Rare decay $\pi_0 \rightarrow e^+ e^$ as a filter for low mass dark matter

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Introduction

Model Independent Approach to $\pi^0 \rightarrow e^+e^-$ Decay and KTeV Data

Other $P \rightarrow I^+ I^-$ decays

 $\pi^{0} \rightarrow e^{+}e^{-}$ Decay and Dark Matter

Conclusions

Introduction

Abnormal people are looking for traces of Extraterrestrial Guests Abnormal Educated people are looking for hints of New Physics

Cosmology tell us that 95% of matter was not described in text-books yet

Two search strategies:

1) High energy physics to excite heavy degrees of freedom. No any evidence till now. LHC era has started on 10.09.08.

2) Low energy physics to produce Rare processes in view of huge statistics.

There are some rough edges of SM.

 $(g-2)_{\mu}$ is very famous example

 $\pi_0 \rightarrow e+e-is$ in the list of SM test after new exp. and theor. results

That's intriguing

Rare Pion Decay $\pi^0 \rightarrow e^+e^-$ **from KTeV** PRD (2007)



One of the simplest process for THEORY

Lowest order diagram

From KTeV E799-II EXPERIMENT at Fermilab experiment (1997-2007)

$$B_{\pi \to e^+ e^-}^{\text{KTeV}} = (7.49 \pm 0.29 \pm 0.25) \cdot 10^{-8}$$

The result is based on observation of 794 candidate $\pi_0 \rightarrow e+e$ - events using $K_L \rightarrow 3\pi_0$ as a source of tagged π_0 s.

Classical theory of $\pi^0 \rightarrow e+e-decay$



Drell (59'), Berman,Geffen (60'), Quigg,Jackson (68')

Bergstrom, et.al. (82') dispersion approach Savage, Luke, Wise (92') χ PT

$$R\left(\pi^{0} \to e^{+}e^{-}\right) = \frac{B\left(\pi^{0} \to e^{+}e^{-}\right)}{B\left(\pi^{0} \to \gamma\gamma\right)} = 2\beta\left(m_{\pi}^{2}\right)\left(\left(\frac{\alpha}{\pi}\frac{m_{e}}{m_{\pi}}\right)^{2}\right)\left(\left(\operatorname{Re}\mathcal{A}\left(m_{\pi}^{2}\right)\right)^{2} + \left(\operatorname{Im}\mathcal{A}\left(m_{\pi}^{2}\right)\right)^{2}\right)\right)$$
$$\beta\left(q^{2}\right) = \sqrt{1 - 4\frac{m_{e}^{2}}{q^{2}}}$$

$$\mathcal{A}(q^{2}) = \frac{2i}{q^{2}} \int \frac{d^{4}k}{\pi^{2}} \frac{q^{2}k^{2} - (qk)^{2}}{(k^{2} + i\varepsilon)\left((k - q)^{2} + i\varepsilon\right)\left((k - p)^{2} - m_{e}^{2} + i\varepsilon\right)} \frac{F_{\pi}\left(k^{2}, (k - q)^{2}\right)}{(k^{2} + i\varepsilon)\left((k - q)^{2} + i\varepsilon\right)\left((k - p)^{2} - m_{e}^{2} + i\varepsilon\right)} \frac{F_{\pi}\left(k^{2}, (k - q)^{2}\right)}{(k^{2} + i\varepsilon)\left((k - q)^{2} + i\varepsilon\right)\left((k - q)^{2} - m_{e}^{2} + i\varepsilon\right)}$$

$$\operatorname{Im} \mathcal{A}(q^{2}) = \frac{\pi}{2\beta(q^{2})} \ln\left(\frac{1-\beta(q^{2})}{1+\beta(q^{2})}\right)$$

The Imaginary part is Model Independent; Unitary limit



I. The Decay Amplitude in Soft limit $q^2 \rightarrow 0$ A.D., M.A.Ivanov **PRD 07**'

inverse moment of Pion Transition FF!!!

The accuracy is of order $O(m_e/m_o)^2$. Thus the amplitude is fully reconstructed!

