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Dense cold matter with special trigger at TWA, Nuclotron, NICA, and FAIR

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for FLINT collaboration @ IBS-XX Dubna, 8 october 2010

Phase diagram of nuclear matter



*p/p₀»1, T/T₀«1(DenseColdMatter): rich structure of the QCD phase diagram - new phenomena are expected!
 **Diagram study not finished-additional new phenomena can be found



*Correlations between signatures could be a signal of the two phase transitions as the density increases and the temperature decreases



Nuclear collisions from 10 to 40 AGeV are the tools to look for the onset of deconfinement (and the critical point?)

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CPOD 2010, Dubna, 27 August 2010

Volker Friese

***Region p/p₀»1, T/T₀«1(DenseColdMatter) hardly
accessible experimentally by standard way

Equation-of-state with and without phase transition



Cassing – Bratkovskaya: Parton-Hadron-String-Dynamics

Perspectives at FAIR/NICA energies partonic energy fraction vs centrality and energy



→Dramatic decrease of partonic phase with decreasing energy and centrality !

****New phase expected to be a small effect – trigger (selection criteria)



*****New phase hardly to be expected as long lived one even for Au+Au collisions

Possible solution for the study of the region $\rho/\rho_0 \gg 1$, $T/T_0 \ll 1$ (DenseColdMatter) : high p_t flucton-flucton interaction





Main idea: to select events with dense matter droplet in the final state with rare cumulative trigger

CUMULATIVE EFFECT



1966 G.A. Leksin: pC →p(137°)X @ 1.0, 6.0 GeV no peaks from pd-, pt-,pHe-... reactions in inclusive spectra. the protons spectra beyond *NN* kinematical limits 1970 A.M. Baldin, V.S.Stavinskiy :

a) Particle production in AA \rightarrow superposition of N+N, N+2N, 2N+N...

b) iN+jN subprocess follow the scaling (the same x-dependence as for N+N int.)

1) X>>1 2) b_{ik}>>1 3) r~r_h A lot of data have been accumulated at JINR , see for example A.M.Baldin et. al.,Sov.J.Part.Nucl.,8(1977)175; JINR comm.E1-8054,Nucl.Phys.A434(1985)695. V.S.Stavinsky, Sov.J.Part.Nucl.,10(1979)949;





For isotopic effects see, also Yad.Fiz. 59 4 694 (1996)

- Some important JINR results: •A-dependence
- •Isotopic effect
- •Spin effects

Cumulative particle production



L. Frankfurt and Strikman Phys. Let. 76B,3 (1978)



M. Braun and V. Vechernin, Nucl. Phys. B **427**, 614 (1994)











aA(flucton) N_{max} ~ 4 Possible solutions : Fragments Lower initial energy not only formally large cumulative number, but also b_{ik}»1(A.M.Baldin)

AA(flucton+flucton): N_{max}~7-8







Fig. 1. Scattering of a virtual photon off a twonucleon correlation, x > 1.5, before (left) and after (right) absorption of the photon.



Fig.ref.:CERN Courier

Fig. 2. Scattering of a virtual photon off a three-nucleon correlation, x > 2, before (left) and after (right) absorption of the photon.





Criterium: r»l

	number of particles	Size(r), fm	free path, (I)fm
Heavy Ions:	1000	10	1
flucton- flucton	10	1	0.1



A.Stavinskiy,8.10.2010,IBS-XX,Dubna

FLINT@ITEP:¹²C + Be $\rightarrow \gamma$ + X

Synchrotron Magnets			
Target 40-70°			
·573	run	2007	2010
	calorimeter	~50 modules	~125 modules
Colorimotor	beam+target	Be+C E _K =3.2 AGeV	Be+C 2,0, 3.2, 4.0 AGeV
Calonineter	intensity	~10 ⁷ N/spill	~2*10 ⁸ N/spill
	trigger	E>1GeV in any glass block	Different triggers
	exposition	1day	20days
	Data	~750K events	~10M events



An estimate of baryon density





FLINT II



Clusterization (close in momentum space) Femtoscopy (close in coordinate space)



. . .

ease)
<mark>lqqqqs)</mark>



Time resolution 0.15ns Neutron efficiency ~ 20%



Neutron detector for FLINT





First prototype:

- Plastic Scintillator 96 * 96 * 128 mm³
- Fiber: KYRARAY,Y-11,d =1mm,
- 4 MRS APD & Amplifier CPTA(Golovin)



Matrix for FLINT 6x6=36 Matrix for NICA ~3000

Neutron detector for FLINT(beam test June 2009,preliminary results)







