

THE CASE FOR THE DIBARYON

d_1^* (1956)

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OUTLINE

- Introduction
- The evidence for the $d_1^*(1956)$ in processes with real photons.
- A new mechanism for dielectron production in NN collisions due to the $d_1^*(1956)$ resonance.
- Contribution of the $pp \rightarrow \gamma^* d_1^* \rightarrow e^+ e^- d_1^*$ mechanism to the $pp \rightarrow e^+ e^- X$ *DLS* data.
- Contribution of the $NN \rightarrow e^+ e^- d_1^*$ mechanism to the *DLS* $AA \rightarrow e^+ e^- X$ data.
- Conclusion

INTRODUCTION

Searches for exotic multiquark states and in particular six-quark states with the baryon number $B=2$ (dibaryons) has a long history.

The results of these searches are believed to be unsuccessful.

The aim of this talk to show you that in the reality this is rather not the case.

There are reasons to believe that one of the exotic states still was observed.
This is a six-quark state that we called .

FIRST EVIDENCE FOR THE d_1^* (1956)

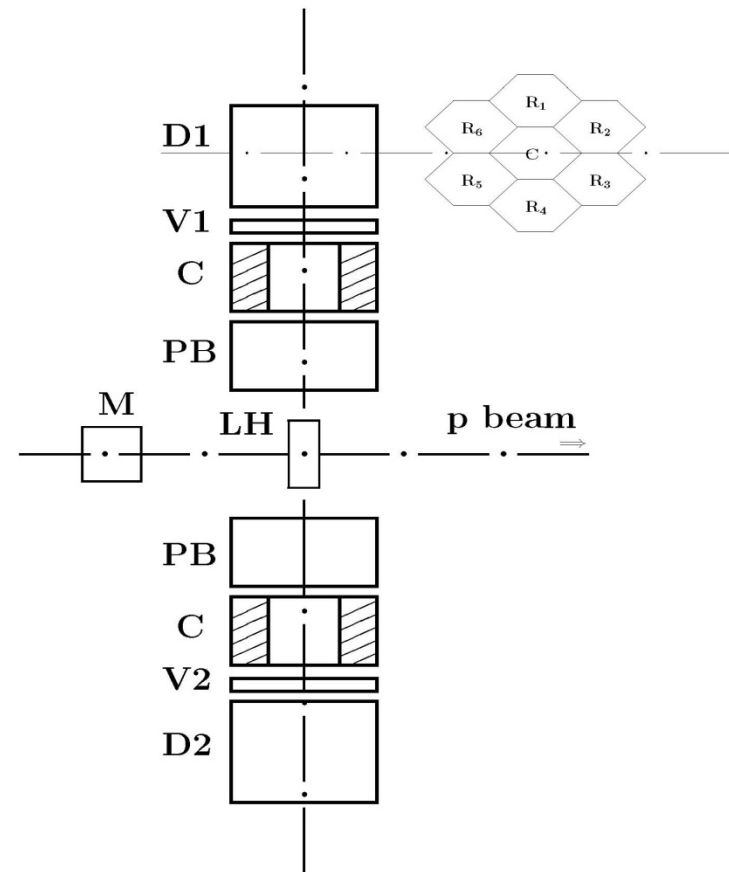
A.S. Khrykin et al., Phys.Rev. C64, 034002(2001)

Reaction : $pp \rightarrow \gamma\gamma X$

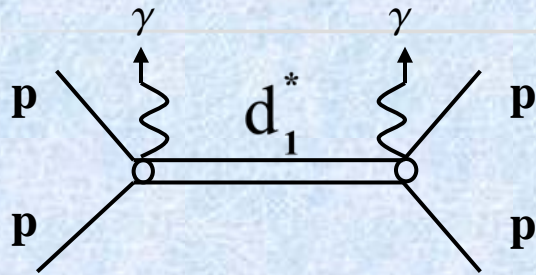
$T_p = 216 \text{ MeV}$; $\theta_\gamma = \pm 90^\circ$

Two photons were detected in coincidence.

Photon energy spectrum was measured.



$$pp \rightarrow \gamma d_1^* \rightarrow \gamma\gamma pp$$



$$E_{\gamma}^F \text{ (cms)} = \frac{S - M_R^2}{2\sqrt{S}}$$

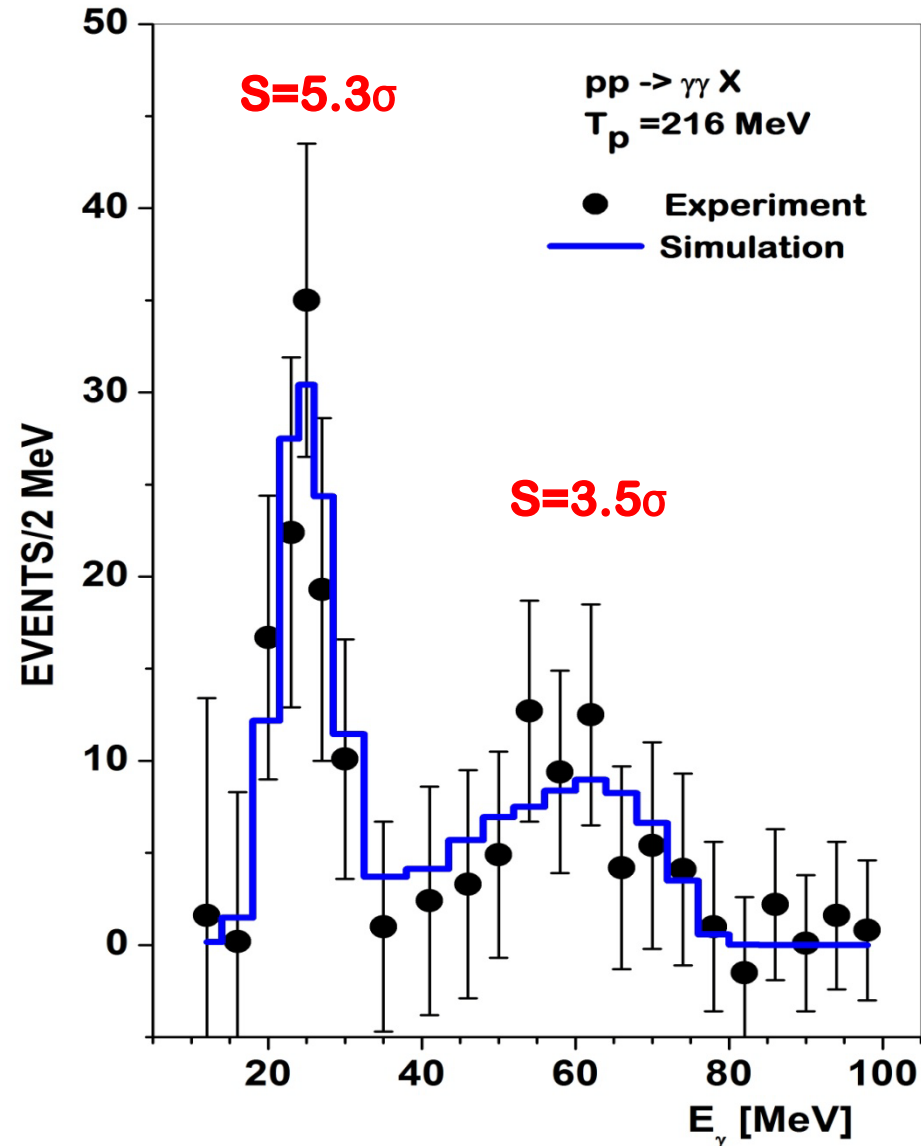
$$E_{\gamma}^D \text{ (rf)} = \frac{M_R^2 - M_{pp}^2}{2M_R}$$

The $d_1^*(1956)$ existence implies the existence of a new mechanism of photon pair production that should be taken into account.

$$M_{\gamma\gamma}^2 = 2 * E_{\gamma}^F * E_{\gamma}^D (1 - \cos(\theta_{12}))$$

$$M_R \sim 1956 \text{ MeV}; \Gamma \sim 8 \text{ MeV}$$

Total $S = 6.3 \sigma$



**The evidence for the d_1^* (1956) in the processes
with real photons.**

Exclusive $pp \rightarrow pp\gamma\gamma$ reaction

A.S.Khrykin and S.B.Gerasimov, in : *Proc. of the MENU2007*, IKP, Forschungszentrum Juelich, Germany, September 10-14, 2007, edited by H. Machner and S. Krewald, eConf C070910(2008),250.