A. E. Dorokhov and M. A. Ivanov, Phys. Rev. D 75 (2007) 114007

- [2] A. E. Dorokhov and M. A. Ivanov, JETP Lett. 87 (2008) 531
- [3] A. E. Dorokhov, E. A. Kuraev, Yu. M. Bystritskiy and M. Secansky, Eur. Phys. J. C 55 (2008) 193

II. CLEO data and Lower Bound on Branching

III. $F_{\pi}(t,t)$ general arguments

Let
$$F_{\pi}(t,t) = \frac{1}{1+t/s_1}$$
 then $\operatorname{Re} \mathcal{A}^{\operatorname{theory}}(q^2 = 0) = -\frac{3}{2} \ln\left(\frac{s_1}{m_e^2}\right) - \frac{5}{4}$
I. From $\left[-\frac{\partial F_{\pi}(t,t)}{\partial t}\right]_{t=0} = -2\frac{\partial F_{\pi}(t,0)}{\partial t}\Big|_{t=0}$ one has $s_1 = s_0/2$
2. From OPE QCD $\left[F_{\pi}^{OPE}(t,0)\Big|_{t\to\infty} = 8\pi^2 f_{\pi}^2 \frac{1}{t}, F_{\pi}^{OPE}(t,t)\Big|_{t\to\infty} = \frac{8\pi^2 f_{\pi}^2 1}{3} \frac{1}{t}\right]$ one has $s_1^{OPE} = s_0^{OPE}/3$
 $\left[F_{\pi}^{OPE}(t,t)\Big|_{t\to\infty} = \frac{8\pi^2 f_{\pi}^2 1}{3} \frac{1}{t}\right]$ one has $s_1^{OPE} = s_0^{OPE}/3$
It follows $\operatorname{Re} A^{\operatorname{theory}}(q^2 = 0) = -21.9 \pm 0.3$
 $B^{\operatorname{theory}}(\pi^0 \to e^+e^-) = (6.2 \pm 0.1) \cdot 10^{-8}$ 3.3σ below data!

It would required change of s_0 scale by factor more then 10!

Now it's intriguing!

B. *F*π(*t*,*t*) VMD parametrization (M.Knecht, A.Nyffeler, EPJC 01')

$$F_{gVMD}\left(s,t\right) = \frac{4\pi^{2}f_{\pi}^{2}}{3} \frac{\left(s+t\right)st - h_{2}st + h_{5}\left(s+t\right) + M_{V}^{4}M_{V1}^{4}h_{7}}{\left(M_{V}^{2}+s\right)\left(M_{V}^{2}+t\right)\left(M_{V1}^{2}+s\right)\left(M_{V1}^{2}+t\right)}$$

$$\left\langle r^2 \right\rangle_{\pi\gamma\gamma^*} = -6 \left. \frac{\partial F\left(t,0\right)}{\partial t} \right|_{t=0} = 6 \left(\frac{1}{M_V^2} + \frac{1}{M_{V1}^2} - \frac{h_5}{M_V^2 M_{V1}^2} \frac{3}{4\pi^2 f_\pi^2} \right) = 0.39 \quad \text{fm}^2$$

$$\operatorname{Re} \mathcal{A}_{gVMD} \left(q^2 = 0 \right) = -3 \ln \left(\frac{M_V}{m_e} \right) + \frac{1}{4} + \frac{3r}{(r-1)^2} - \frac{3}{2} \frac{3r-1}{(r-1)^3} \ln r + \frac{4\pi^2 f_{\pi}^2}{M_V^2 (r-1)^3} \left(\frac{h_2}{2M_V^2} \left((r+1) \ln r - 2 \left(r-1 \right) \right) - \left(1 + \frac{h_5}{M_{V1}^2 M_V^2} \right) \left(r^2 - 1 - 2r \ln r \right) \right) \neq -21.94$$

$$r = (M_{V1}/M_V)^2.$$

Nicely confirms general arguments! Parameters for LbL

Re
$$A^{\text{theory}} \left(q^2 = 0 \right) = -21.9 \pm 0.3$$

A. $F_{\pi}(t,t)$ **QCD** sum rules (V.Nesterenko, A.Radyushkin, YaF 83')

$$F_{\pi}^{\text{QCDsr}}\left(t,t\right) = 2 \int_{0}^{s_{0}^{\text{QCDsr}}} ds \int_{0}^{1} dx \frac{x\left(1-x\right)t^{2}}{\left[x\left(1-x\right)s+t\right]^{3}} + \text{v.c.},$$

From
$$-\frac{\partial F_{\pi}(t,t)}{\partial t}\Big|_{t=0} = -2\frac{\partial F_{\pi}(t,0)}{\partial t}\Big|_{t=0}$$

and
$$\left\langle r^2 \right\rangle_{\pi^{0\gamma^*\gamma^*}}^{\text{QCDsr}} = -6 \left. \frac{\partial F_{\pi}^{\text{QCDsr}}\left(t,t\right)}{\partial t} \right|_{t=0} = \frac{12}{s_0^{\text{QCDsr}}} \text{ one has } \left[s_0^{\text{QCDsr}} = s_0^{\text{CLEO}} \right]_{t=0}$$

Re
$$A^{\text{QCDsr}} \left(q^2 = 0\right) = -\frac{3}{2} \ln \left(\frac{s_0^{\text{CLEO}}}{m_e^2}\right) + \frac{1}{4} = -21.7 \pm 0.1,$$

 $s_1^{\text{QCDsr}} = \frac{s_0^{\text{QCDsr}}}{e}$

Nicely confirms general arguments!

$$\operatorname{Re} A^{\operatorname{theory}} \left(q^2 = 0 \right) = -21.9 \pm 0.3$$



What is next? It would be very desirable if Others will confirm KTeV result Also, Pion transition FF need to be more accurately measured.