Rate estimate (preliminary)



Estimated possible data sample (based on ITEP experimental data):

- ITEP:10⁶sec*10⁷int/sec*0.2ster ~5*10³events (CC) for Q1+Q2~5.5
- current experiment (see next report by G.Sharkov for FLINT collaboration)
- Nuclotron-M: 10⁶sec*10⁸int/sec*0.3ster ~10⁴ events (CC) for Q1+Q2~6
- dedicated experimental set-up, large beam intensity and close to optimal initial energy(upper limit of Nuclotron M)
- NICA: 10⁶sec*10⁵int/sec*10ster

 $\sim 3^{*}10^{3}$ events(CC) for Q1+Q2 ~ 5.5

 $\sim 10^4$ events(CC) for Q1+Q2 ~ 6.5

- ~4 π detector (correlations), close to optimal initial energy(lower limit of NICA)
- CBM@SIS100: 10⁶sec*10⁸int/sec*10ster
- large universal detector +beam intensity

The Status of CBM at FAIR



Dense baryon system region for T/A =6.0 GeV



DCM @ NICA



Detector \ Task	ECAL (trigger)	Neutron detector	TOF	Magnet	Vertex detector	Tracking detector
Trigger	•					
Cluster	•	•	•	•		•
Femtoscopy	•	•	•	•		•
Isosymmetri- sation	•	•	•	•	•	•
Strangeness	•	•	•	•	•	•
Vector mesons	•		•	•	•	•
Exotics	•	•	•	•	•	•
Multi-bosons/ fermions	•	•	•	•	•	•

Conclusion



- 1. New phase diagram sector proposed for the study. It potentially create new wide experimental program.
- 2. Possible trigger is proposed; its efficiency is under experimental test
- 3. A lot of work must to be done-cooperation are welcome.









A rendition of the structure and phases of a neutron star (courtesy of Dany Page)

nucl-th/0901.4475

The discovery of neutron stars in the form of pulsars has been a major stimulus to dense matter studies Observables:

gravitational red shift central density of the star moments of inertia pulsar timing

These informations can be inferred from the photons, randing from the radio waves to X-rays, and also those involving neutrinos and gravity waves

[T.K.Jha,nucl-th/0902.0262]

Outlook



- Physical motivation
- Program status @ FLINT
- Future with FAIR-CBM, NICA-MPD
- Conclusions



Arguments for FF: slope parameter(see figure) (also min (Ni+Nj)&angular dependence)





RESULTS(2008)





How many constituents needed to say about matter?







1000-100-10? It depends on density!



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FLINT SUBSYSTEMS





Celorimeters & LED-monitoring system







VETO-system for FLINT





Plastic Scintillator 105 * 105 * 5 mm^3

Fiber: KYRARAY,Y-11,d =1mm, wavelength shift

MRS APD & Amplifier - CPTA(Golovin)





VETO system supermodule



Electonics(VME)









Figure 14: Dependence of symmetry energy of nuclear matter on density. CT[1] and CT[2] refer to the results of refs. [112] and [113] obtained within the covariant T-matrix theory. V refers to the variational result of ref. [114] which includes a three-body force; BBG refers to the results based on the BBG theory with two-body forces [115]; BBG(1) includes in addition a microscopically derived three-body force [116, 117, 118]; BBG(2) is the same as BBG(1) but with a phenomenological three-body force [119, 120].

The physics of dense ha	dronic matter and compact stars	[115] X. R. Zhou, G. F. Burgio, U. Lombardo, HJ. Schulze, and W. Zuo, Phys. Rev. C 69 (2004) 018801.
		[116] S. A. Coon et al., Nucl. Phys. A 317 (1979) 242.
Arme	en Sedrakian	[117] P. Grangé, A. Lejeune, M. Martzolff, and JF. Mathiot, Phys. Rev. C 40 (1989) 1040.
[112] H. Müther, M. Prakash, and T. L. Ainswe	orth, Phys. Lett. B 199 (1987) 469.	[118] W. Zuo, A. Lejeune, U. Lombardo, and JF. Mathiot, Eur. Phys. J. A 14 (2002) 469; Nucl. Phys.
[113] CH. Lee, T. T. S. Kuo, G. Q. Li, and G.	. E. Brown, Phys. Rev. C 57 (1998) 3488.	A 706 (2002) 418.
[114] A. Akmal, V. R. Pandharipande, and D.	G. Ravenhall, Phys. Rev. C 58 1804 (1998).	[119] B. S. Pudliner, V. R. Pandharipande, J. Carlson, and R. B. Wiringa, Phys. Rev. Lett. 74 (1995) 4396.
08.10.2010	A Stavinskiy 8 10 20	1201 B. S. Pudliner, V. R. Pandharipande, J. Carlson, S. C. Pieper, and R. B. Wiringa, Phys. Rev.