Experiment

CELSIUS-WASA Collaboration

Bashkanov et al. *Int. Jour. of Mod.*

Phys. A20,554(2005); hep-ex/0406081

$pp \rightarrow pp\gamma\gamma$

The two-photon invariant mass spectra of the reaction were measured at two energies:

$T_p=1.36$ GeV and $T_p=1.2$ GeV

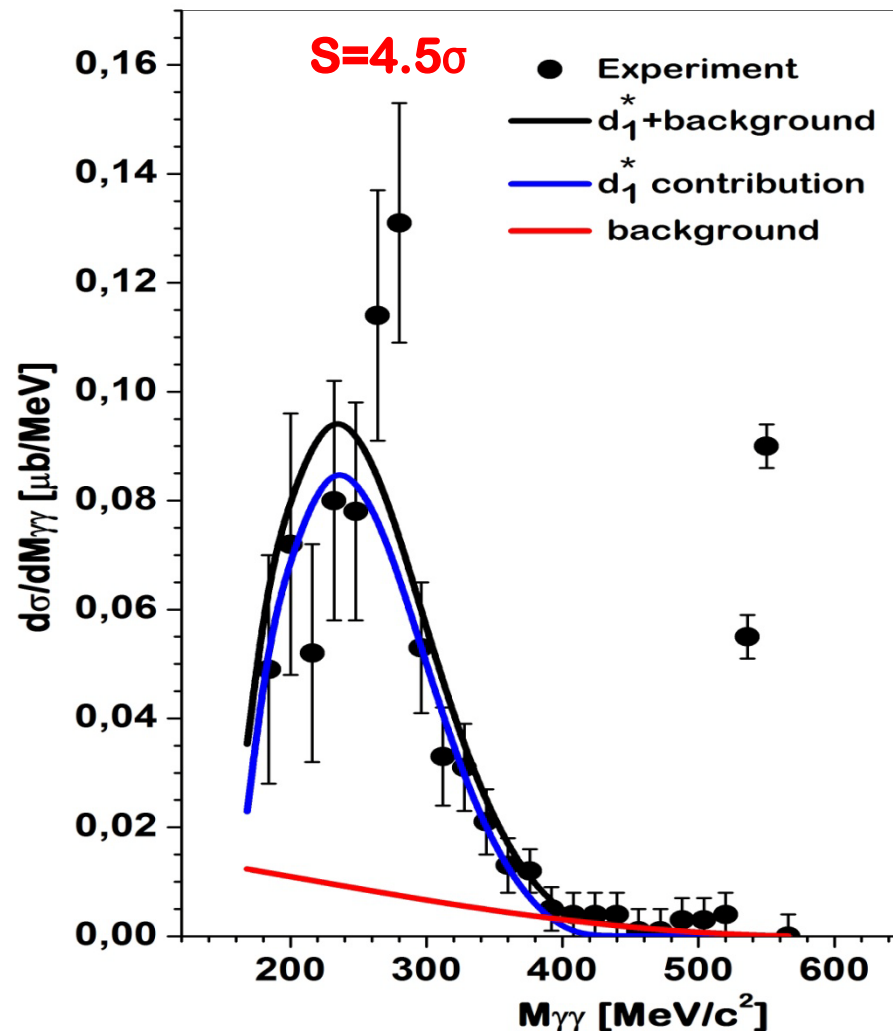
St.sign.= $N_S/(N_S+2N_B)^{1/2}$: 4.5σ & 3.5σ

Calculations: $|M(NN \rightarrow \gamma d_1^* \rightarrow pp \gamma\gamma)|^2$

$\Rightarrow d_1^*(1956) \rightarrow (N\Delta)_{\text{bound}}$

$M_{\gamma\gamma}^2 = (k_1+k_2)^2 = 2E_{\gamma 1} * E_{\gamma 2} * (1-\cos\theta_{12})$

$\chi^2=1.1/\text{dof}$



Why the Celsius-Wasa Collaboration did not find the dibaryon signal in their $pp \rightarrow pp\gamma\gamma$ data at 310 MeV?

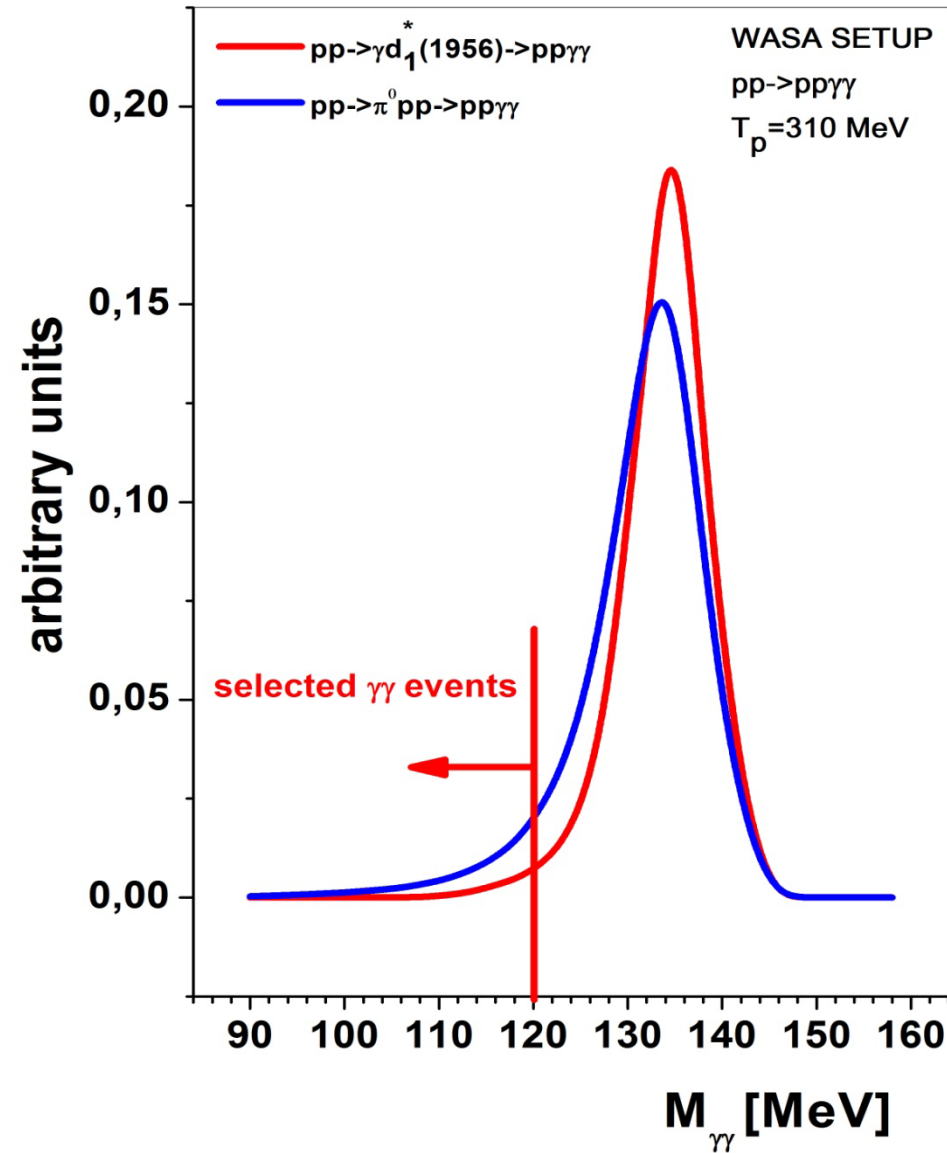
$pp \rightarrow \gamma pp$ at $T_k = 310$ MeV

H.Calen et al., Phys. Lett. B427, 248(1998).

To remove the γ -background due to the $pp \rightarrow \pi^0 pp \rightarrow \gamma\gamma pp$ process $\gamma\gamma$ events also were collected. From these events were removed those connected with the π^0 - decay to search for the d^*_1

Two-photon inv. mass spectra were calculated for the $pp \rightarrow \gamma d^*_1(1956) \rightarrow pp\gamma\gamma$ and $pp \rightarrow pp\pi^0 \rightarrow pp\gamma\gamma$ channels of the reaction $pp \rightarrow pp\gamma\gamma$ for the geometry and kinematic of the experiment PLB427,248 (1998).

So, all events (at least most of them) associated with the $d^*_1(1956)$ were removed!



