

**FORMATION OF THE 1S_0 DIPROTON IN THE
REACTION $pp \rightarrow \{pp\}_s \pi^0$**

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Motivation

- Bulk of the data with the deuteron $AB \rightarrow dC$ at high Q ($pp \leftrightarrow d\pi^+$, $pn \leftrightarrow d\gamma$ $pd \rightarrow dp, \dots$) in the GeV region. Testing ground for NN models
- Reactions with the 1S_0 diproton $\{pp\}_s$ can give more insight into underlying dynamics due to difference in quantum numbers
 $J_d^\pi = 1^+, T_d = 0 \implies J_{pp}^\pi = 0^+, T_{pp} = 1$

deuteron $\implies (^1S_0)$ pn singlet deuteron or
 $\implies (^1S_0)$ -diproton, $\{pp\}_s$

$pd \rightarrow dp \implies p\{NN\}_s \rightarrow dN$ in $A(p,Nd)B$
suppression of the Δ - and N^* -excitations as 1 : 9 inverse

channel: $pd \rightarrow \{pp\}_s n$

/O.Imambekov , Yu.N.U., In: "Quark-Hadronic Systems"
(Dubna,1987); Yad. Fiz. 52 (1990) 1361/

$$\pi^- \{pp\}_s \rightarrow pn \quad \text{in } {}^3He \quad /M.A. Moinsster, PRL, 1984/$$
$$\gamma \{pp\}_s \rightarrow pp \quad \text{in } {}^3He \quad /J. Laget, NPA, 1989/$$

Δ -contribution is also suppressed as compared to the d channel

- ★ Δ mechanism masks short-range NN properties in the d -channel, and these may reveal themselves in the $\{pp\}_s$ channel.
- ★ Spin structure of the transition m.e. is much simpler for $\{pp\}_s$

Diproton physics at ANKE-COSY, 1999-2010

$$pd \rightarrow \{pp\}_s n \quad 0.5 - 2.0 \text{ GeV}$$

$$pp \rightarrow \{pp\}_s \pi^0$$

$$pp \rightarrow \{pp\}_s \gamma$$

$$pp \rightarrow \{pp\}_s \pi\pi$$

$$pn \rightarrow \{pp\}_s \pi^- \quad (\text{in progress})$$

COSY DATA in the Δ -region

pd \rightarrow {pp}_sn at 0.5-2.0 GeV:

V. Komarov et al., Phys. Lett. B553 (2003) 179;

S.Yaschenko et al., Phys.Rev. Lett. 94 (2005) 072304;

S. Dymov et al., Phys. Rev. C 81 (2010) 044001

pp \rightarrow {pp}_s π^0 at 0.5-2.0 GeV:

S. Dymov et al., Phys. Lett. B 635 (2006) 270.

V. Kurbatov et al., Phys. Lett. B 661 (2008) 22.

pp \rightarrow {pp}_s γ at 0.3-0.8 GeV:

V. Komarov et al., Phys.Rev. Lett. 101 (2008) 102501;

D. Tsirkov et al., J.Phys. G:Nucl.Part.Phys. 37 (2010);

Very simple spin structure of the $pp \rightarrow \{pp\}_s \pi^0$

$$\underline{\frac{1}{2}^+ + \frac{1}{2}^+ \rightarrow \mathbf{0}^+ + \mathbf{0}^-}$$

$$F_{fi}^P = \chi_{\sigma_2}^{(T)}(2) \frac{i\sigma_y}{\sqrt{2}} \left\{ A\sigma \cdot \mathbf{l} + B\sigma \cdot \mathbf{m} \right\} \chi_{\sigma_1}(1), \quad (1)$$

