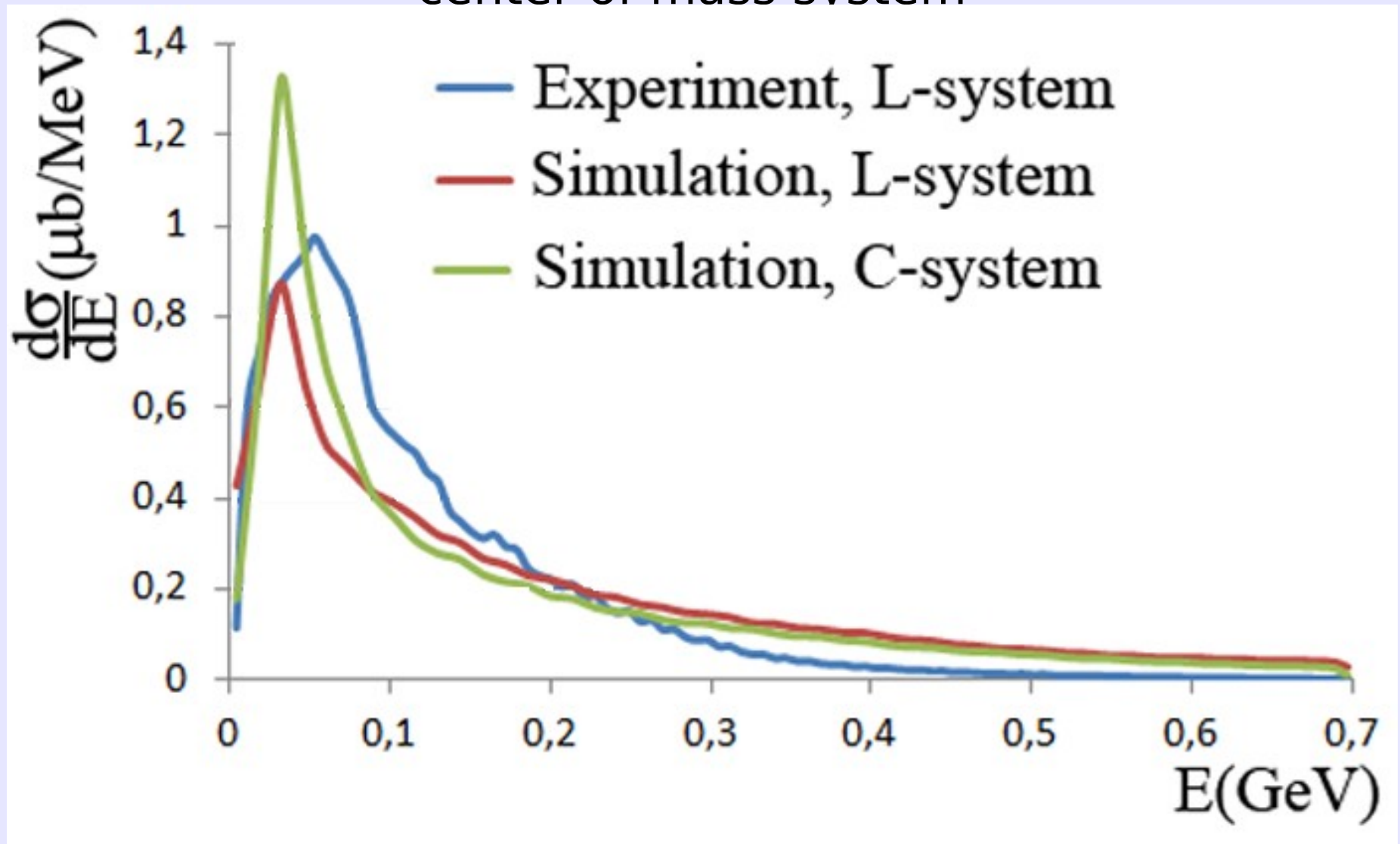
Back-up slides



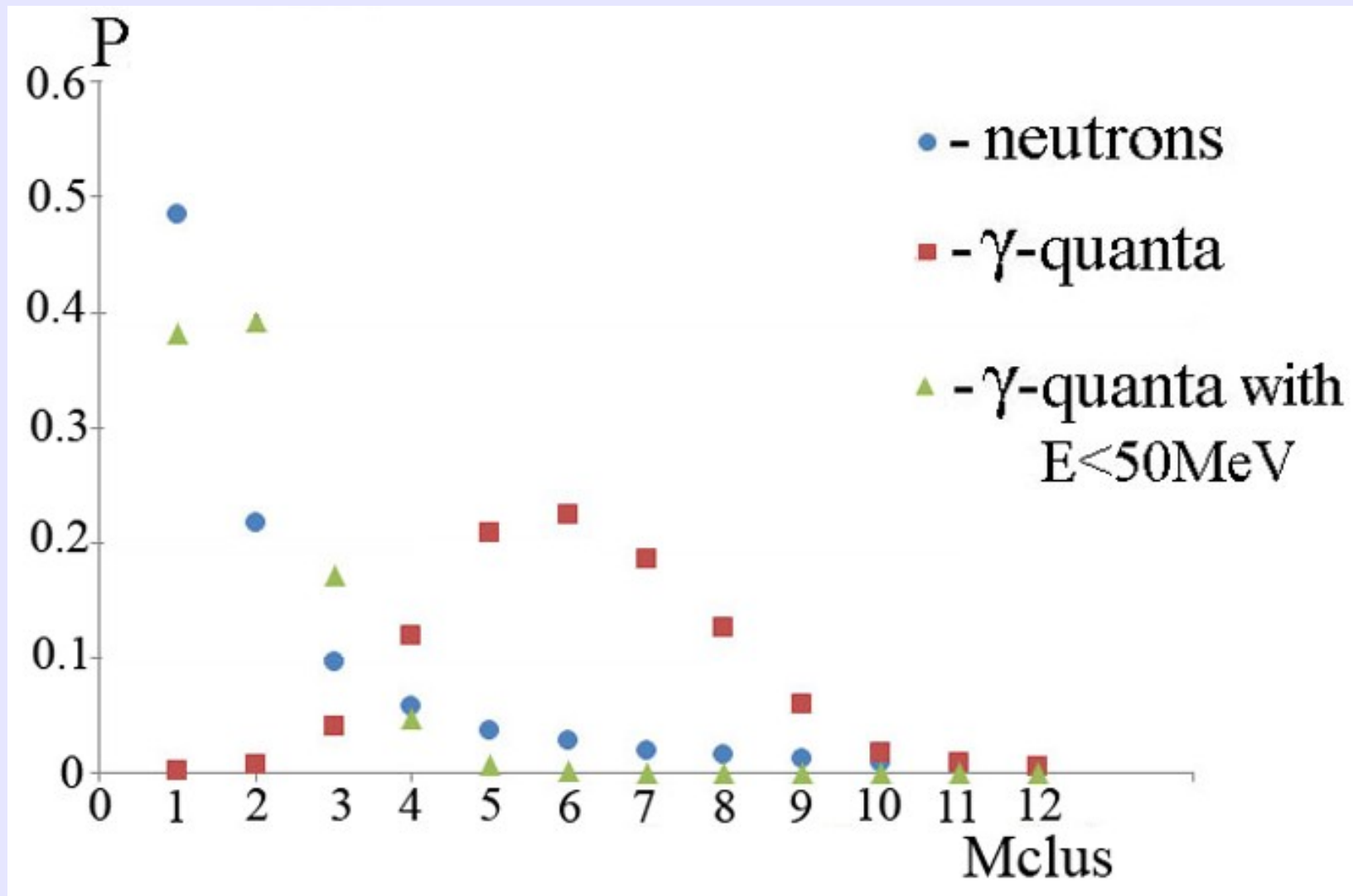
Distinguishing of charged fragments and neutrons from pions and photons in the forward detector. Simulated (a) and detected (b) events are shown.



Energy distributions of nucleons produced in photodisintegration of ^{12}C . Measured (blue line) and calculated (red line) distributions in the laboratory system and calculated distribution (green line) for the center of mass system



Probability to have a given number of fired crystals (cluster size) for neutrons, photons and low energy photons which hit the BGO ball.