TWO PHOTON INVARIANT MASS SPECTRUM

Experiment

Reaction: $p + {}^{12}\text{C} \rightarrow \gamma X$

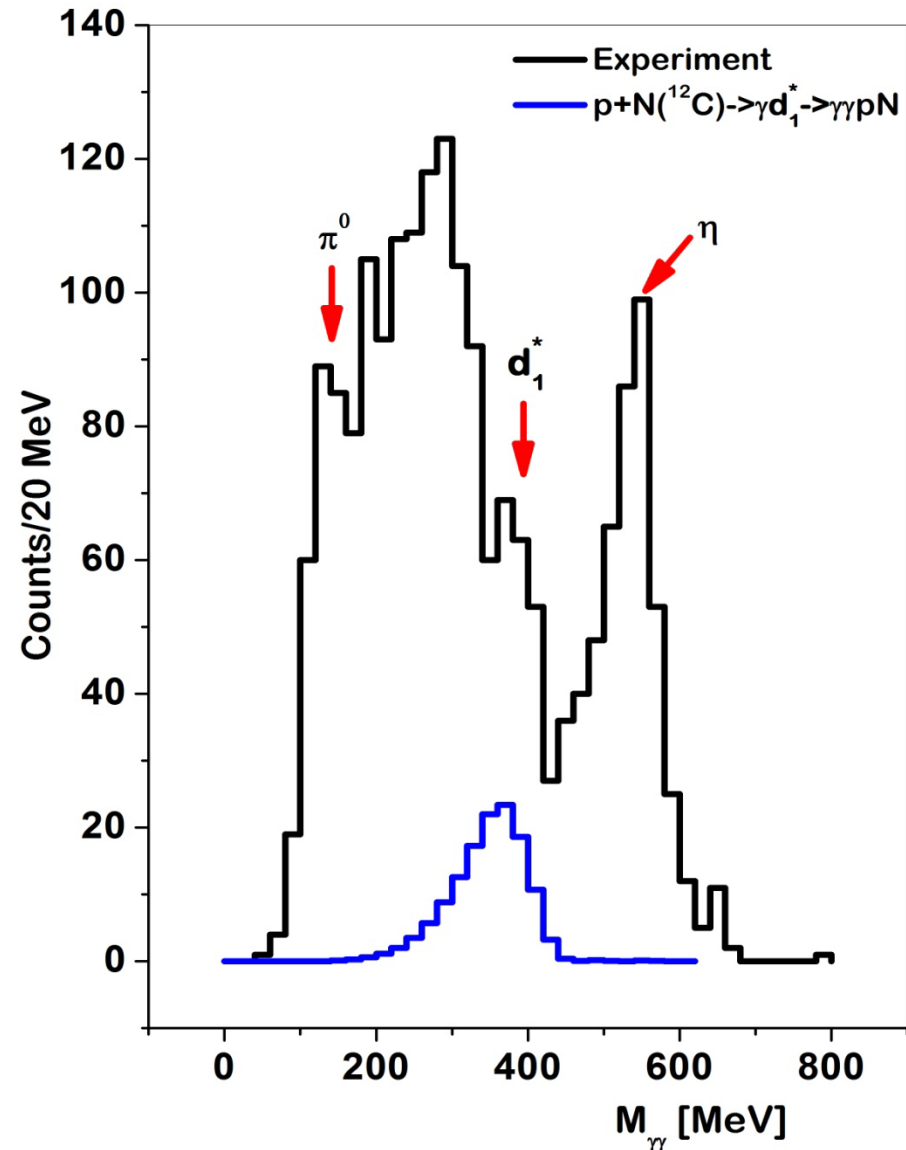
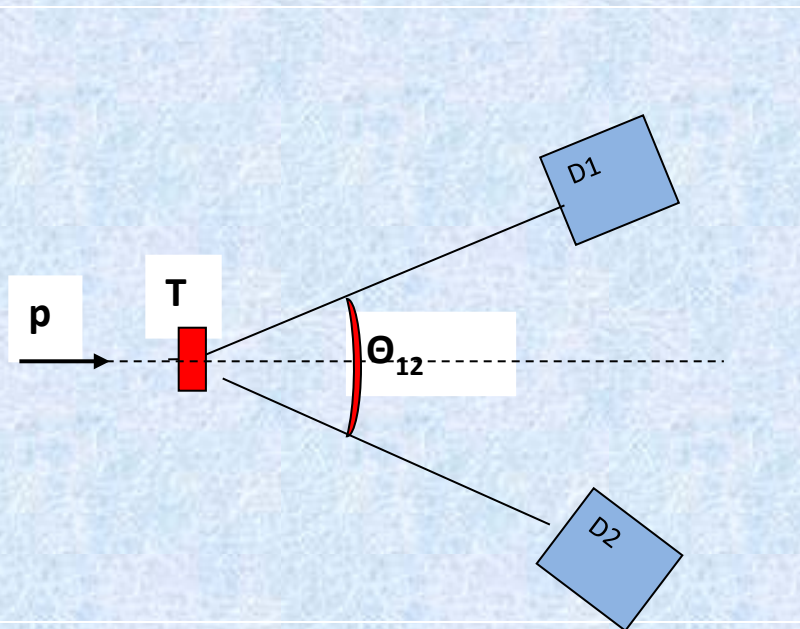
$T_p = 1300 \text{ MeV}$

PINOT Spectrometer $\Theta_{12} = 55^\circ$

Exp. Data: C. De Oliveira Martins et al.

Braz. Jour. Of Phys. V.31,

n.31, p. 533(2001)



TWO PHOTON INVARIANT MASS SPECTRUM

Experiment

Experiment

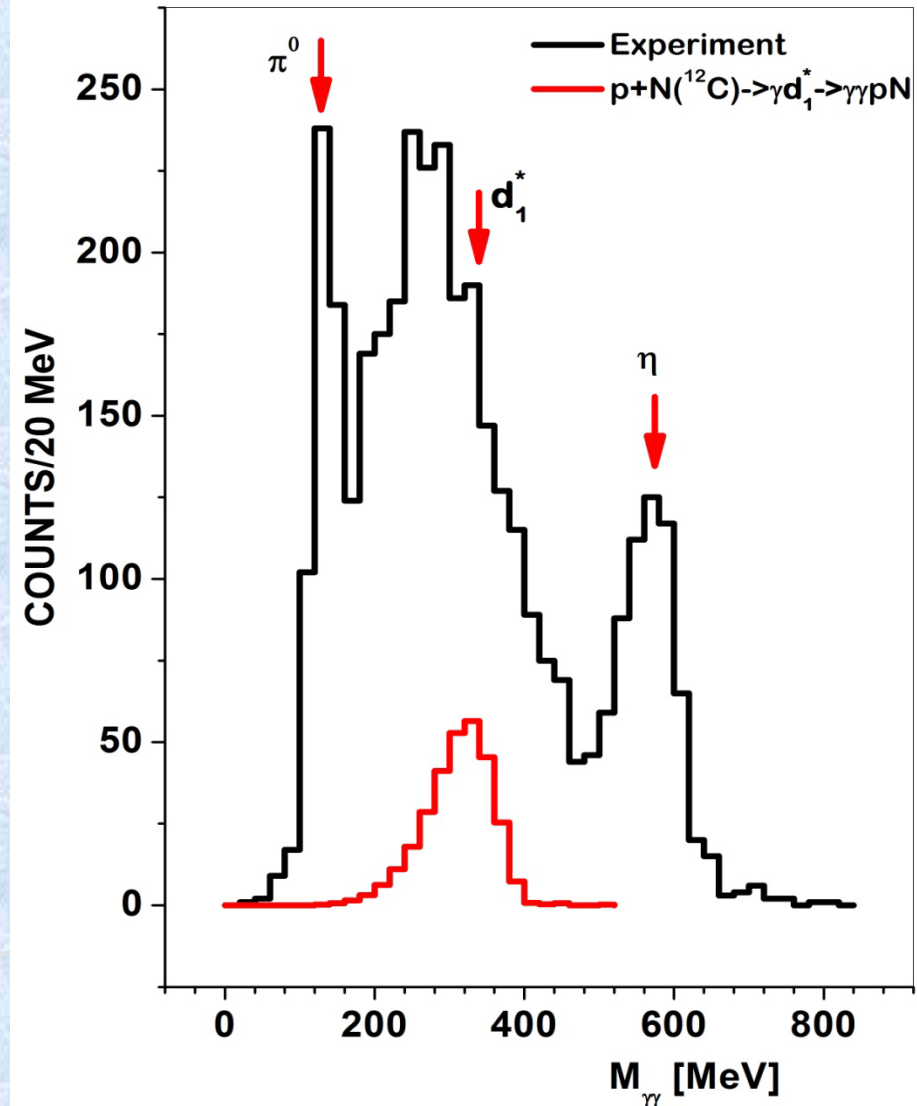
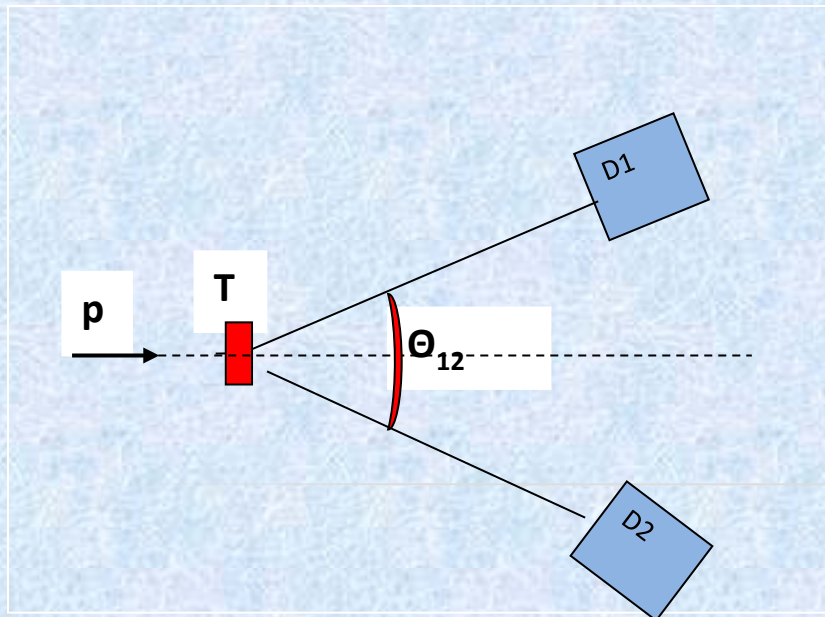
Reaction: $p + {}^{12}\text{C} \rightarrow \gamma X$

$T_p = 1150 \text{ MeV}$

PINOT Spectrometer

$\Theta_{12} = 66^\circ$

Exp. Data: E. Chiavassa et al., Europhys. Lett. V.41, p. 365(1998)



Inclusive photon energy spectrum for the $pd \rightarrow \gamma X$ reaction.