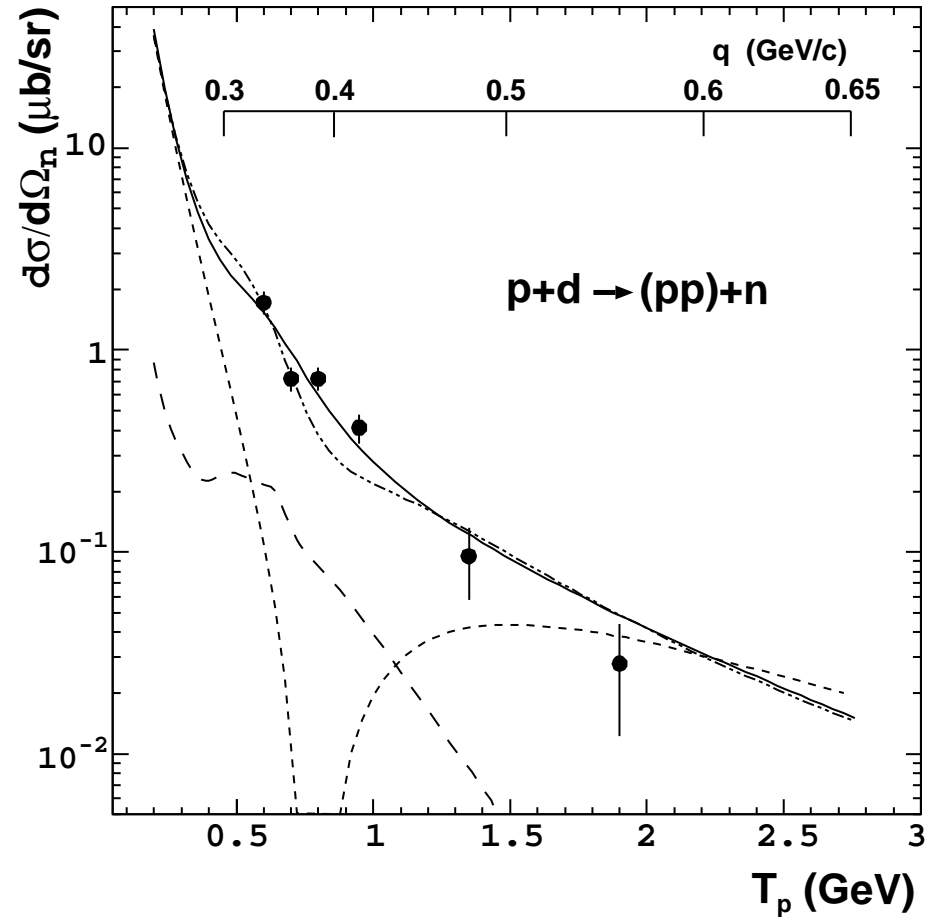
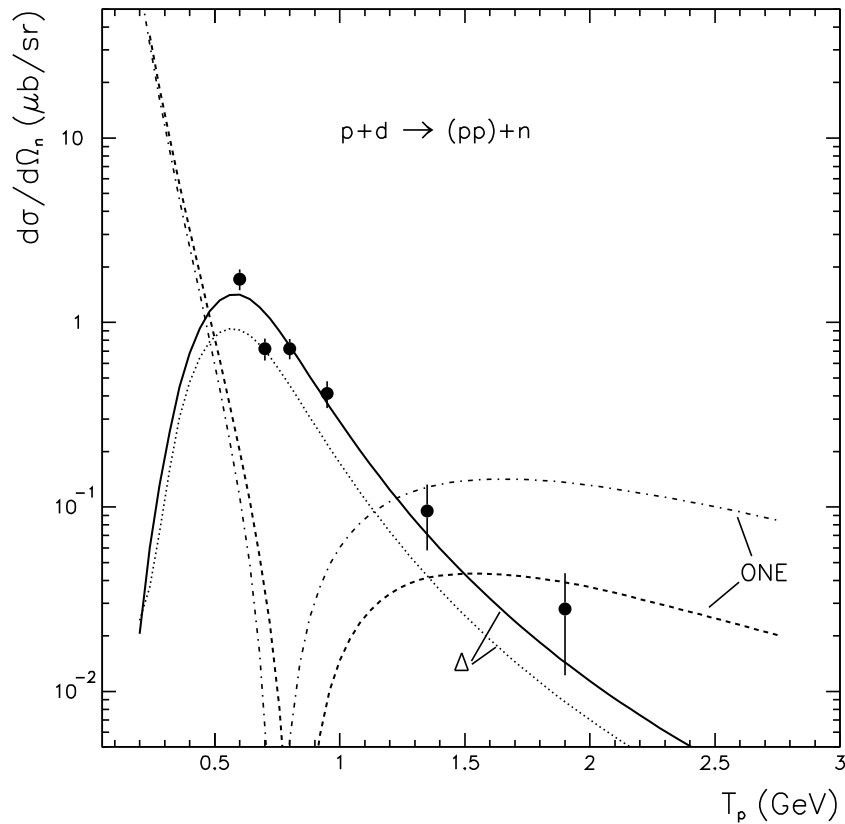
$$d\sigma_0 = \frac{1}{4}(|A|^2 + |B|^2), \quad A_y^b = A_y^t = -\frac{2\text{Im}(A^*B)}{|A|^2 + |B|^2},$$

$$C_{x,x} = -C_{z,z} = \frac{|B|^2 - |A|^2}{|A|^2 + |B|^2}, \quad C_{y,y} = 1,$$

$$C_{x,z} = C_{z,x} = -\frac{2\text{Re}A B^*}{|A|^2 + |B|^2}, \quad (2)$$

COMPLETE POLARIZATION EXPERIMENT:
 $d\sigma_0, A_y, C_{x,z}, (C_{x,x})$. **Contact d-term** $NNNN\pi$ **ChEFT**

$pd \rightarrow (pp)_s n$



ONE+ Δ +SS calculation (*J.Haidenbauer, Yu.Uzikov, Phys.Lett. B562(2003)227*)
When changing hard V_{NN} (RSC, Paris) to the soft V_{NN} (CD Bonn), ONE decreases and Δ -increases providing agreement with the COSY data.
New COSY data confirm this conclusion, PRC 81 (2010)

Allowed transitions in $pp \rightarrow d\pi^+$ & $pp \rightarrow \{pp\}_s\pi^0$

1. $pp \rightarrow d\pi^+$ & $pp \rightarrow \{pp\}_s\pi^0$

1S_0 diproton: $J^\pi = 0^+$, $T = 1$, $S = 0$, $L = 0$

deuteron: $J^\pi = 1^+$, $T = 0$, $S = 1$, $L = 0, 2$

- $(-1)^{L+S+T} = -1$ (Pauli principle)

- Spin-parity conservation:

- ★ $pp \rightarrow \{pp\}_s\pi^0$ L – odd ($L = 1, 3, \dots$) $T = 1$, $S = 1$
 $\implies \Delta N$ in S-wave (or N^*N) $\pi = +1$ - *verboden*

- ★ $pp \rightarrow d\pi^+$ L-odd and even, $T = 1$, $S = 1$ and $S = 0$
 $\implies \Delta N$ in S-wave (N^*N) $\pi = +1$ - *not verboden*

