

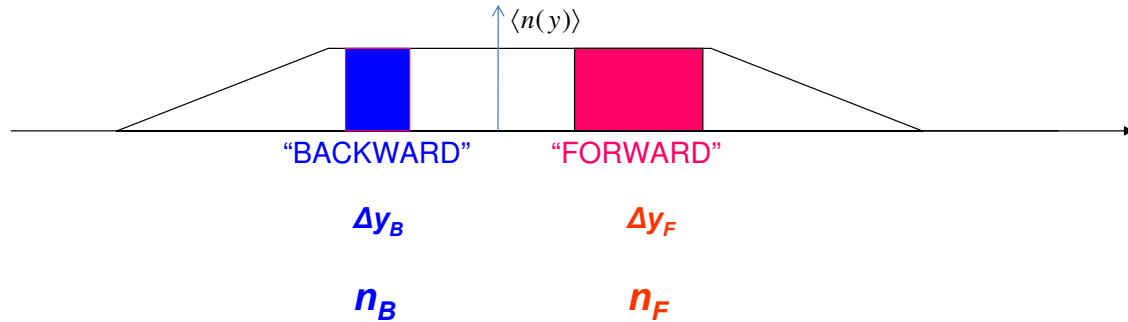
Long-Range Rapidity Correlations in the Model with Independent Emitters

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Long Range Rapidity Correlations
two rapidity intervals separated by a gap

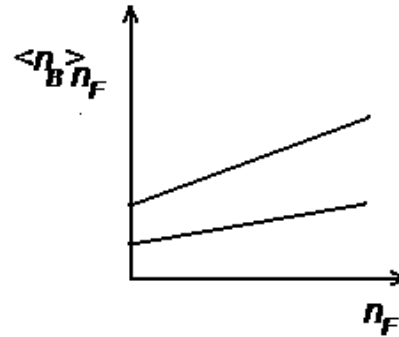
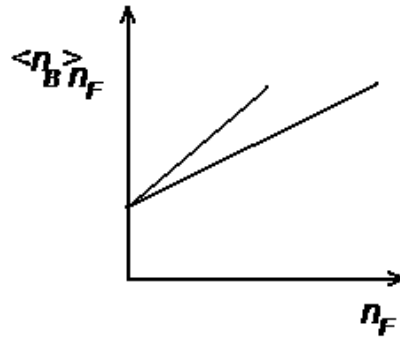


- the event multiplicity in the **BACKWARD** or **FORWARD** rapidity windows.

$\langle n_B \rangle_{n_F} \equiv f(n_F)$ – the correlation function (regression)

The **linear** correlation function (linear regression):

$$\langle n_B \rangle_{n_F} = a^{abs} + b^{abs} n_F$$



$$\frac{\langle n_B \rangle_{n_F}}{\langle n_B \rangle} = a^{rel} + b^{rel} \frac{n_F - \langle n_F \rangle}{\langle n_F \rangle} = a^{rel} + b^{rel} \left(\frac{n_F}{\langle n_F \rangle} - 1 \right)$$

$$b^{rel} = \frac{\langle n_F \rangle}{\langle n_B \rangle} b^{abs}, \quad a^{rel} = \frac{\langle n_B \rangle_{n_F = \langle n_F \rangle}}{\langle n_B \rangle}$$

For a **nonlinear** correlation function $\langle n_B \rangle_{n_F} = f(n_F)$ (nonlinear regression), expanding in powers of $[n_F - \langle n_F \rangle]$ we have

$$\langle n_B \rangle_{n_F} \equiv f(n_F) = f_0 + f_1[n_F - \langle n_F \rangle] + f_2[n_F - \langle n_F \rangle]^2 + f_3[n_F - \langle n_F \rangle]^3 + \dots$$

$$b^{abs} \equiv \left. \frac{d\langle n_B \rangle_{n_F}}{dn_F} \right|_{n_F=\langle n_F \rangle} = f_1, \quad b^{rel} = \left. \frac{d\langle n_B \rangle_{n_F}/\langle n_B \rangle}{dn_F/\langle n_F \rangle} \right|_{n_F=\langle n_F \rangle} = \frac{\langle n_F \rangle}{\langle n_B \rangle} b^{abs}$$