A.S. Khrykin, Nucl. Phys. A721, 625c (2003)

Reaction : $p+d \rightarrow \gamma X$

Experiment

$T_p = 195 \text{ MeV}$; $\theta_\gamma = 90^\circ$

Michigan State group

J. Clayton et al., Phys. Rev. C45, 1810 (1992).

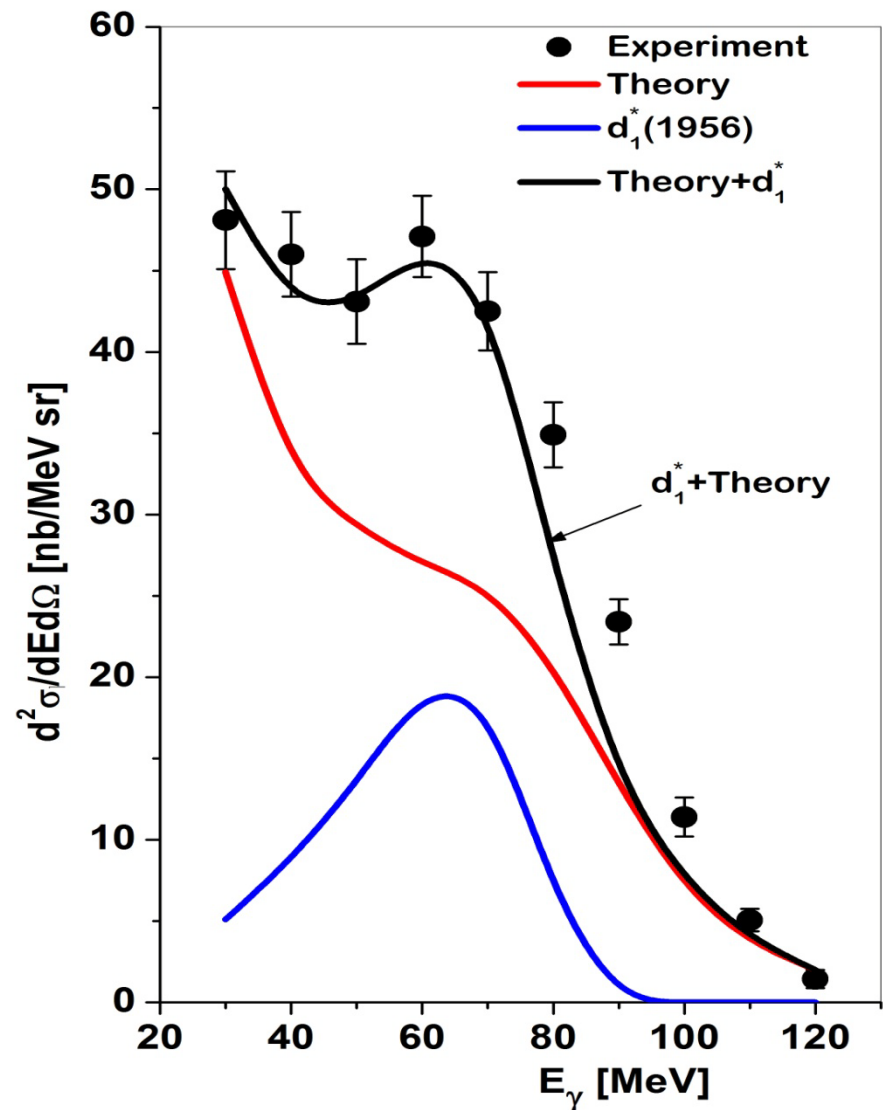
$T_p = 200 \text{ MeV}$; $\theta_\gamma = 90^\circ$

Grenoble group

J.A. Pinston et al.,
Phys.Lett. B249, 402(1990)

Theory

K.Nakayama,
Phys.Rev. C45, 2039 (1992).



Inclusive photon energy spectrum for the $np \rightarrow \gamma X$ reaction

A.S. Khrykin, Nucl. Phys. A721, 625c (2003)

Reaction: $n+p \rightarrow \gamma X$

Experiment

At the Saturne National Laboratory in Saclay.

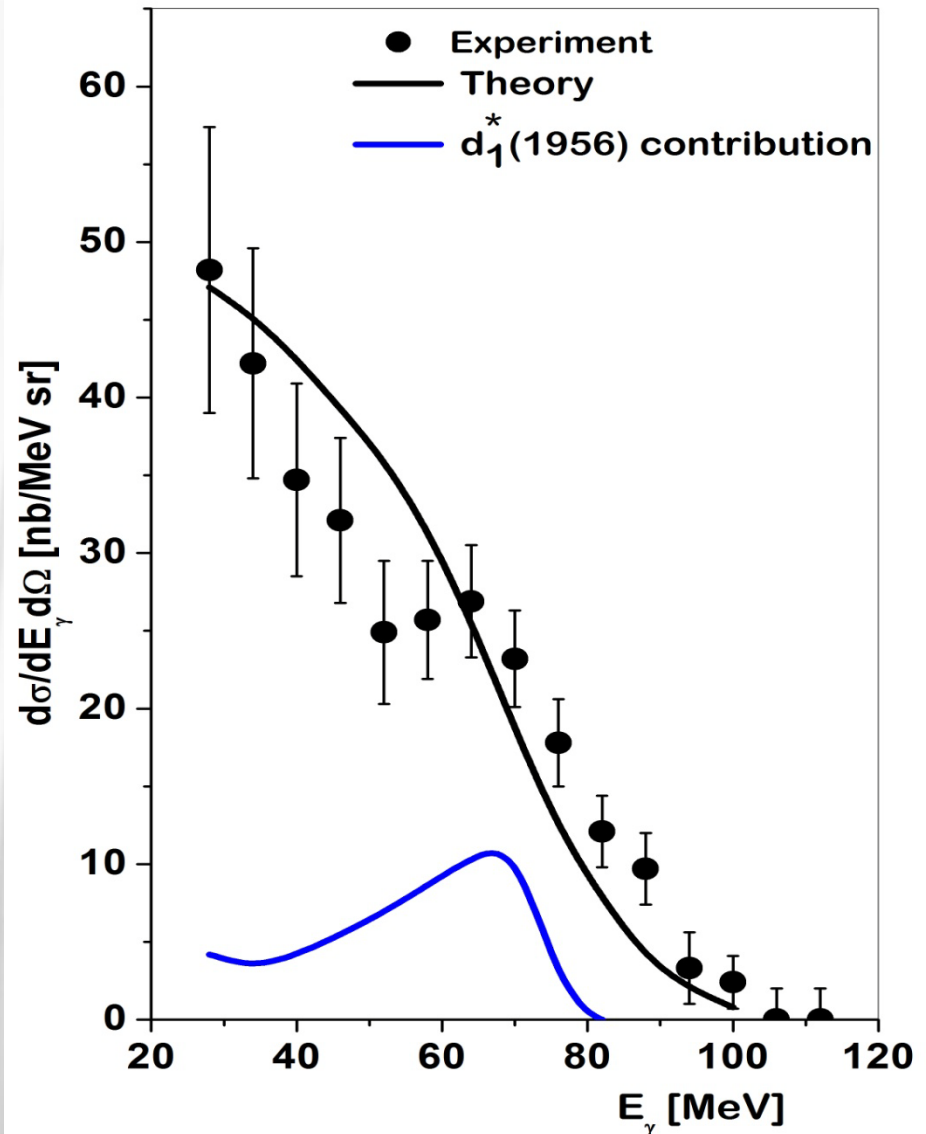
$T_n = 170 \pm 35$ MeV; $\theta_\gamma = 90^\circ$

F. Malek et al., Phys. Lett. B266, 255 (1991).

Theory

M. Schafer et al.

Z. Phys. A 339, 391 (1991).



Inclusive photon spectrum for the $p + {}^{12}\text{C} \rightarrow \gamma X$ reaction

A.S. Khrykin, Nucl. Phys. A721, 625c (2003)

Reaction: $p + {}^{12}\text{C} \rightarrow \gamma X$

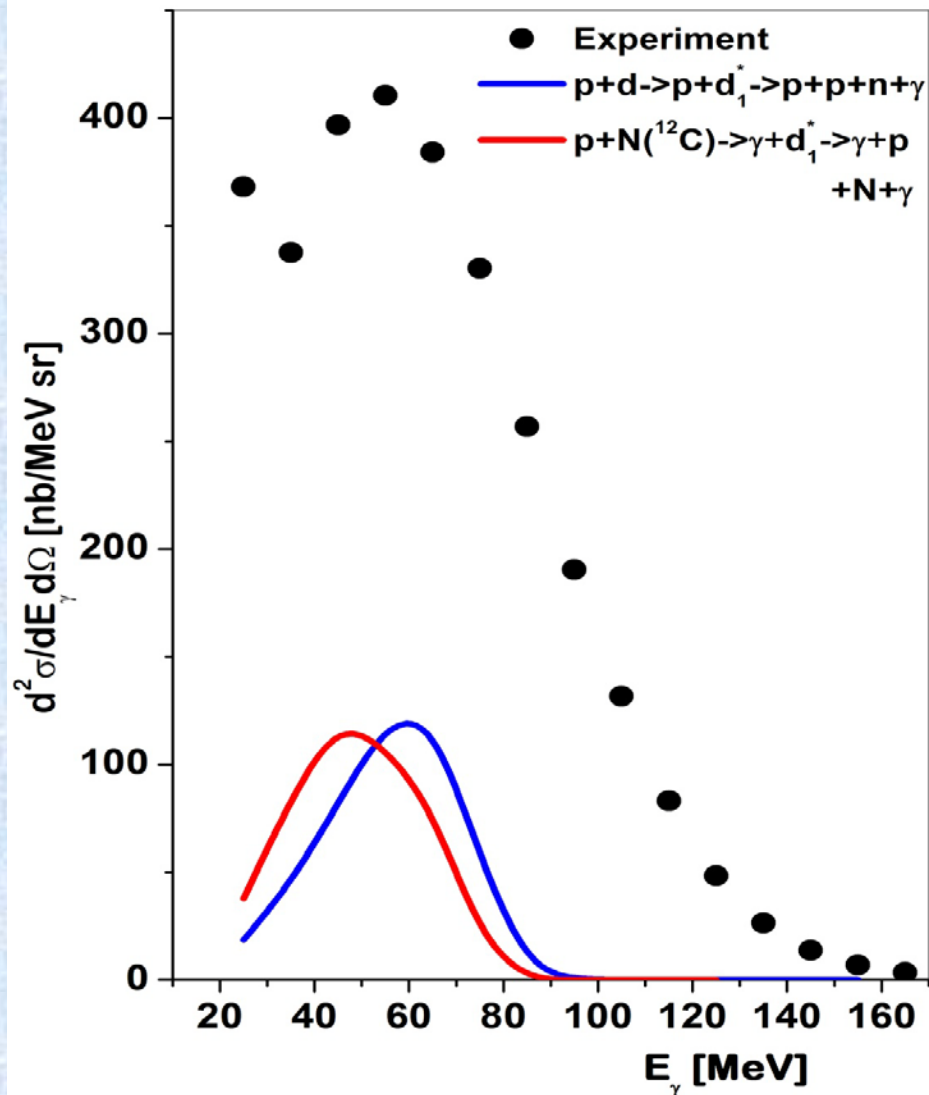
Experiment

At the Orsay synchrocyclotron.