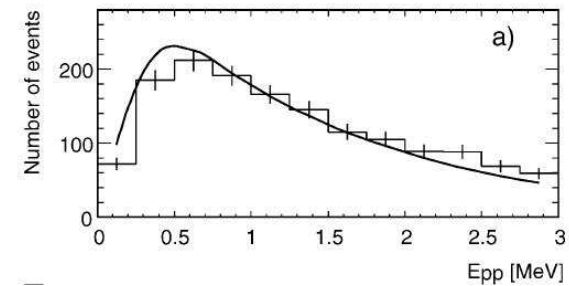
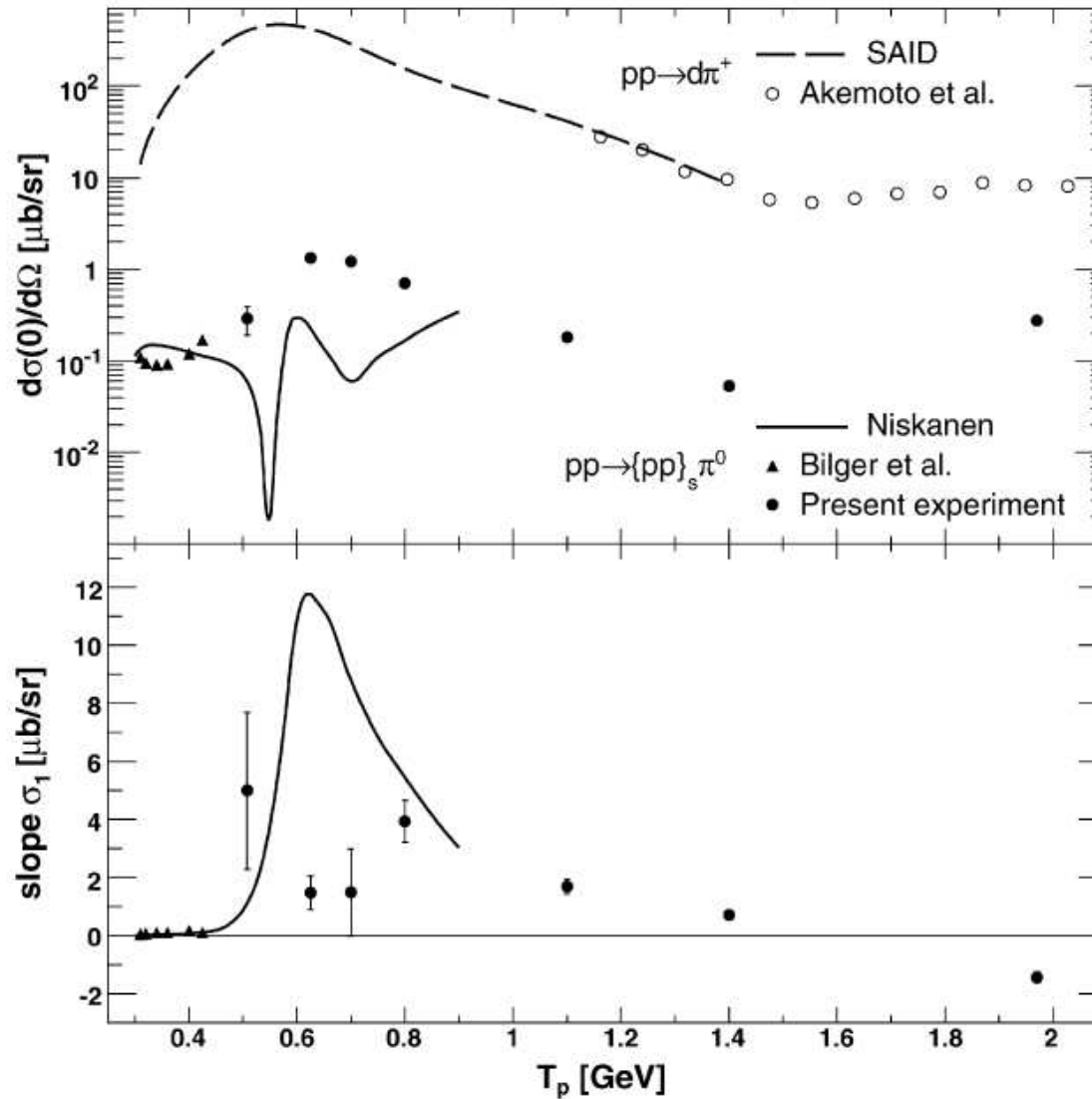
$\Delta(1232)$ dominates in the $pp \rightarrow d\pi^+$ at ≈ 600 MeV

$pp \rightarrow pn\pi^+$ LAMPF data 800 MeV

J.Hudomalj-Gabitzch et al. PRC 18 (1978) 2666

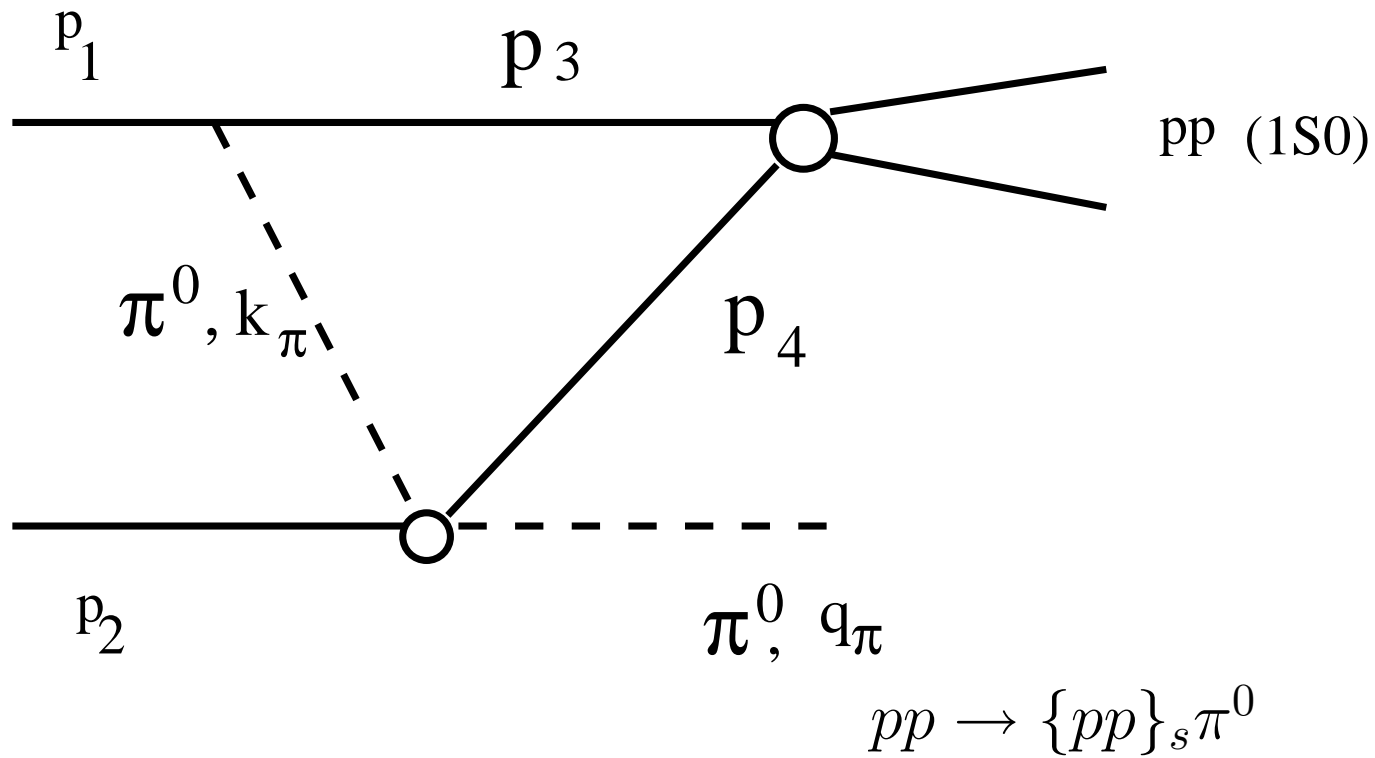
singlet-to-triplet $\xi < \text{few \%}$ in $pp \rightarrow \{pn\}_{s,t}\pi^+$

