Study of the reaction $np \rightarrow np\pi^+\pi^$ at intermediate energies

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<u>Outlook</u>

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1. Introduction: study of inelastic np interactions at accelerator facility of LHEP JINR

- Quasimonochromatic neutron channel: δP≈2.5%, P₀=1.25; 1.43; 1.73; 2.23; 3.10; 3.83; 4.10 and 5.20 GeV/c, 4π geometry.
- Cross-sections of inelastic np interactions: (black circles – our data)



The unique of fullness and precision data are obtained. It permits to carry out the detailed study of inelastic np interactions in a wide region of energies

2. Reaction $np \rightarrow np\pi^+ \pi^-$ at $P_0 > 3$ GeV/c

It is characterized by:

- plentiful production of the Δ resonance,



- large peripherality of the secondary nucleons

To study the mechanism of the reaction it was chosen the model of reggeized π-exchange (OPER), developed in ITEP [L.Ponomarev. Part. and Nucl., v.7(1), pp. 186-248, 1976, JINR, Dubna].

The advantages of OPER model are:

- small number of free parameters (3 in our case),
- wide region of the described energies (2 ÷ 200 GeV),
- calculated values are automatically normalized to the reaction cross-section.

The following main diagrams correspond to the reaction $np \rightarrow np\pi^+ \pi^-$ within the framework of OPER model:



Fig.1

Matrix element for the diagrams 1a,1b и 1c is written in the following form:

$$\mathsf{M}_{1}= T_{\boldsymbol{\pi}\boldsymbol{N}\to\boldsymbol{\pi}\boldsymbol{N}} \cdot \mathbf{F}_{2} \cdot T_{\boldsymbol{\pi}\boldsymbol{N}\to\boldsymbol{\pi}\boldsymbol{N}} / (\mathsf{t} - \mathsf{m}_{\pi}^{2}),$$

where $T_{\pi N \to \pi N}$ – amplitude of elastic $\pi N \to \pi N$ scattering off mass shell, F_2 – form-factor, going away off mass shell of $T_{\pi N \to \pi N}$ amplitudes, $1 / (t - m_{\pi}^2) - \pi$ -meson propagator.

The data of elastic $\pi N \rightarrow \pi N$ were taken from *PWA* [R.A. Arndt et al. IJMP A18(3), 2003, p. 449]

The analysis shows, that interference between diagrams 1a, 1b и 1c is negligible [A.P.Jerusalimov et al. JINR Rapid Communications, v.35(2) pp.21-26, 1989, JINR, Dubna].

The study has shown that it is not necessary to take into account the contribution of the «hanged» diagrams into the reaction cross-sections at $P_0 < 10$ GeV/c



It was shown in [G.W. van Appeldorn et al, NP B156 (1979),pp. 110-125] that the use of some specific cuts permits to select the kinematic region of the reaction $np \rightarrow np\pi^+ \pi^-$ in which the contribution of the diagrams **1a**, **1b** μ **1c** consists up to 95 % at P₀ > 3 GeV/c.

Fig.2 shows some distributions for the reaction np \rightarrow np $\pi^+ \pi^-$ for this region at P₀ = 5.20 GeV/c (blue curves – results of calculations using OPER model)



But the diagrams shown in Fig.1 are insufficient to describe totally the characteristics of the reaction $np \rightarrow np\pi^+ \pi^-$. It is necessary to take into account the diagrams of the following type:



the matrix element for which is written in the following form:

$$\begin{split} & \mathsf{M}_3 = G \; u(q_N) \gamma_5 u(\mathsf{Q}_N) \cdot \mathbf{F}_1 \cdot \mathcal{T}_{\pi N \to \pi \pi N} / (\mathsf{t} - \mathsf{m}_{\pi}^{-2}), \\ & \text{where } \mathsf{T}_{\pi N \to \pi \pi N} \text{ - off mass shell amplitudes of inelastic } \pi N \to \pi \pi N \text{ scattering} \\ & \text{that are known much worse than } \mathsf{T}_{\pi N \to \pi \pi N} \text{ amplitudes.} \\ & \text{Therefore it is necessary to do a parametrization of the inelastic } \pi N \to \pi \pi N \text{ scattering} \\ & (\text{see Appendix}). \end{split}$$

It permits to get a good description of the experimental characteristics of the reaction $np \rightarrow np\pi^+ \pi^-$ at P₀=5.20 GeV/c (Fig. 5) taking into account OPER diagrams shown in Fig.1 and Fig.3 :



Fig. 4

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3. Reaction pbar $p \rightarrow pbar p \pi^+ \pi^-$ at P₀ = 7.23 GeV/c

Using OPER model we try to describe the experimental distributions from the reaction **pbar p** \rightarrow **pbar p** $\pi^+ \pi^-$ at P₀ = 7.23 GeV/c

[G.W. van Appeldorn et al, NP B156 (1979),pp. 110-125]



It is observed a good agreement between experimental data and theory.