$T_p = 200 \text{ MeV}$; $\theta_\gamma = 90^\circ$

J.A.Pinston et al., Phys.Lett.
B249, 402(1990).

CNWF:O.Benhar et al.,
Phys. Lett. B177,135(1986)



Inclusive photon spectrum for the $p + {}^{12}\text{C} \rightarrow \gamma X$ reaction

Reaction: $p + {}^{12}\text{C} \rightarrow \gamma X$

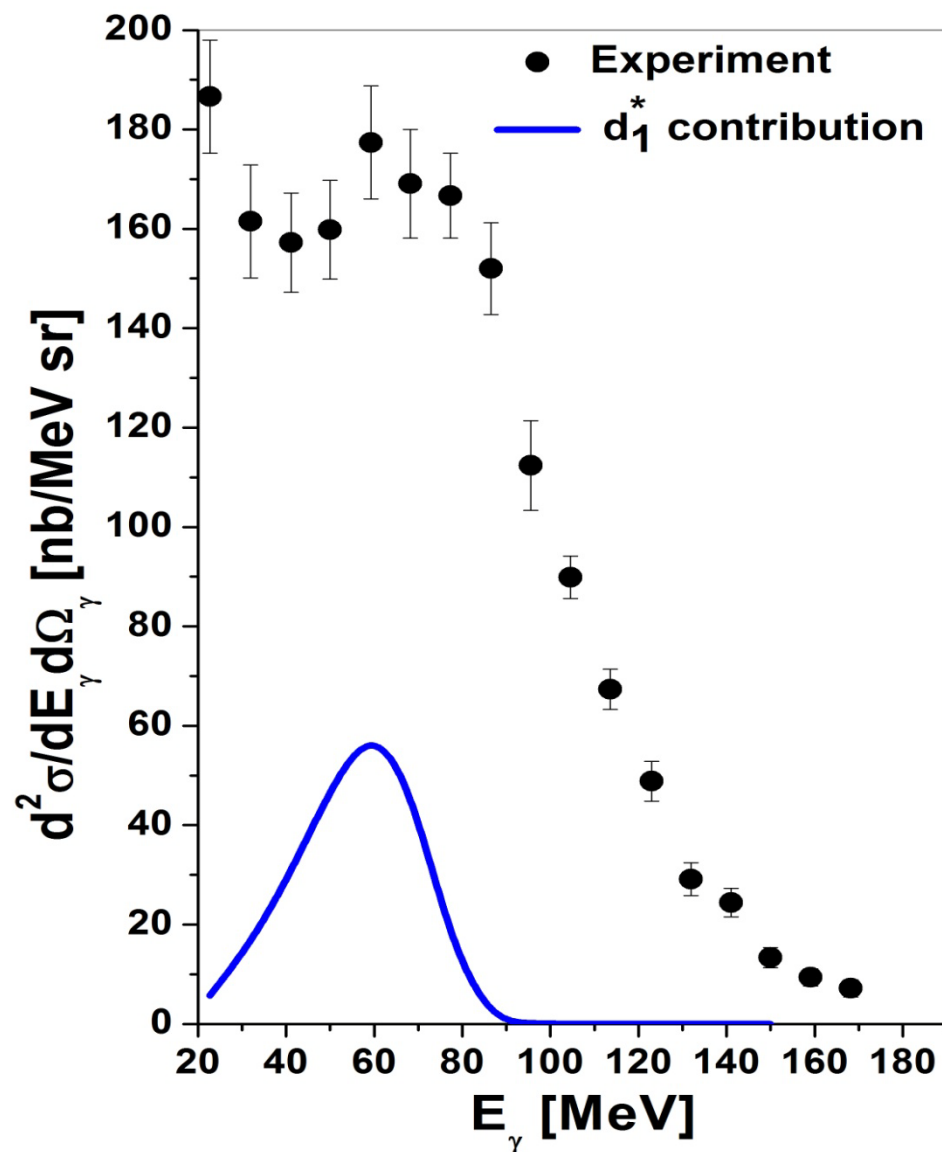
Experiment

At the AGOR facility of the KVI
Groningen.

Photon Spectrometer TAPS

$T_p = 190 \text{ MeV}$; $\theta_\gamma = 75^\circ$

M.J. van Goethem et al., Phys.Rev.
Lett 88,122302(2002).



A new mechanism for dielectron production in NN collisions.

A. S. Khrykin, in Proceedings of the 12th International Conference on Meson-Nucleon Physics and the Structure of the Nucleon, College of William and Mary, Williamsburg, Virginia USA, May 31-June 4, 2010, will be published by the AIP.

DLS Measurements: $A+A \rightarrow e^+e^-X$ for C+C and Ca+Ca at 1.04 A GeV \Rightarrow **Substantial Excess** of the e^+e^- pair yield in the mass region $0.2 < M_{ee} < 0.6$ GeV

HADES Measurements: $A+A \rightarrow e^+e^-X$ spectra for C+C collisions at 1A and 2A GeV. The DLS finding was confirmed.

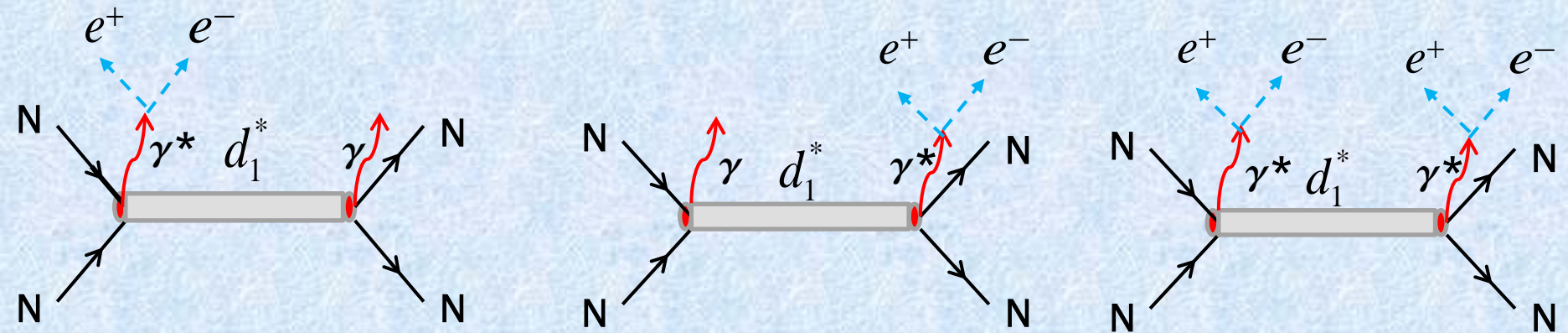
Conventional mechanisms of dielectron production at 1-2 GeV per nucleon:

- Dalitz decay of the π^0 -, η -, and ω mesons and baryon resonances $\Delta(1232)$, $N^*(1520)$,...
- Direct decay of the π^0 -, ρ -, and ω - mesons
- Bremsstrahlung in NN and $\pi\pi$ collisions
- Pion annihilation

Up to now there is no clear conception of the origin of such an **excess**.

If the resonance d_1^* (1956) really exists then together with its formation (decay) channel with a real photon the same channels with a virtual, massive photon should also take place. Conversion of such photons into e^+e^- pairs would give rise to a new source of dielectrons.

Dibaryon mechanism for dielectron production in NN Collisions



The kinematical region for a photon mass overlaps with that where the shortage of dielectrons of interest was observed.

We decided to examine whether a new source $NN \rightarrow \gamma^* d_1^*(1956) \rightarrow e^+e^- d_1^*$ can supply the observed shortages of e^+e^- pairs in the DLS data in the mass region $0.2 < M_{ee} < 0.6 \text{ GeV}$.

$$NN \longrightarrow \gamma^* d_1^* (1956) \longrightarrow e^+ e^- d_1^* (1956)$$

$$p_a + p_b = p_1 + p_2 + p_3$$

$p = p_a + p_b$, p_a - and p_b - the four momenta of colliding nucleons,
 $p^2 = s$ – the total energy of the colliding nucleons inc.m.s.