Yu.N.U, C.Wilkin, PLB 551 (2001)191



$pp \rightarrow \{pp\}_s \pi^0$, theory: J.Niskanen, PLB 642 (2006) 34 /full lines/

The OPE model



The OPE is similar to that for $pd \rightarrow \{pp\}_s n$
/Yu.N.U., J. Haidenbauer, C. Wilkin, PRC **75** (2007) 014008/

$$\begin{aligned} \mathbf{A}^{\text{dir}}(\mathbf{p}_1, \sigma_1, \mathbf{p}_2, \sigma_2) &= \sqrt{3} \frac{f_{\pi\text{NN}}}{m_\pi} \mathbf{N}_{\text{pp}} 2m_p \mathbf{F}_{\pi\text{NN}}(\mathbf{k}_\pi^2) \times \\ &\times \sum_{\sigma_3 \sigma_4 \mu} \left(\frac{1}{2} \sigma_3 \frac{1}{2} \sigma_4 | 00 \right) \left(1 \mu \frac{1}{2} \sigma_3 | \frac{1}{2} \sigma_1 \right) \mathbf{J}^\mu(\tilde{\mathbf{p}}, \gamma) \mathbf{A}_{\sigma_2}^{\sigma_4}(\pi^0 \mathbf{p} \rightarrow \pi^0 \mathbf{p}), \end{aligned} \quad (3)$$

where

$$J^\mu(\mathbf{Q}) = \int \frac{\mathbf{Q}^\mu \Psi_{\mathbf{k}_{\text{pp}}}^{(-)*}(\mathbf{q})}{m_\pi^2 - \mathbf{k}_{\text{pp}}^2 - i\varepsilon} \frac{d^3\mathbf{q}}{(2\pi)^3} \quad (4)$$

$$\mathbf{Q} = \sqrt{\frac{E_1 + m_p}{E_3 + m_p}} \mathbf{p}_3 - \sqrt{\frac{E_3 + m_p}{E_1 + m_p}} \mathbf{p}_1 \quad (5)$$

$$\Psi_k^{(-)*}(\mathbf{q}) = (2\pi)^3 \delta^{(3)}(\mathbf{q} - \mathbf{k}) - \frac{\mathbf{m} \langle \Psi_{\mathbf{k}}^{(-)} | \mathbf{V}(^1\text{S}_0) | \mathbf{q} \rangle}{\mathbf{q}^2 - \mathbf{k}^2 - i\varepsilon} \quad (6)$$

$$J^\mu(\tilde{\mathbf{p}}, \gamma) = \sqrt{\frac{E_1 + m_p}{2m_p} \frac{m_p}{E_1}} \left\{ R^\mu F_0(\tilde{\mathbf{p}}, \gamma) - i\hat{\tilde{\mathbf{p}}}^\mu \Phi_{10}(\tilde{\mathbf{p}}, \gamma) \right\}, \quad (7)$$

where

$$\mathbf{F}_0(\tilde{\mathbf{p}}, \gamma) = \int_0^\infty \mathbf{drr} j_0(\tilde{\mathbf{p}}\mathbf{r}) \psi_{\mathbf{k}}^{(-)*}(\mathbf{r}) \exp(-\gamma\mathbf{r}), \quad (8)$$

$$\Phi_{10}(\tilde{\mathbf{p}}, \gamma) = \mathbf{i} \int_0^\infty \mathbf{dr} j_1(\tilde{\mathbf{p}}\mathbf{r}) \psi_{\mathbf{k}}^{(-)*}(\mathbf{r}) (\mathbf{1} + \gamma\mathbf{r}) \exp(-\gamma\mathbf{r}), \quad (9)$$

$$\psi_{\mathbf{k}}^{(-)}(\mathbf{r}) \rightarrow \frac{\sin(\mathbf{k}\mathbf{r} + \delta)}{\mathbf{k}\mathbf{r}}. \quad (10)$$

$$\gamma^2 = \frac{T_1^2}{(E_1/m_p)^2} + \frac{m_\pi^2}{E_1/m_p}, \quad \mathbf{R} = -\mathbf{p}_1 \frac{m_p T_1}{(E_1 + m_p) E_1}, \quad \tilde{\mathbf{p}} = \frac{\mathbf{p}_1}{E_1/m_p}, \quad (11)$$

where E_1 , \mathbf{p}_1 and $T_1 = E_1 - m_p$ are the total energy, 3-momentum and kinetic energy of the initial proton p_1 ,

$$\frac{d\sigma}{d\Omega} = \frac{1}{(4\pi)^5} \frac{p_f}{s_{pp} p_i} \int_0^{k^{\max}} dk^2 \frac{k}{\sqrt{m_p^2 + k^2}} \frac{1}{2} \int d\Omega_k |\overline{\mathbf{A}_{fi}}|^2, \quad (12)$$