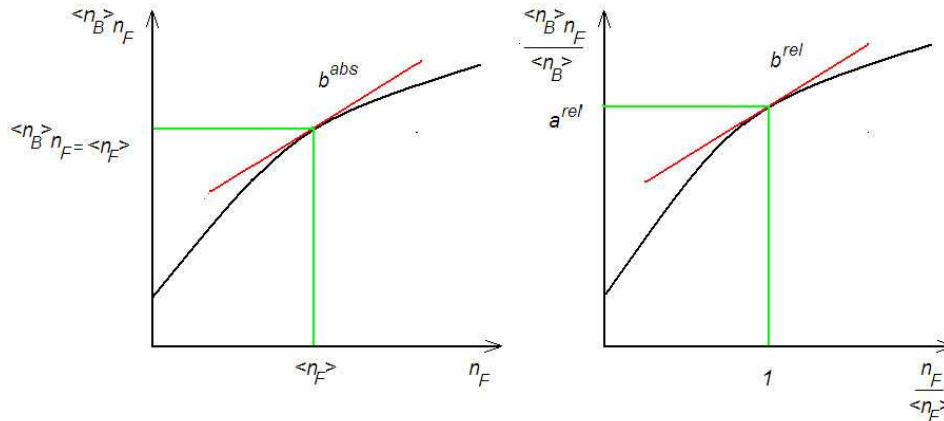
$$\langle n_B \rangle_{n_F=\langle n_F \rangle} = f(\langle n_F \rangle) = f_0, \quad a^{rel} = \frac{\langle n_B \rangle_{n_F=\langle n_F \rangle}}{\langle n_B \rangle}$$

To exclude the trivial dependence on the lengths of the forward Δy_F and backward Δy_B rapidity windows we define the correlation coefficient b^{rel} using the scaled variables:

Definition 1 :

$$b^{rel} \equiv \frac{d\langle n_B \rangle_{n_F} / \langle n_B \rangle}{dn_F / \langle n_F \rangle} \Big|_{n_F = \langle n_F \rangle} = \frac{\langle n_F \rangle}{\langle n_B \rangle} \frac{d\langle n_B \rangle_{n_F}}{dn_F} \Big|_{n_F = \langle n_F \rangle} = \frac{\langle n_F \rangle}{\langle n_B \rangle} b^{abs}$$

where $\langle n_F \rangle$ and $\langle n_B \rangle$ are the mean multiplicities in the forward and backward rapidity windows. The $\langle n_B \rangle_{n_F}$ is the correlation function (regression) - the mean multiplicity in the backward window Δy_B as a function of the multiplicity in the forward window Δy_F .



In the framework of the model with independent emitters in paper [1] using methods developed in [2] under some very specific assumptions the following formula for the defined correlation coefficient b^{rel} was obtained:

$$b^{rel} = \frac{\kappa \bar{\mu}_F}{\kappa \bar{\mu}_F + 1} .$$

Here the κ is the ratio of two scaled variances:

$$\kappa = \frac{V_N}{V_{\mu_F}} , \quad V_N = \frac{D_N}{\langle N \rangle} , \quad V_{\mu_F} = \frac{D_{\mu_F}}{\bar{\mu}_F} ,$$

$\langle N \rangle$ and $D_N = \langle N^2 \rangle - \langle N \rangle^2$ - the mean number of emitters and the event-by-event variance of the number of emitters.

$\bar{\mu}_F$ and $D_{\mu_F} = \overline{\mu_F^2} - \bar{\mu}_F^2$ - the mean multiplicity produced by *one* emitter in the *forward* window and the corresponding variance.