$p_1(E_1, \vec{p}_1), p_2(E_2, \vec{p}_2)$ and $p_3(E_3, \vec{p}_3)$ -the four momenta of dielectrons and resonance.

$$\frac{d\sigma}{dM} = \frac{(2\pi)^4}{f} \int \prod_{i=1}^3 \frac{d^3 \vec{p}_i}{2E_i (2\pi)^3} |\mathcal{M}|^2 \delta(p - \sum_{i=1}^3 p_i) \cdot \delta(M - M(\vec{p}_1, \vec{p}_2))$$

$$f = 4\sqrt{(p_a p_b)^2 - m_a^2 m_b^2}$$

M - the invariant mass of the e^+e^- - pair

$|\mathcal{M}|^2$ – the matrix element for the transition $NN \rightarrow e^+e^-d_1^*$

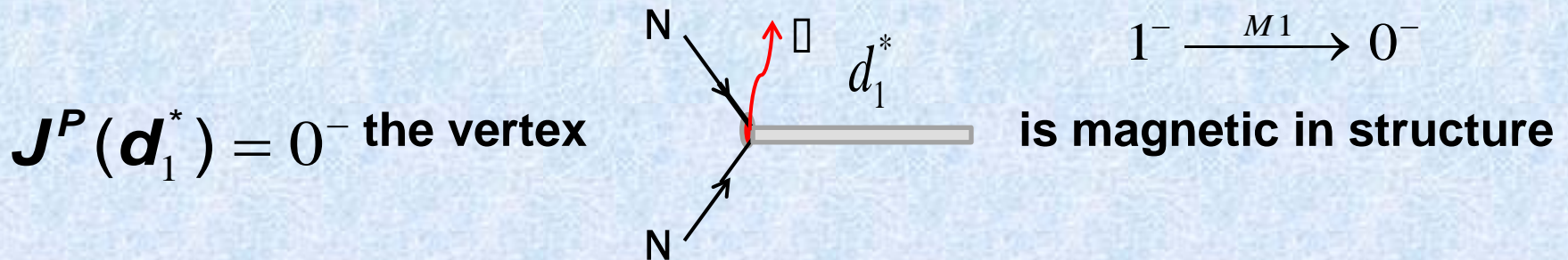
$$\mathcal{M} = \frac{e^2}{k^2} j^\mu J_\mu, \quad j_\mu = \langle e^+ e^- | \hat{j} | \mathbf{0} \rangle, \quad J_\mu = (J_0, \vec{J}) = \langle d_1^* | \hat{J}_\mu | NN \rangle$$

$$\vec{J} = \vec{J}_T + \vec{J}_L$$

$$|\mathcal{M}|^2 = \frac{1}{4} \frac{e^4}{M^4} \frac{1}{2m_e^2} \left\{ M^2 |\vec{J}_T|^2 - |\vec{J}_T \vec{q}_T|^2 + \frac{M^2}{k_0^2} \left(1 - \frac{M^2}{k_0^2} |\vec{J}_L \vec{q}_L|^2 \right) \right.$$

$$\left. - 2 |\vec{J}_L \vec{q}_L| |\vec{J}_T \vec{q}_T| \right\}, \quad k = (k_0, \vec{k}) = p_1 + p_2, \quad q = (q_0, \vec{q}) = p_1 - p_2$$

O. Scholton et al. PRC71,034005(2005)



$$|\mathcal{M}_E|^2 = e^2 |\vec{J}_T|^2 = C \cdot (p_a \cdot k)(p_b \cdot k)$$

Hadronic current for the magnetic transition is transverse, $\vec{J}_L = 0$

$$|\mathcal{M}|^2 = \frac{1}{4} \frac{e^4}{M^2} \frac{1}{2m_e^2} \left(|\vec{J}_T|^2 - \frac{1}{M^2} |\vec{J}_T \vec{q}_T|^2 \right) \cong \frac{1}{4} \frac{e^4}{M^2} \frac{1}{2m_e^2} |\vec{J}_T|^2$$

$$\left| \mathcal{M}_{NN \rightarrow e^+ e^- d_1^*} \right|^2 = \frac{N}{M^2} |\mathcal{M}_E|^2 |F(M^2)|^2$$

$$\left| F(M^2) \right|^2 = \frac{m_\rho^4 - m_\rho^2 \Gamma_\rho^2}{(m_\rho^2 - M^2)^2 + m_\rho^2 \Gamma_\rho^2}$$

N-is the normalization constant

The calculations \Rightarrow Monte Carlo method. Event generator \Rightarrow GENBOD. It used to randomly generate four momenta of the outgoing particles of the explored reaction. The probability of any event has been given its weight:

$$WT = \left| \mathcal{M}_{NN \rightarrow e^+ e^- d_1^*} \right|^2$$

Energy resolution: by procedure of a spectrum smearing with a Gaussian distribution with the corresponding σ .

$$\frac{\sigma_{\text{tot}}^{ee}}{\sigma_{\text{tot}}^\gamma} = \frac{\alpha}{\pi} \left[\frac{2}{3} \ln\left(\frac{\Delta M}{m_e}\right) - \frac{5}{9} + \frac{1}{3} I_1 + O\left(\left(\frac{m_e}{\Delta M}\right)^2\right) \right], \quad I_1 \approx 1$$

**B.E.Lautrup and J.Smith,
PRD3,1122(1971)**

**Contribution of the $pp \rightarrow e^+ e^- d_1^*$ mechanism to the
DLS $pp \rightarrow e^+ e^- X$ data**

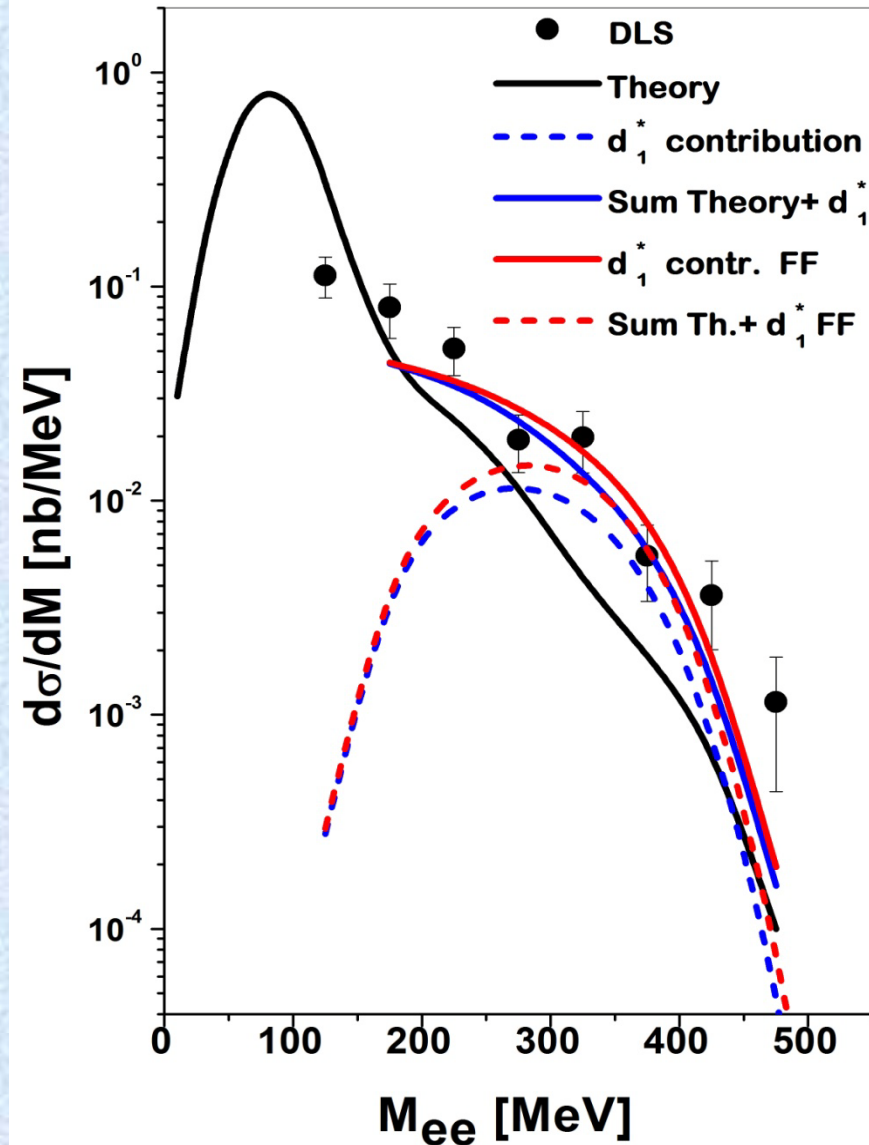
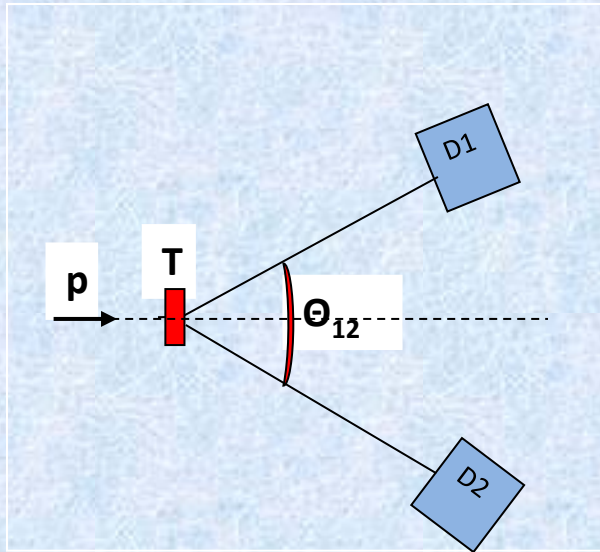
DLS data: W.K. Wilson et al., Phys. Rev. C **57**, 1865(1998)

Theoretical data: Amand Faessler et al., J. Phys. G **29**, 603(2003)