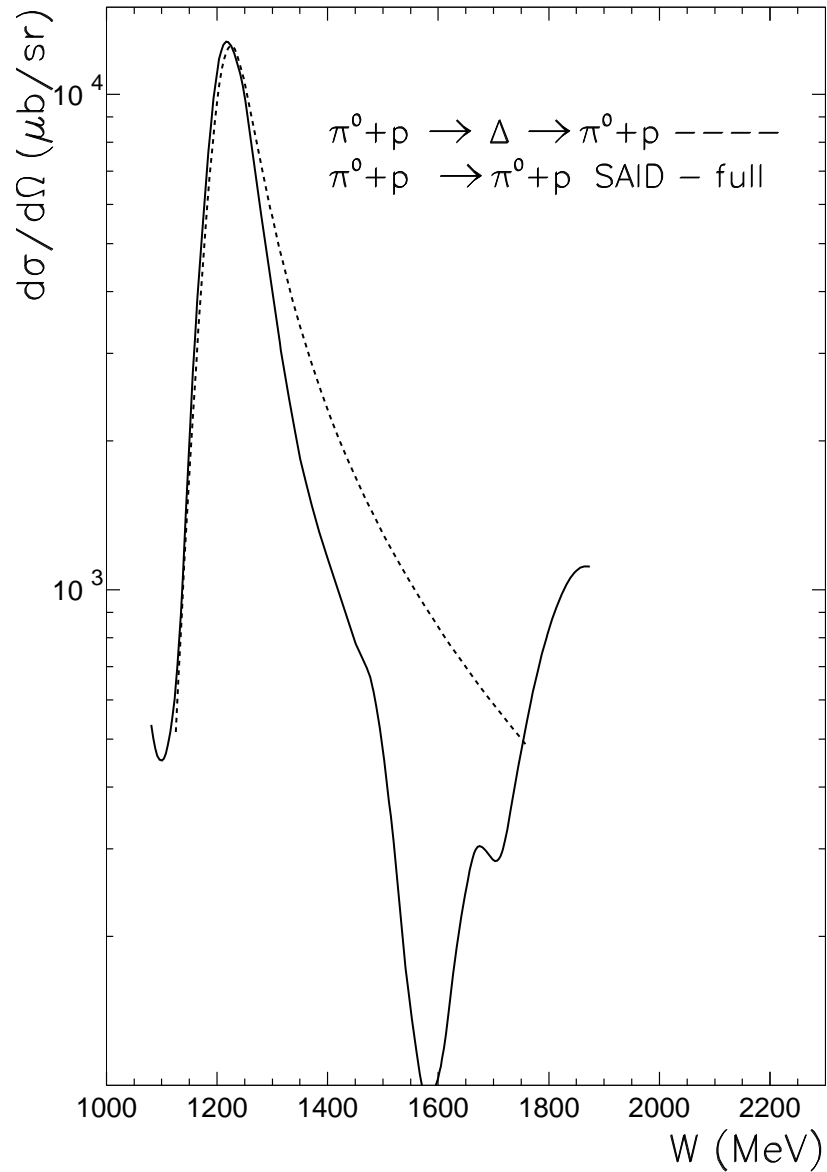
$$\begin{aligned} \frac{d\sigma}{d\Omega_\theta} (pp \rightarrow \{pp\}_s \pi^0) &= \frac{1}{8\pi^2} \frac{p_f}{p_i} \frac{s_{\pi p}}{s_{pp}} \left[\frac{f_{\pi NN}}{m_\pi} N_{pp} m_p F_{\pi NN}(k_\pi^2) \right]^2 \times \\ &\times \int_0^{k^{\max}} dk \frac{2k^2}{\sqrt{m_p^2 + k^2}} \left\{ |J^{\mu=0}(Q)|^2 + 2|J^{\mu=1}(Q)|^2 \right\} \frac{d\sigma}{d\Omega_\phi} (\pi^0 p \rightarrow \pi^0 p) \end{aligned} \quad (13)$$

$$f_{\pi NN}^2/4\pi = 0.0796, \quad F_{\pi NN}(k_\pi^2) = \frac{\Lambda^2 - m_\pi^2}{\Lambda^2 - k_\pi^2}$$

$$k_\pi^2 = 2m_p^2 + p_i p_f \cos \theta - \sqrt{m_p^2 + p_i^2} \sqrt{M_{pp}^2 + p_f^2}, \quad (14)$$