1. **V.V. Vechernin, R.S. Kolevatov, hep-ph/0304295 (2003); Vestnik SPbU, ser.4, no.2, 12 (2004).**
2. **M.A. Braun, C. Pajares and V.V. Vechernin, Phys. Lett. B493, 54 (2000).**

For Poisson distributions $V_N = V_{\mu_F} = 1$ and $\kappa = 1$. Clear that the $\bar{\mu}_F$ is depends on the length of the forward rapidity window. In a first approximation we can assume

$$\bar{\mu}_F = \mu_{0F} \Delta y_F$$

where μ_{0F} is the average multiplicity produced by *one* emitter in the *forward* window per a *unit* of rapidity.

$$b^{rel} = \frac{\kappa \mu_{0F} \Delta y_F}{\kappa \mu_{0F} \Delta y_F + 1} .$$

So the multiplicity correlation coefficient b^{rel} even *defined for scaled variables* nevertheless depends through μ_F on the length of the *forward* rapidity window Δy_F and does not depend on the length of the backward one Δy_B .

This is because the regression procedure is being made by the forward window. One can find the physical discussion of this phenomenon in ref.:

V.V. Vechernin, R.S. Kolevatov, hep-ph/0304295 (2003); Vestnik SPbU, ser.4, no.2, 12 (2004).

For a **linear** correlation function:

$$\langle n_B \rangle_{n_F} = a^{abs} + b^{abs} n_F, \quad \frac{\langle n_B \rangle_{n_F}}{\langle n_B \rangle} = a^{rel} + b^{rel} \left(\frac{n_F}{\langle n_F \rangle} - 1 \right)$$

we have *exactly*:

$$b^{abs} = \frac{\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle}{\langle n_F^2 \rangle - \langle n_F \rangle^2} = \frac{\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle}{D_{n_F}}, \quad b^{rel} = \frac{\langle n_F \rangle}{\langle n_B \rangle} b^{abs},$$

$$a^{rel} = \frac{\langle n_B \rangle_{n_F = \langle n_F \rangle}}{\langle n_B \rangle} = 1$$

So we can take as

Definition 2 :

$$b^{abs} = \frac{\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle}{\langle n_F^2 \rangle - \langle n_F \rangle^2}, \quad b^{rel} = \frac{\langle n_F \rangle}{\langle n_B \rangle} b^{abs},$$

Note that for a **nonlinear** correlation function $\langle n_B \rangle_{n_F} = f(n_F)$

Definition 1 \neq Definition 2

If we use the Definition 2

we can obtain the above formula for b^{rel} at very general assumptions.

Because to calculate such defined correlation coefficient

we *need not to calculate the correlation function* $\langle B \rangle_F = f(F)$,

but only four averages: $\langle F \rangle$, $\langle B \rangle$, $\langle BF \rangle$ and $\langle F^2 \rangle$.

Calculation of the correlation coefficient

Simplified notations:

$$\langle B \rangle_F = a + bF, \quad F \equiv n_F, \quad B \equiv n_B$$

$$P(B, F) - \text{basic}, \quad \sum_{B, F} P(B, F) = 1, \quad \langle BF \rangle \equiv \sum_{B, F} BF P(B, F)$$

$$P(F) = \sum_B P(B, F), \quad \sum_F P(F) = 1, \quad \langle F \rangle \equiv \sum_F FP(F) = \sum_{B, F} FP(B, F)$$

$$P(B) = \sum_F P(B, F), \quad \sum_B P(B) = 1, \quad \langle B \rangle \equiv \sum_B BP(B) = \sum_{B, F} BP(B, F)$$

$$P(B, F) = P(F)P_F(B) \Rightarrow P_F(B) = P(B, F)/P(F) \quad \langle B \rangle_F \equiv \sum_B BP_F(B)$$

For independent identical emitters:

$$P(B, F) = \sum_N w(N) \sum_{B_1, \dots, B_N} \sum_{F_1, \dots, F_N} \delta_{B \ B_1 + \dots + B_N} \delta_{F \ F_1 + \dots + F_N} \prod_{i=1}^N p(B_i, F_i)$$

For LRC:

$$p(B_i, F_i) = p_B(B_i) p_F(F_i)$$

Clear that for identical emitters:

$$\begin{aligned} \sum_{F_i} p_F(F_i) &= 1, & \sum_{F_i} F_i p_F(F_i) &= \bar{\mu}_F, & \sum_{F_i} F_i^2 p_F(F_i) &= \overline{\mu_F^2} \\ \sum_{B_i} p_B(B_i) &= 1, & \sum_{B_i} B_i p_B(B_i) &= \bar{\mu}_B, & \sum_{B_i} B_i^2 p_B(B_i) &= \overline{\mu_B^2} \end{aligned}$$