4. Reaction np \rightarrow np $\pi^+ \pi^-$ at P₀ < 3 GeV/c

The study of effective mass spectra of np – combinations at $P_0=1.73$ and 2.23 GeV/c shows the clear peack close the threshold ($M_{np}=m_n+m_p$) that can not be described within the framework of OPER-model c using the diagrams 1a, 1b, 1c, 3a, 3b, 3c μ 3d.



Р₀=2,23 ГэВ/с

Р₀=1,73 ГэВ/с

The model of Regge poles with baryon exchange and nonlinear trajectories, suggested in [A.B. Kaydalov and A.F. Nilov. **Ya**F, v.**41**(3),pp. 768-776, 1985 ; **YaF**, v.**52**(6), pp. 1683-1696, 1990] was used to describe these features.

The following diagrams of one baryon exchange (OBE) were taken into account within the framework of this model:





The vertex function of elastic $np \rightarrow np$ scattering was calculated using the data from [NN and ND interactions – a compilation. UCRL-20000 NN, august 1970].

The vertex functions of $\Delta N \rightarrow np$, $NN \rightarrow \Delta N \lor \Delta N \rightarrow \Delta N$ scattering were calculated corresponding to [B.C. Барашенков и Б.Ф. Костенко. 4-84-761, JINR, Dubna, 1984]. In result one can get the good description of the experimental distribution from the reaction $np \rightarrow np\pi^+\pi^-$ at P₀ =1.73 and 2.23 GeV/c.



5. OPER model and other reactions

The other reactions of **np** interactions are scheduled to study by means of OPER model:

- $np \rightarrow pp\pi^{-}$ vertex functions 1 x 2- $np \rightarrow pp\pi^{-}\pi^{0}$ vertex functions 2 x 2 and 1 x 3- $np \rightarrow pp\pi^{+}\pi^{-}\pi^{-}$ vertex functions 2 x 3
- np \rightarrow pp $\pi^+\pi^-\pi^-\pi^0$ vertex functions 3 x 3
- np \rightarrow np $\pi^+ \pi^- \pi^-$ vertex functions 3 x 3

Similar reactions of p p -, pbap p - and $\pi N -$ interactions also can be described by OPER model.

The following reactions were simulated for HADES experiment:

 $\begin{array}{l} pp \rightarrow pp \ \pi^{+} \pi^{-} \ \text{at } T_{kin} = 3.5 \ \text{GeV} \\ np \rightarrow np \pi^{+} \ \pi^{-} \ \text{at } T_{kin} = 1.25 \ \text{GeV} \\ np \rightarrow np e^{+} \ e^{-} \ \text{at } T_{kin} = 1.25 \ \text{GeV} \\ \end{array}$ with vertex function of $\pi N \rightarrow N \ e^{+} \ e^{-}$

Since the $\pi N \rightarrow \pi N$ and $\pi N \rightarrow \pi \pi N$ vertex functions are taken in helicity representation it seems to be perspective to use OPER model for description of the reaction with polarized particles.

6. Conclusion

Reaction np \rightarrow np $\pi^+ \pi^-$ at <u>P₀ > 3 GeV/c</u> is characterized by the plentiful production of the Δ resonance and the large peripherality of the secondary particles. The experimental data are successfully described by the further development of OPER – model.

However at $\underline{P_0} < 3 \text{ GeV/c}$ it is necessary to take into account another mechanism of the reaction (such as OBE).

OPER – model permits to describe another N(barN)-N reactions with the production of some π -mesons.

The further development of OPER – model can be very promising to describe the production of e^+e^- -pairs in hadronic interactions.

OPER – model can be used as an effective tool to simulate various reactions of hadronic interactions.

<u>Appendix</u>: Parametrization of $\pi N \rightarrow \pi \pi N$ reactions

Within the framework of Generalized Isobar Model (GIM) [D.J. Herndon et al. PR D11, 3165 (1975); D.M.Manley and E.M. Saleski, PR D45, 4002 (1992)] $\pi N \rightarrow \pi \pi N$ reactions are described as quasi-two body ones: $\pi N \rightarrow N^*(\Delta^*) \rightarrow \Delta \pi$, $\rightarrow N^*(\Delta^*) \rightarrow N \rho$, $\rightarrow N^*(\Delta^*) \rightarrow N \varepsilon$, ε - S-wave of $\pi\pi$ scattering with I=0 ightarrow N*(Δ *) ightarrow N $^{*}_{1440}$ π , $\Delta \rightarrow N \pi$. with the consequent $ho
ightarrow \pi \pi$, decays : $\mathcal{E} \rightarrow \pi \pi$, $N^*_{1440} \rightarrow N \pi$. N*(1440) P11 D*(1600) P33 N*(1520) D13 D*(1620) S31 The parameters of the following N*(1675) D15 D*(1700) D33 The spin and N*(1680) F15 D*(1900) S31 isospin relations resonances (**** and ***) N*(1720) P13 D*(1905) F35 were taken were taken N*(2000) F15 D*(1910) P31 account from **RPP** N*(2080) D13 D*(1920) P33 N*(2190) G17 D*(1940) D33 D*(1950) F37