DLS Experiment

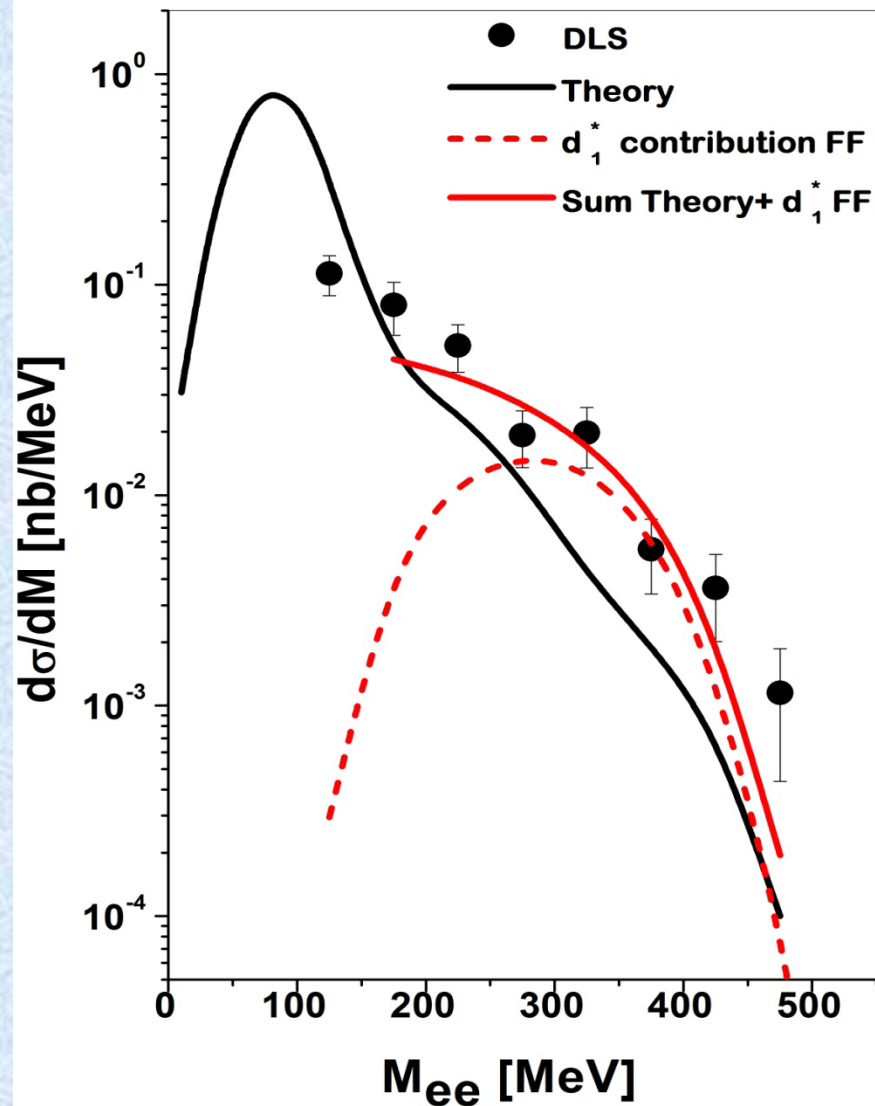
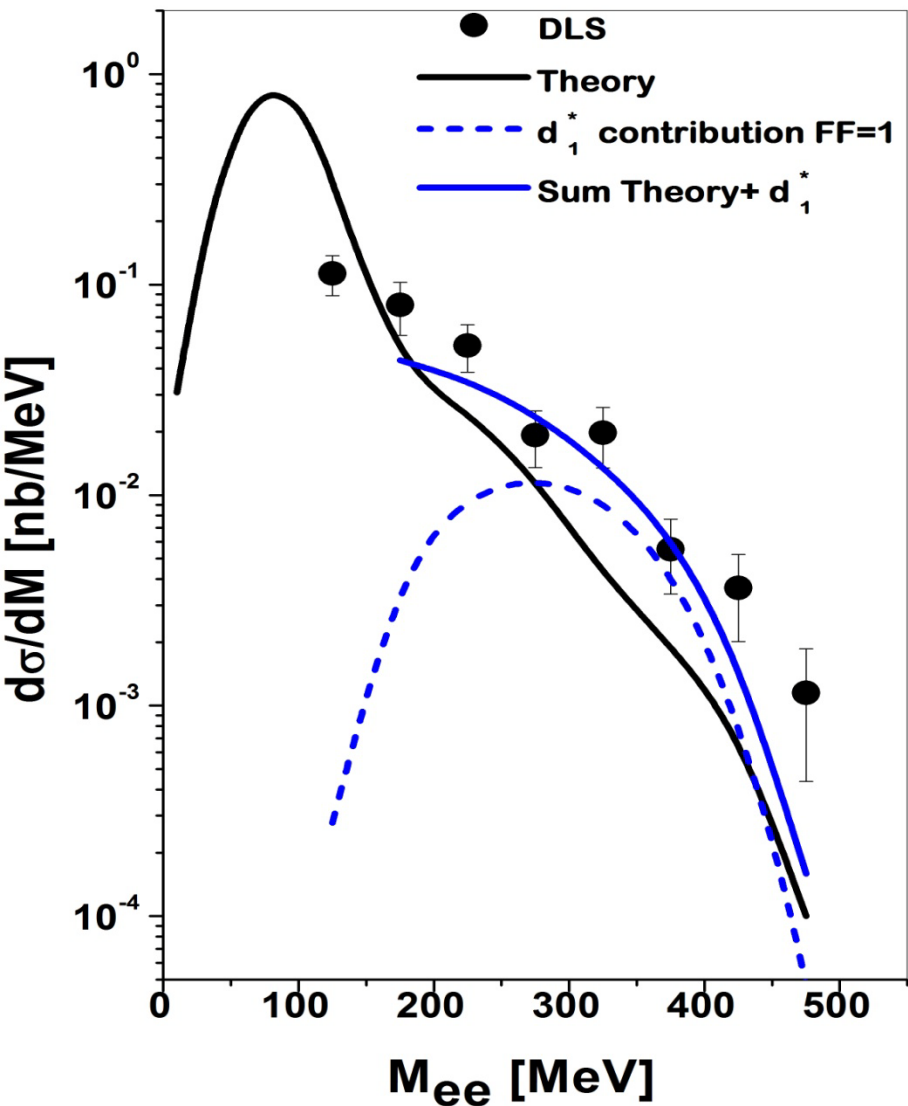
Reaction: $p+p \rightarrow e^+ e^- X$

$T_p = 1.04$ GeV



DLS data: W.K. Wilson et al., Phys. Rev. C **57**, 1865 (1998)

Theoretical data: Amand Faessler et al., J. Phys. G **29**, 603 (2003)



***Contribution of the $NN \rightarrow e^+e^-d_1^*$ mechanism to the
DLS $^{12}\text{C} + ^{12}\text{C} \rightarrow e^+e^-X$ data***

DLS data: R.J. Porter et al., Phys. Rev. Lett. **79**, 1229 (1997)

Theoretical data: E.L. Bratkovskaya et al., Nucl. Phys. **A634**, 168 (1998)

DLS_Experiment

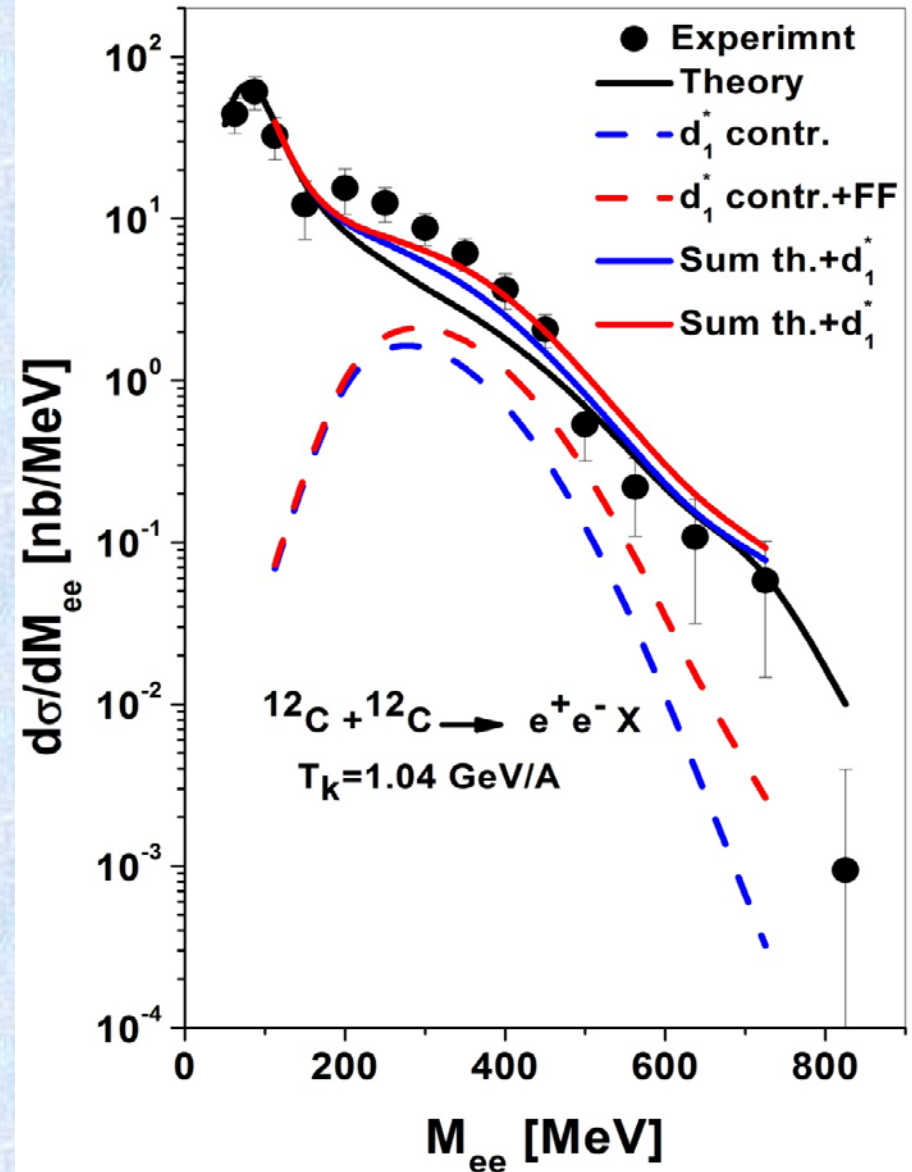
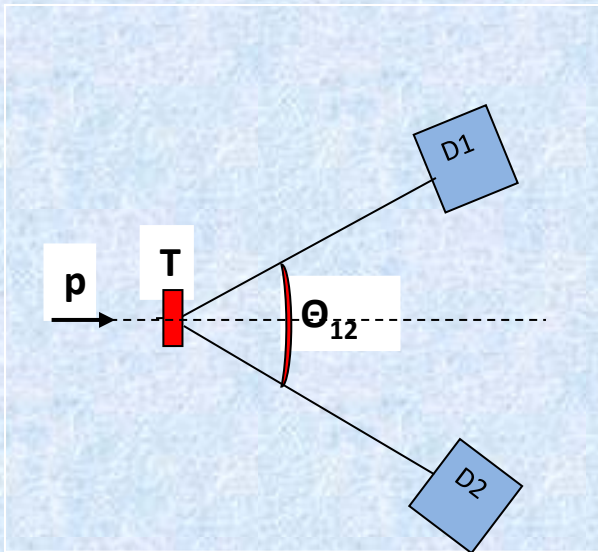
Reaction: $^{12}\text{C} + ^{12}\text{C} \rightarrow e^+ e^- X$