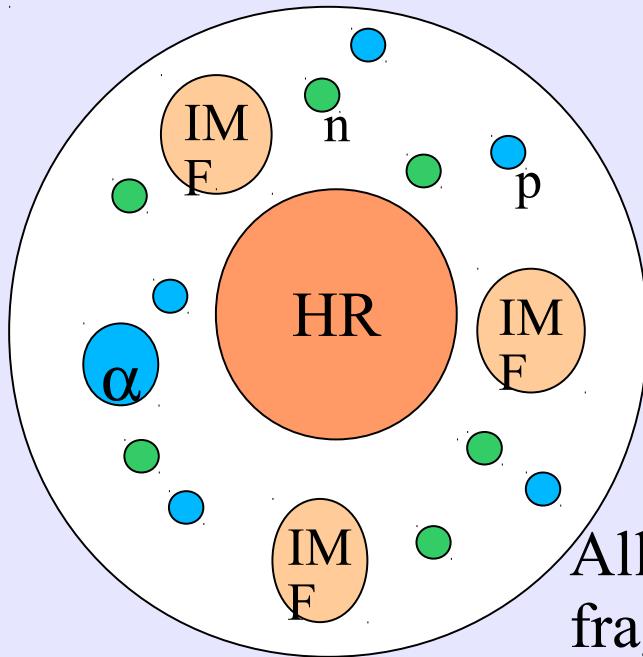
Average numbers of protons and neutrons measured by the BGO ball and forward detectors

	Protons	Neutrons
BGO ball	$2,05 \pm 0,03$	$0,57 \pm 0,01$
Forward direction	$0,35 \pm 0,01$	$0,04 \pm 0,01$

Statistical description of nuclear break-up: SMM

J.P. Bondorf, R. Donangelo, I.N. Mishustin, et al., Nucl. Phys. A443 (1985) 321; A444 (1985) 460;
 J.P. Bondorf, A.S. Botvina, A.S. Iljinov, I.N. Mishustin, K. Sneppen, Phys. Rep. 257 (1995) 133

Ensemble of nucleons and fragments
 in thermal equilibrium characterized by



neutron number N_0
 proton number $Z_0, N_0 + Z_0 = A_0$
 excitation energy $E^* = E_0 - E_{CN}$
 break-up volume $V = (1 + \kappa)V_0$

All break-up channels are enumerated by the sets of
 fragment multiplicities or partitions, $f = \{N_{AZ}\}, M_f = \sum N_{AZ}$

- Baryon number and charge conservation
- in micro-canonical description: $\sum A N_{AZ} = A_0, \sum Z N_{AZ} = Z_0$
 or in macro canonical: $\sum A \langle N_{AZ} \rangle = A_0, \sum Z \langle N_{AZ} \rangle = Z_0$
- Statistical distribution of probabilities: $W_f \sim \exp \{S_f(A_0, Z_0, E^*, V)\}$