We denote also

$$\sum_N w(N) = 1, \quad \sum_N N w(N) = \langle N \rangle, \quad \sum_N N^2 w(N) = \langle N^2 \rangle$$

The variances:

$$D_N = \langle N^2 \rangle - \langle N \rangle^2, \quad D_{\mu_F} = \overline{\mu_F^2} - \bar{\mu}_F^2$$

and the scaled variances:

$$V_N = D_N / \langle N \rangle, \quad V_{n_F} = D_{\mu_F} / \bar{\mu}_F$$

Calculation of $\langle n_F^2 \rangle \equiv \langle F^2 \rangle$ as an example

$$\begin{aligned}
 \langle n_F^2 \rangle &\equiv \langle F^2 \rangle \equiv \sum_F F^2 P(F) = \sum_F F^2 \sum_N w(N) \sum_{F_1, \dots, F_N} \delta_{F, F_1 + \dots + F_N} \prod_{i=1}^N p_F(F_i) = \\
 &= \sum_N w(N) \sum_{F_1, \dots, F_N} (F_1 + \dots + F_N)^2 \prod_{i=1}^N p_F(F_i) = \\
 &= \sum_N w(N) \sum_{F_1, \dots, F_N} \left[\sum_{i=1}^N F_i^2 + \sum_{i \neq j=1}^N F_i F_j \right] \prod_{i=1}^N p_F(F_i) = \\
 &= \sum_N w(N) [N \overline{\mu_F^2} + (N^2 - N) \overline{\mu_F^2}] = \langle N \rangle \overline{\mu_F^2} + (\langle N^2 \rangle - \langle N \rangle) \overline{\mu_F^2} = \\
 &= \langle N \rangle (\overline{\mu_F^2} - \overline{\mu_F^2}) + \langle N^2 \rangle \overline{\mu_F^2}
 \end{aligned}$$

So we find

$$\langle n_F^2 \rangle = \langle N \rangle D_{\mu_F} + \langle N^2 \rangle \overline{\mu_F^2}$$

and

$$D_{n_F} \equiv \langle n_F^2 \rangle - \langle n_F \rangle^2 = \langle N \rangle D_{\mu_F} + \langle N^2 \rangle \overline{\mu_F^2} - \langle N \rangle^2 \overline{\mu_F^2} = \langle N \rangle D_{\mu_F} + D_N \overline{\mu_F^2}$$

Gathering we find

$$b^{abs} \equiv \frac{\langle BF \rangle - \langle B \rangle \langle F \rangle}{\langle F^2 \rangle - \langle F \rangle^2} = \frac{\langle BF \rangle - \langle B \rangle \langle F \rangle}{D_{n_F}} = \frac{D_N \bar{\mu}_B \bar{\mu}_F}{\langle N \rangle D_{\mu_F} + D_N \bar{\mu}_F^2}$$

and

$$b^{rel} = \frac{\langle n_F \rangle}{\langle n_B \rangle} b^{abs} = \frac{\langle N \rangle \bar{\mu}_F}{\langle N \rangle \bar{\mu}_B} b^{abs} = \frac{\bar{\mu}_F}{\bar{\mu}_B} b^{abs} = \frac{D_N \bar{\mu}_F^2}{\langle N \rangle D_{\mu_F} + D_N \bar{\mu}_F^2} = \frac{\kappa \bar{\mu}_F}{\kappa \bar{\mu}_F + 1} ,$$

where the κ is the ratio of two scaled variances:

$$\kappa = \frac{V_N}{V_{\mu_F}} , \quad V_N = \frac{D_N}{\langle N \rangle} , \quad V_{\mu_F} = \frac{D_{\mu_F}}{\bar{\mu}_F} ,$$

$$D_N = \langle N^2 \rangle - \langle N \rangle^2 , \quad D_{\mu_F} = \overline{\mu_F^2} - \bar{\mu}_F^2$$