Prog.Part.Nucl. Phys.58(2007)158 nucl-th/0601086

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At densities around the saturation density the only baryonic degrees of freedom are protons and neutrons, which form an iso-duplet whose approximate free-space SU(2) symmetry is largely broken in matter. At larger densities the number of stable baryons increases. These include the isospin 3/2 nucleon resonances Δ^{\pm} , Δ^{0} and the strangeness carrying baryons (hyperons). The hyperonic states can be

perons and nucleons, the threshold for hyperons to become stable is determined by comparison of the hyperon mass to the largest available energy scale - the neutron Fermi energy [121]. The Σ^- hyperons can appear in matter through the weak hyperonic (inverse) beta-decay reactions $e^- + n \rightarrow \Sigma^- + \nu_e$ and hadronic weak decay $n + \pi^- \rightarrow \Sigma^-$. The energy balance in the first reaction implies $2\mu_n \simeq M_{\Sigma^-} = 1197$ MeV, where μ_n is the chemical potentials of neutrons (we used the fact that the chemical potential of neutrons and electrons are almost equal in matter under β equilibrium, see subsection 2.7). The r. h. side of the second reaction is $O(\mu_n)$, therefore it is negligible compared to the first reaction. Similar arguments apply to other hyperons which are stabilized either through the hyperonic β decays or hadronic weak decays. For example for the lightest hyperon Λ^0 one finds

$$p + e^{-} \rightarrow \Lambda^{0} + \nu_{e}, \quad p + \pi^{-} \rightarrow \Lambda^{0}, \quad O(\mu_{p}) \\ n + \pi^{0} \rightarrow \Lambda^{0} \qquad O(\mu_{n})$$

$$= M_{\Lambda} = 1116 \text{ MeV}$$

$$(124)$$

The reactions in the first line being $O(\mu_p)$, where $\mu_p \ll \mu_n$ is the proton chemical potential, can be neglected and Λ^0 appear primarily through the weak hadronic process in the second line. Since the r. h. side of this reaction is $O(\mu_n)$ and the mass difference $M_{\Sigma^-} - M_{\Lambda} < \mu_n$ at relevant densities, $\Sigma^$ hyperons appear first.

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Angular dependence







linearity & energy resolution



TAPS ${}^{12}C+{}^{12}C\rightarrow\pi^{\circ}(\eta)X$ @ 0.8, 1.0 & 2.0 AGEV







•Постановка задачи реалистична •Необходимы бо́льшие iN+jN при бо́льших Е₀

Кумулятивное число



$$\begin{cases} E_0 + M = E + E_X \\ p_0 = p \cdot \cos \theta + p_X \cdot \cos X \\ 0 = p \cdot \sin \theta + p_X \cdot \sin X \\ m_X^2 = (m_0 + M)^2 \end{cases} \implies \begin{cases} E_X^2 = E_0^2 + M^2 + E^2 + 2E_0 M - 2EM - 2E_0 E \\ p_X^2 = E_0^2 + p^2 - 2p_0 p \cos \theta \\ m_X^2 = (m_0 + M)^2 \end{cases}$$

$$X \cdot m_N = M = \frac{E_0 E - p_0 p \cos \theta - m^2/2}{T_0 - E}$$

1+N	4,0	1+(N+1)	5,0
2+N	1,9	2+(N+2)	3,9
3+N	1,6	3+(N+3)	4,6

спектр в гипотезе флуктонфлуктонного взаимодействия



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Fermi motion



•разыгрываются π°, в направлении аксептанса

 ${}^{\bullet}\pi^{\circ} \to \gamma\gamma$

•в жёсткой части наклон γ уменьшается ightarrow ассиметричные распады π°



Is it possible to create dense cold matter in lab? At large scale - probably not (left figure) Al relatively small scale – probably yes (right one)



Next steps



- Neutron detector beam test
 - Time resolution

Nuclotron M, MPD-NICA

- Neutron detector beam test
 - Efficiency study
- $AA \rightarrow \gamma + X$
 - as high Q₁+Q₂ as possible with FLINT@Nuclotron-M
- $AA \rightarrow \gamma + B + X$
 - dense baryon matter study @ MPD-NICA

Phase diagram of nuclear matter