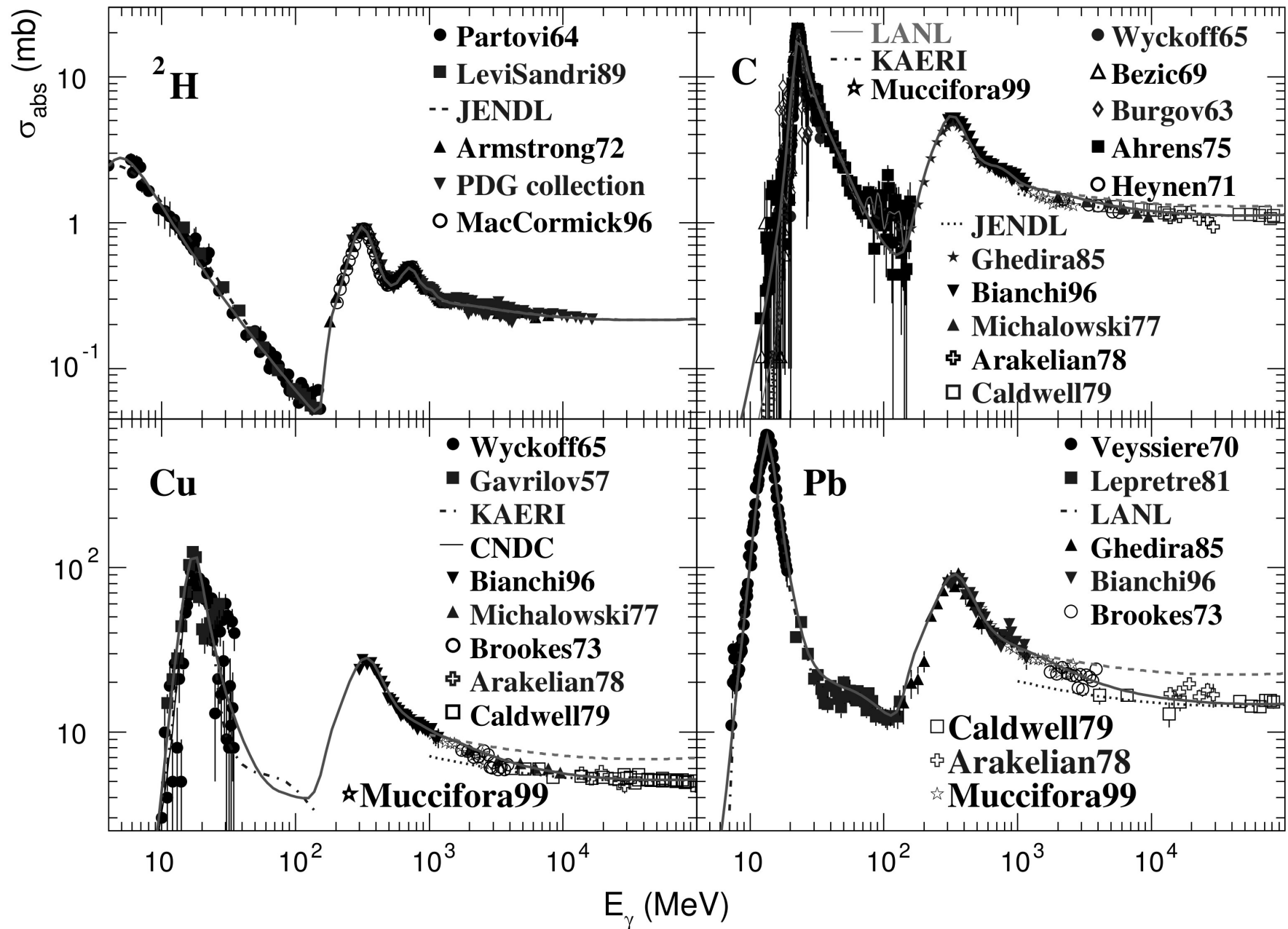
*To find more on the **RELDIS** model:*

A.S. Iljinov et al.,
Nucl. Phys. **A616** (1997) 575

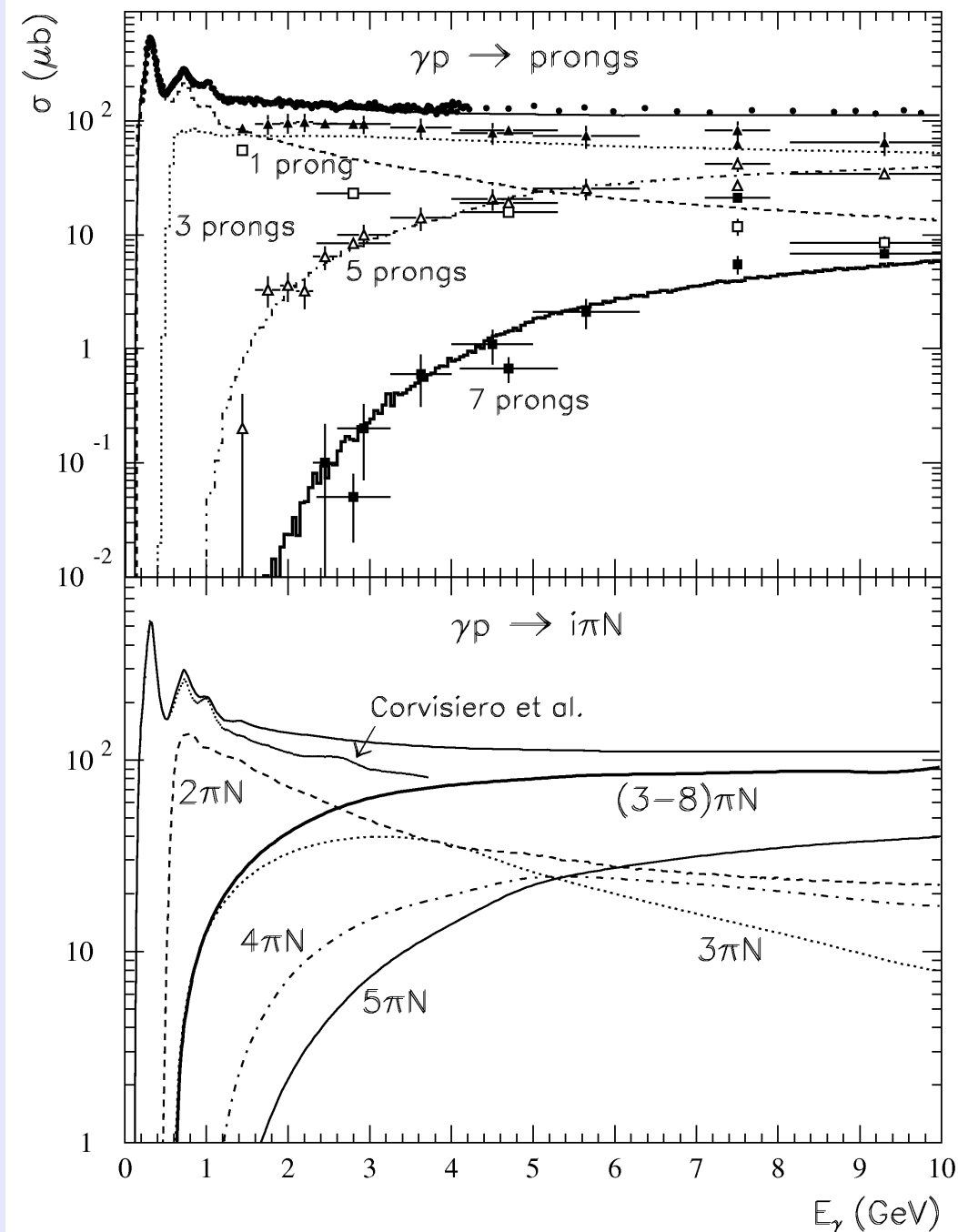
I.A.Pshenichnov et al.,
Phys.Rev. **C57** (1998) 1920;
Phys. Rev. **C60** (1999) 044901;
Phys. Rev. **C64** (2001) 024903

I.A. Pshenichnov,
Phys. Part. Nuclei **42** (2011) 21

Total photoabsorption cross section on nuclei



Meson photoproduction on nucleons



γp	γn
$\gamma p \rightarrow \pi^+ n$	$\gamma n \rightarrow \pi^- p$
$\gamma p \rightarrow \pi^0 p$	$\gamma n \rightarrow \pi^0 n$
$\gamma p \rightarrow \pi^- \Delta^{++}$	$\gamma n \rightarrow \pi^- \Delta^+$
$\gamma p \rightarrow \pi^0 \Delta^+$	$\gamma n \rightarrow \pi^0 \Delta^0$
$\gamma p \rightarrow \pi^+ \Delta^0$	$\gamma n \rightarrow \pi^+ \Delta^-$