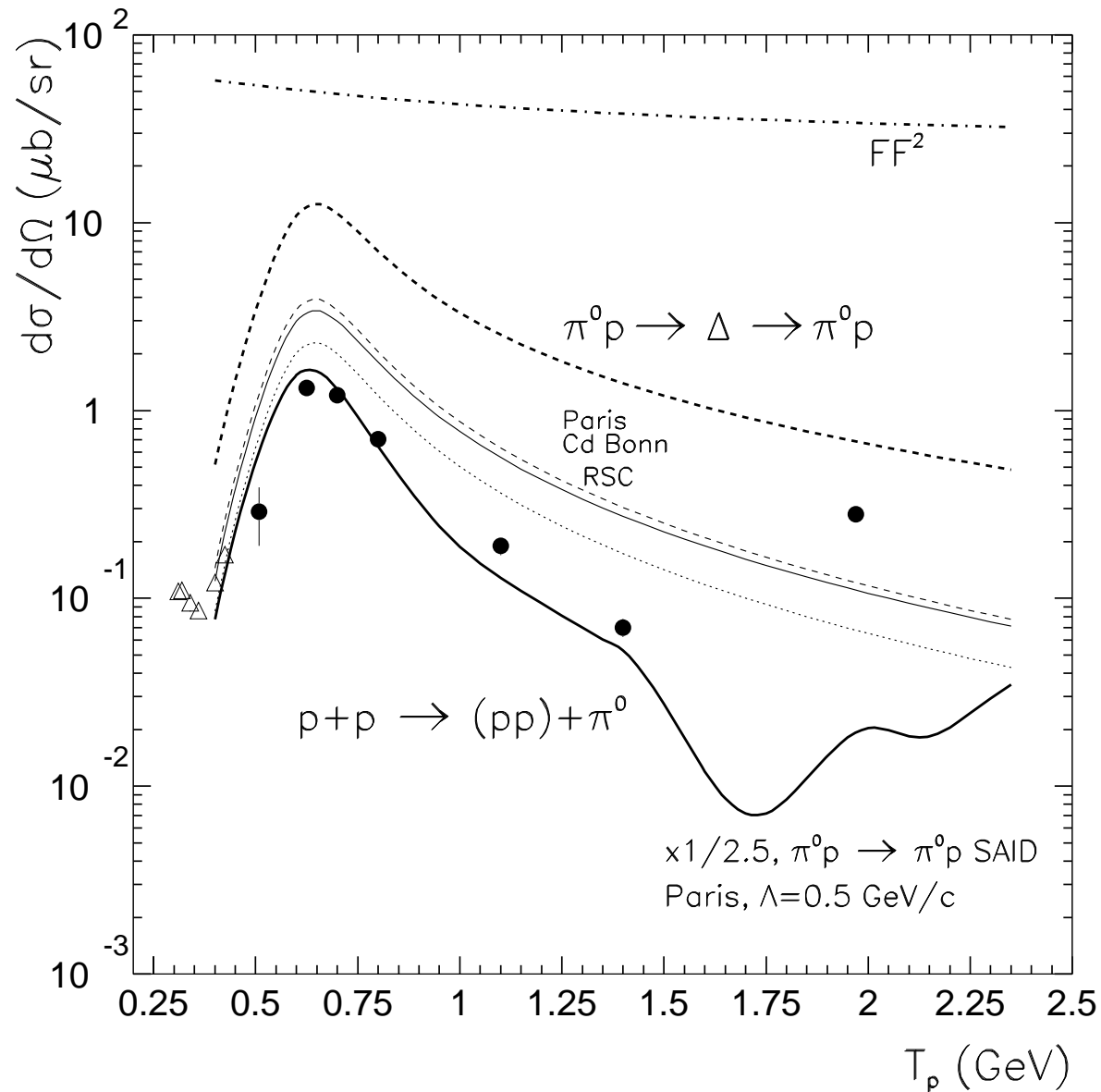
M_{pp} – mass of the diproton. $\Lambda = 0.5$ GeV/c from fit by O.Imambekov, Yu.N.Uzikov Yad.Fiz. (1988) to the $pp \rightarrow pn\pi^+$ LAMPF data at 800 MeV in the Δ -region

$$\pi^0 p \rightarrow \Delta(1232) \rightarrow \pi^0 p$$



$\pi^0 p \rightarrow \pi^0 p$; via Δ mechanism
 (dashed) and from SAID (full)

The OPE results with $\pi^0 p \rightarrow \Delta(1232) \rightarrow \pi^0 p$



$$pp \rightarrow \{pp\}_s \pi^0 ;$$

COSY data: ● – V.Kurbatov et. al PLB 661 (2008) 33

How to exclude Δ from $\pi^0 p \rightarrow \pi^0 p$?

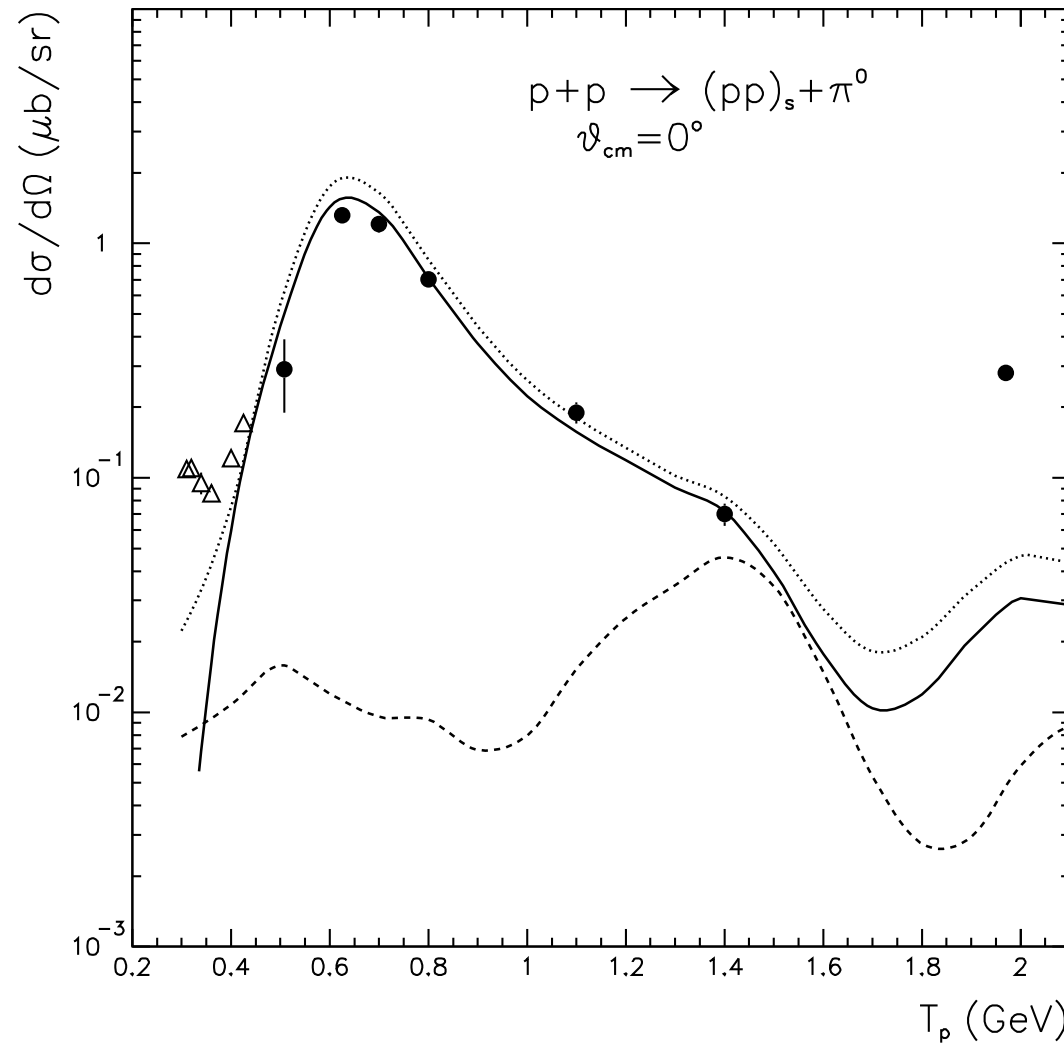
$$\mathbf{A}(\pi^0 \mathbf{p} \rightarrow \pi^0 \mathbf{p}) = \frac{1}{3} \left(\mathbf{a}_{\frac{1}{2}} + 2\mathbf{a}_{\frac{3}{2}} \right), \quad (15)$$