$T = 1.04 \text{ GeV/A}$

Mass resolution $\Delta M/M = 10\%$

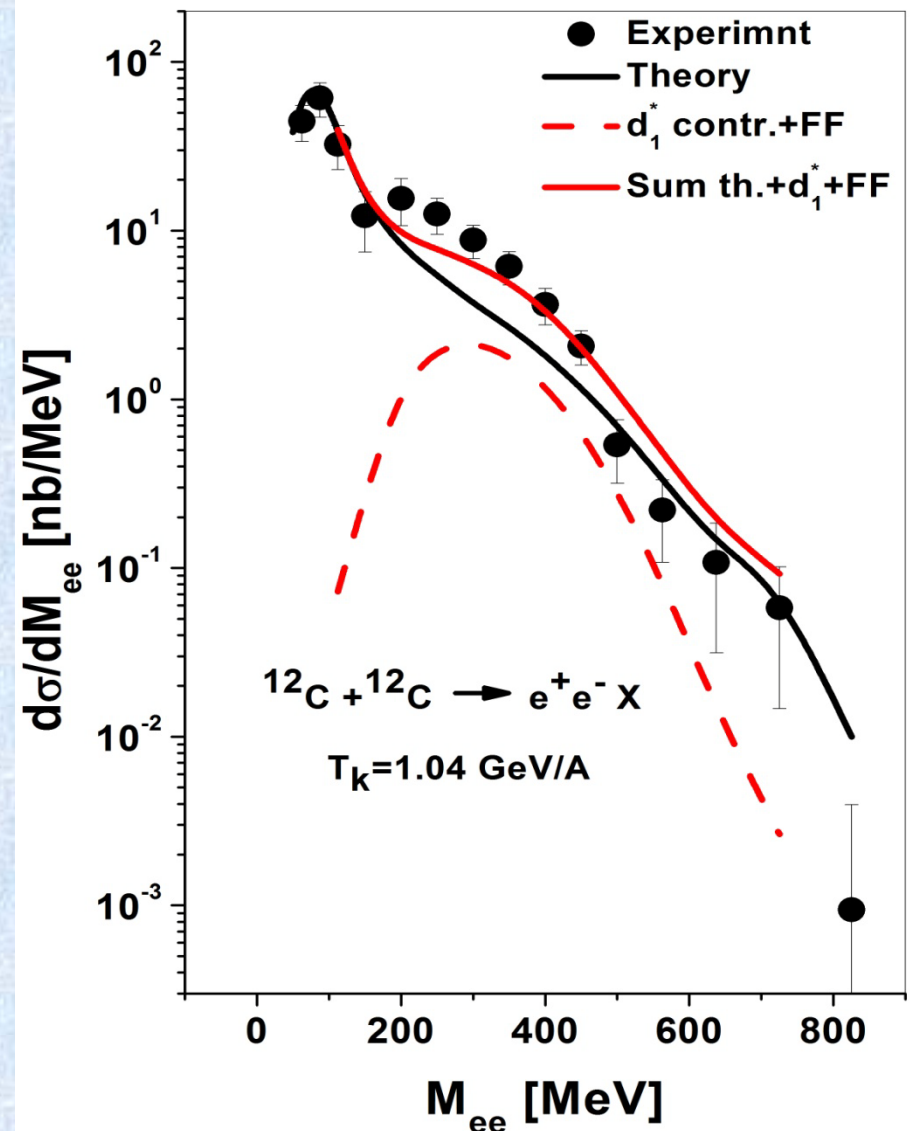
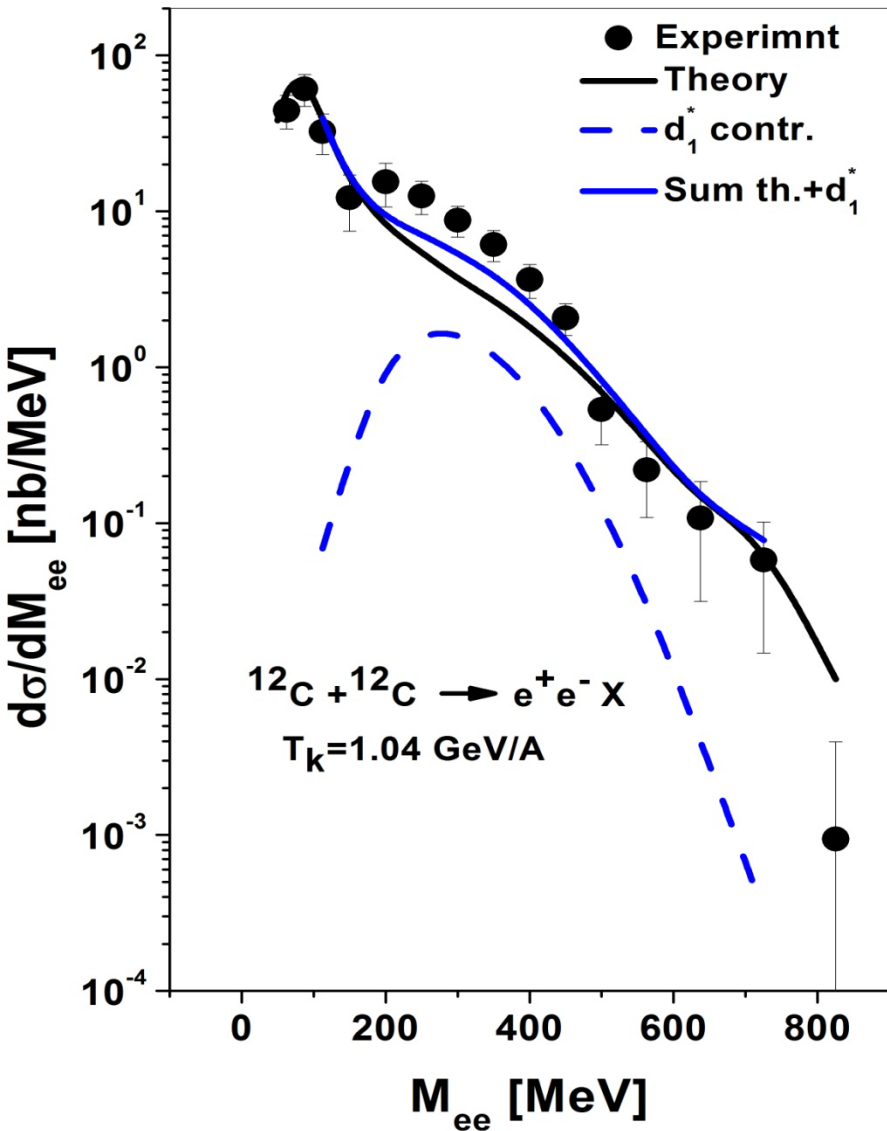
Filter v4.1p, CNWF: O. Benhar et al.,

Phys. Lett. **B177**, 135 (1986)



DLS data: R.J. Porter et al., Phys. Rev. Lett. **79**, 1229 (1997)

Theoretical data: E.L. Bratkovskaya et al., Nucl. Phys. **A634**, 168 (1998)



Conclusions

- ❖ The presented data strongly suggest the existence of the NN-decoupled dibaryon resonance $d_1^*(1956)$.
- ❖ The idea of the existing of this resonance can be confirmed or refuted by the direct measurement of the photon energy spectrum of the reaction $pp \rightarrow \gamma\gamma pp$ at energy below the pion production threshold or the missing mass spectrum of the reactions $pp \rightarrow e^+e^-X$.