Comparing the definitions

In the case of **nonlinear** regression:

$$\langle n_B \rangle_{n_F} \equiv f(n_F) = f_0 + f_1[n_F - \langle n_F \rangle] + f_2[n_F - \langle n_F \rangle]^2 + f_3[n_F - \langle n_F \rangle]^3 + \dots$$

Definition 1 \neq Definition 2

By **Definition 1** :

$$\bar{b}^{abs} \equiv \left. \frac{d\langle n_B \rangle_{n_F}}{dn_F} \right|_{n_F=\langle n_F \rangle} = f_1, \quad \bar{b}^{rel} = \frac{\langle n_F \rangle}{\langle n_B \rangle} \bar{b}^{abs},$$

By **Definition 2** :

$$b^{abs} = \frac{\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle}{\langle n_F^2 \rangle - \langle n_F \rangle^2} = \frac{\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle}{D_{n_F}}, \quad b^{rel} = \frac{\langle n_F \rangle}{\langle n_B \rangle} b^{abs},$$

$$b^{abs} - \bar{b}^{abs} = \frac{f_2 \langle [n_F - \langle n_F \rangle]^3 \rangle + f_3 \langle [n_F - \langle n_F \rangle]^4 \rangle + \dots}{D_{n_F}}$$

Coefficient a^{rel}

$$a^{rel} = \frac{\langle n_B \rangle_{n_F = \langle n_F \rangle}}{\langle n_B \rangle} = \frac{f(\langle n_F \rangle)}{\langle f(n_F) \rangle}$$

$$\langle n_B \rangle_{n_F} \equiv f(n_F) = f_0 + f_1[n_F - \langle n_F \rangle] + f_2[n_F - \langle n_F \rangle]^2 + f_3[n_F - \langle n_F \rangle]^3 + \dots$$

$$\langle n_B \rangle_{n_F = \langle n_F \rangle} = f(\langle n_F \rangle) = f_0$$

$$\langle n_B \rangle = \langle f(n_F) \rangle = f_0 + f_2 \langle [n_F - \langle n_F \rangle]^2 \rangle + f_3 \langle [n_F - \langle n_F \rangle]^3 \rangle + \dots$$

$$\langle n_B \rangle - \langle n_B \rangle_{n_F = \langle n_F \rangle} = f_2 D_{n_F} + f_3 \langle [n_F - \langle n_F \rangle]^3 \rangle + \dots .$$

So $a^{rel} = 1$ for linear correlation function.

In the next (quadratic) approximation:

If $a^{rel} < 1$ - the correlation function is convex downwards: $f_2 > 0$.

If $a^{rel} > 1$ - the correlation function is convex upwards: $f_2 < 0$.

Conclusions

- ◇ The formula for the long-range multiplicity correlation coefficient in the model with independent emitters is obtained at very general assumptions:

$$b^{rel} = \frac{\kappa \bar{\mu}_F}{\kappa \bar{\mu}_F + 1},$$

where the κ is the ratio of two scaled variances: $\kappa = \frac{V_N}{V_{\mu_F}}$, $V_N = \frac{D_N}{\langle N \rangle}$, $V_{\mu_F} = \frac{D_{\mu_F}}{\bar{\mu}_F}$
and $\bar{\mu}_F$ - the mean multiplicity produced by *one* emitter in the *forward* window.

- ◇ The multiplicity correlation coefficient *defined for scaled variables* nevertheless depends through μ_F on the length of the *forward* rapidity window Δy_F and does not depend on the length of the backward one Δy_B :

$$b^{rel} = \frac{\kappa \mu_{0F} \Delta y_F}{\kappa \mu_{0F} \Delta y_F + 1}, \quad \bar{\mu}_F = \mu_{0F} \Delta y_F$$

where μ_{0F} is the average multiplicity produced by *one* emitter in the *forward* window per a *unit* of rapidity.

- ◇ This is due to the regression procedure is being made by the forward window.

One can find the physical discussion of this phenomenon in ref.:

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