$\gamma p \rightarrow \eta p$	$\gamma n \rightarrow \eta n$
$\gamma p \rightarrow \omega p$	$\gamma n \rightarrow \omega n$
$\gamma p \rightarrow \rho^0 p$	$\gamma n \rightarrow \rho^0 n$
$\gamma p \rightarrow \rho^+ n$	$\gamma n \rightarrow \rho^- p$
$\gamma p \rightarrow \pi^+ \pi^- p$	$\gamma n \rightarrow \pi^+ \pi^- n$
$\gamma p \rightarrow \pi^0 \pi^+ n$	$\gamma n \rightarrow \pi^0 \pi^- p$
$\gamma p \rightarrow \pi^0 \pi^0 \pi^0 p$	$\gamma n \rightarrow \pi^0 \pi^0 \pi^0 n$
$\gamma p \rightarrow \pi^+ \pi^- \pi^0 p$	$\gamma n \rightarrow \pi^+ \pi^- \pi^0 n$
$\gamma p \rightarrow \pi^+ \pi^0 \pi^0 n$	$\gamma n \rightarrow \pi^- \pi^0 \pi^0 p$
$\gamma p \rightarrow \pi^+ \pi^+ \pi^- n$	$\gamma n \rightarrow \pi^+ \pi^- \pi^- p$
$\gamma p \rightarrow i\pi N (4 \leq i \leq 8)$	$\gamma n \rightarrow i\pi N (4 \leq i \leq 8)$
(35)	(35)