Possible explanations of the effect

1) Radiative corrections

KTeV used in their analysis the results from Bergstrom 83'.

A.D., Kuraev, Bystritsky, Secansky (EJPC 08') confirmed Numerics.

2) Mass corrections (tiny)

A.D., M.Ivanov (ZhETPh Lett 08') Dispersion relations are corrected by small power corrections $(m_{\pi}/m_{\rho})^n$

3) New physics Kahn, Schmidt, Tait 07' Low mass dark matter particles

4) Experiment wrong Waiting for new results from KLOE, NA48, WASA@COSY, BESIII,...

Enhancement in Rare Pion Decays from a Model of MeV Dark Matter (Boehm&Fayet) was considered by Kahn, Schmitt and Tait (2007)



The anomalous 511 keV γ -ray signal from Galactic Center observed by INTEGRAL/SPI (2003) is naturally explained



FIG. 1. Dark matter annihilations into e^+e^- pairs [12,13]. The first diagram corresponds to the pair annihilation of spin- $\frac{1}{2}$ LDM particles χ (which may be self-conjugate, or not); and the second one to the case of spin-0 particles φ .

Rare decay $\pi 0 \rightarrow e+e-as$ a sensitive probe of light CP-odd Higgs in Next-to-Minimal SuperSymmetric Model (NMSSM) (Qin Chang, Ya-Dong Yang, 2008)



Figure 1: Relevant Feynman diagram within NMSSM.

They find the combined constraints from $Y \rightarrow \gamma A_1^0$, a_μ and $\pi 0 \rightarrow e+e-$ point to a very light A_1^0 with $m_{A01} \simeq 135 MeV$ and |Xd| = 0.10 + 0.08



Other $P \rightarrow I+I-$ decays

$$F_{\eta}^{\text{CLEO}}\left(k^{2}, q^{2}=0\right) = \frac{1}{1+k^{2}/s_{0\eta}^{\text{CLEO}}}, \qquad s_{0\eta}^{\text{CLEO}} = \left(774 \pm 29 \quad \text{MeV}\right)^{2},$$

TABLE II. Values of the branchings $B(P \rightarrow l^+ l^-)$ obtained in our approach and compared with the available experimental results.

В	Unitary bound	CLEO bound	CLEO + OPE	Experiment
$B(\pi^0 \rightarrow e^+ e^-) \times 10^8$	≥ 4.69	$\geq 5.85 \pm 0.03$	6.23 ± 0.09	7.49 ± 0.38 [1]
$B(\eta \rightarrow \mu^+ \mu^-) \times 10^6$	≥ 4.36	$\leq 6.23 \pm 0.12$	5.11 ± 0.20	5.8 ± 0.8 [7,32]
$B(\eta \rightarrow e^+e^-) \times 10^9$	≥ 1.78	$\geq 4.33 \pm 0.02$	4.60 ± 0.06	

 $B(\eta \rightarrow e^+e^-) \le 2.7 \times 10^{-5}$ Celsius/Wasa 2008

Summary

- The processes P → I+I- are good for test of SM. Radiative and mass corrections are well under control. New measurments of the transition form factors are welcome.
- 2) At present there is 3.3σ disagreemnt between SM and KTeV experiment for $\pi 0 \rightarrow e+e-$ CLOE, WASA@COSY, NA48 are interested in new measurments
- 3) If effect found persists it might be evidence for the SM extensions with low mass (10-100 MeV) particles (Dark Matter, NSSM)

Anomalous magnetic moment of muon

From BNL E821 experiment (1999-2006)

 $a_{\mu}^{\text{BNL}} = 11\ 659\ 208.0(6.3) \cdot 10^{-10}$

Standard Model

$$a_{\mu}^{SM} = a_{\mu}^{QED} + a_{\mu}^{EW} + a_{\mu}^{Hadr} = 11\ 659\ 178.5(6.1) \cdot 10^{-10}$$

predicts the result which is **3.4** σ below the experiment (since 2006)