For quasi two-body reactions like $\mathbf{a} + \mathbf{b} \rightarrow \mathbf{c} + \mathbf{d}$ $\pi N \rightarrow N^*(\Delta^*) \rightarrow \Delta \pi$, $\rightarrow N^*(\Delta^*) \rightarrow N \rho$, $\rightarrow N^*(\Delta^*) \rightarrow N \epsilon$, $\rightarrow N^*(\Delta^*) \rightarrow N^*_{1440} \pi$,

one can write

$$d\sigma = \frac{1}{(2S_a+1)(2S_b+1)} \left(\frac{2\pi}{p}\right)^2 \sum_{\lambda_i} |\langle \lambda_d \lambda_c | T | \lambda_b \lambda_a \rangle|^2 \times dPS$$

where

$$<\lambda_d \lambda_c |T| \lambda_b \lambda_a > = \frac{1}{4\pi} \sum_j (2j+1) < \lambda_d \lambda_c |T_j| \lambda_b \lambda_a > e^{i(\lambda-\mu)\varphi} d^j_{\lambda\mu}(\vartheta)$$

 $\Lambda = \Lambda_a - \Lambda_b$, $\mu = \Lambda_c - \Lambda_d$ – helicity variables, $d'_{\lambda\mu}(\theta)$ – rotation matrixes.

The polarization components of the particles **c** and **d** from the reaction $\mathbf{a} + \mathbf{b} \rightarrow \mathbf{c} + \mathbf{d}$ is suiteable to express through the elements of the spin density matrix (for example, for particle **d**):

$$\rho_{mm'}^{d} = \frac{1}{N} \sum_{\lambda_c \lambda_b \lambda_a} \langle m' \lambda_c \mid T \mid \lambda_b \lambda_a \rangle^* \langle m \lambda_c \mid T \mid \lambda_b \lambda_a \rangle$$

where

normalization factor for $Sp \rho = 1$.

$$N = \sum_{m\lambda_c\lambda_b\lambda_a} | < m\lambda_c \mid T \mid \lambda_b\lambda_a > |^2$$

Example:

$$\pi + N \to N^*_{1680} \to \Delta + \pi \to (N + \pi) + \pi$$
$$< \lambda_{\Delta} |T| \lambda_N > = C_{3,0;\frac{1}{2},-\lambda_{\Delta}}^{\frac{5}{2},-\lambda_{\Delta}} C_{1,0;\frac{3}{2},-\lambda_{\Delta}}^{\frac{5}{2}} d_{-\lambda_N,\lambda_{\Delta}}^{\frac{5}{2}}(\vartheta) \times R_J$$

 R_{J} is taken in Breight-Wigner form

Then it is easy to get the angular distribution of Δ (in CMS):

$$\frac{d\sigma(s,t)}{d\Omega} \sim (1 + 2\cos^2\vartheta_{\Delta})|R_J|^2 = (1 + 2\cos^2\vartheta_{\Delta})BW(\sqrt{s}, M_R, \Gamma_R)$$

If particle **d** is unstable: $\mathbf{d} \to \mathbf{\alpha} + \mathbf{\beta} \ (\mathbf{d} \to \mathbf{\Delta} + \mathbf{\pi})$ then in the rest system of the particle $\mathbf{d} \ (\mathbf{d} \to \mathbf{\Delta} + \mathbf{\pi})$

$$W_{\Delta}(\vartheta,\varphi) = \frac{3}{4\pi} \left\{ \rho_{33} \sin^2 \vartheta + \frac{1}{3} \rho_{11} (1 + 3\cos^2 \vartheta) - \frac{2}{\sqrt{3}} Re \rho_{3-1} \sin^2 \vartheta \cos 2\varphi - \frac{2}{\sqrt{3}} Re \rho_{31} \sin 2\vartheta \cos \varphi \right\}$$

is the normalized angular distribution of the decay products.

Comparison with experimental data

The following cross-sections were calculated using **GIM**:



One can see a satisfactory description of cross-sections, except $\pi^* \rho \rightarrow n \pi^* \pi^*$ May be it is necessary to put into **GIM S**-wave of $\pi^*\pi^*$ scattering with I=2.

Some distributions of the reaction $\pi \bar{\rho} \rightarrow n \pi^{\dagger} \pi \bar{\rho}$ were calculated at various energies to study a quality of the application of GIM :



[J. Dolbeau et al. NP B78, 233(1974)]

It is observed a good agreement between experimental data and theory.