$$\mathbf{d}\sigma(\pi^0 \mathbf{p} \rightarrow \pi^0 \mathbf{p}) = \frac{1}{2} \left\{ \mathbf{d}\sigma(\pi^+ \mathbf{p}) + \mathbf{d}\sigma(\pi^- \mathbf{p}) - \mathbf{d}\sigma(\pi^0 \mathbf{n} \rightarrow \pi^- \mathbf{p}) \right\}, \quad (16)$$

If the amplitude $a_{\frac{3}{2}}$ is excluded from Eq. (15)

$$\mathbf{d}\tilde{\sigma}(\pi^0 \mathbf{p} \rightarrow \pi^0 \mathbf{p}) = \frac{1}{18} \left\{ 3\mathbf{d}\sigma(\pi^- \mathbf{p}) - \mathbf{d}\sigma(\pi^+ \mathbf{p}) + 3\mathbf{d}\sigma(\pi^0 \mathbf{n} \rightarrow \pi^- \mathbf{p}) \right\}. \quad (17)$$

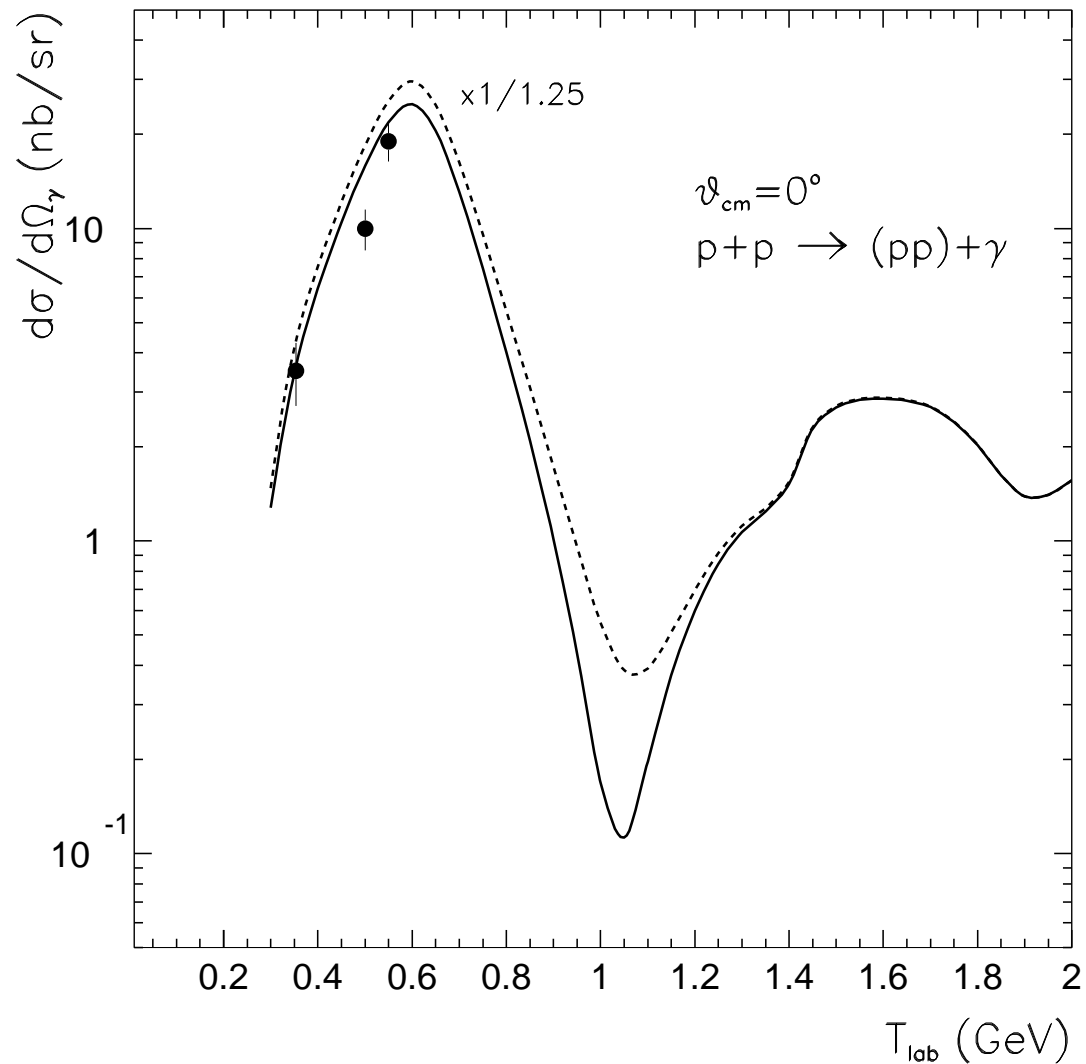
The OPE results with and without $\Delta(1232)$



COSY data: ● – V.Kurbatov et. al PLB 661 (2008) 33

$pp \rightarrow \{pp\}_s \pi^0$;

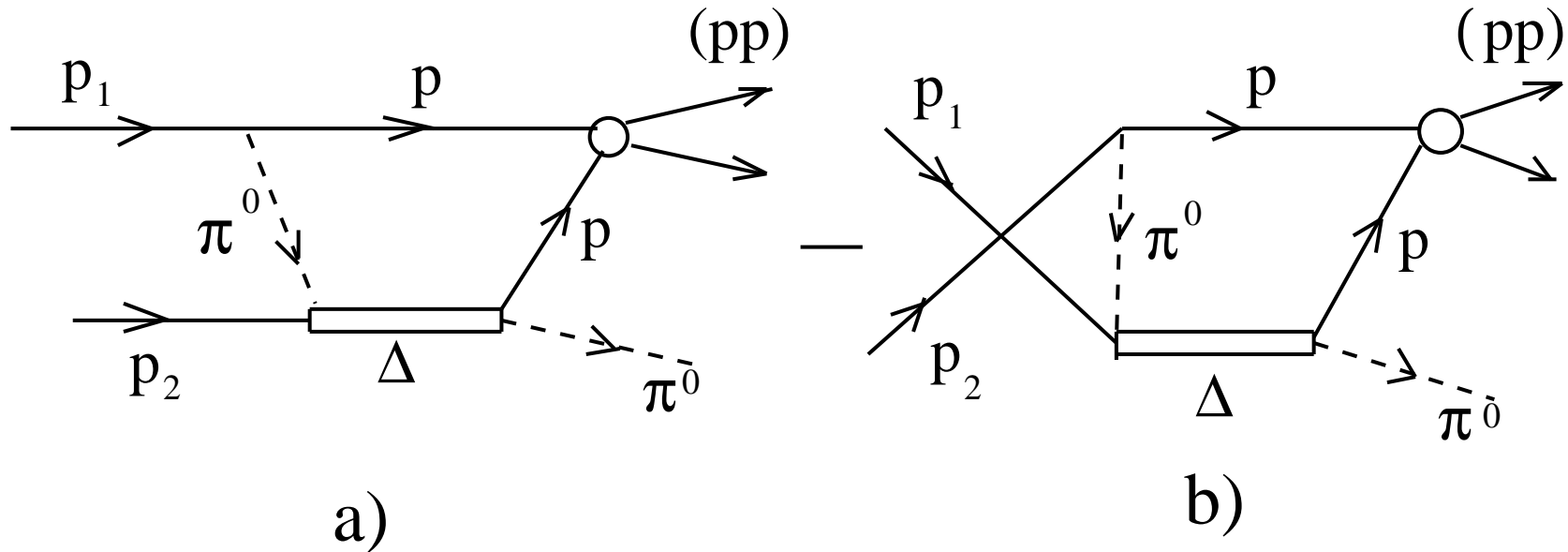
OPE: $pp \rightarrow \{pp\}_s \gamma$



COSY data ● —

V. Komarov et al., PRL (2008)

The BOX-diagramm with Δ for $p\pi^0 \rightarrow p\pi^0$



$$A_{\sigma_1\sigma_2}^{dir} = -8m_\Delta m_p^2 N_{pp} \left(\frac{f_{\pi NN}}{m_\pi} \right) \left(\frac{f_{\pi N\Delta}}{m_\pi} \right)^2 \frac{2}{3} \frac{i}{\sqrt{2}} G_{\sigma_1\sigma_2}^{dir} \times$$

$$\times \int \frac{F_{\pi NN}(k_\pi^2)}{(m_\pi^2 - k_{\pi_a}^2 - i\varepsilon)} \frac{F_{\pi N\Delta}(k_\pi^2)}{(m_\Delta^2 - k_{\Delta_a}^2 - im_\Delta\Gamma)} \frac{\langle \Psi_k^{(-)} | V(^1S_0) | \mathbf{q} \rangle}{(k_{pp}^2 - q^2 + i\varepsilon)} \frac{d^3\vec{q}}{(2\pi)^3} \quad (18)$$

In progress!

Conclusion

- Comparison of d - and $\{pp\}_s$ - channels is very instructive.
- The OPE is an initial step of analysis, explains the shape of $d\sigma/d\Omega(0^\circ)$ for $pp \rightarrow \{pp\}_s\pi^0$ and roughly its absolute value at 0.5-1.5 GeV
- $\Delta(1232)$ contribution is still very important in the $pp \rightarrow \{pp\}_s\pi^0$ (and in $pp \rightarrow \{pp\}_s\gamma$) in spite of strong suppression by spin-parity conservation.
- A similar Δ - dominance was found in $pd \rightarrow \{pp\}_sn$ at 0.5-1 GeV within **ONE+ Δ +SS** & **OPE** models (**softness of NN?**).
- **Outlook: explicite Δ consideration + ONE** \implies
 $\implies A_y, C_{x,x} ? \quad pn \rightarrow \{pp\}_s\pi^-$