

Recent results on a solving Bethe–Salpeter equations for spinor particles

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Bethe-Salpater equation

$$\mathcal{G}(p) = ig^2 \int \frac{d^4 k}{(2\pi)^4} V(p, k) \Gamma(1) S(k_1) \mathcal{G}(k) \tilde{S}(k_2) \tilde{\Gamma}(2),$$

$$S(k) = \frac{\hat{k} + m}{k^2 - m^2 + i\epsilon}, \quad V(p, k) = \frac{1}{(p - k)^2 - \mu^2 + i\epsilon}.$$

$$\Gamma(1) = 1 \qquad \tilde{\Gamma}(2) = -1 \qquad \text{scalar}$$

$$\Gamma(1) = \gamma_5 \qquad \tilde{\Gamma}(2) = -\gamma_5 \qquad \text{pseudoscalar}$$

$$\Gamma(1) = \gamma_\mu - \frac{i\kappa}{2m} \sigma_{\mu\rho} q^\rho \qquad \tilde{\Gamma}(2) = \gamma_\nu + \frac{i\kappa}{2m} \sigma_{\nu\rho} q^\rho \qquad \text{vector}$$

Vertex function structure

$$\mathcal{G}(p_0, \mathbf{p}) = \sum_{\alpha} g_{\alpha}(p_0, |\mathbf{p}|) \Gamma_{\alpha}(\mathbf{p})$$

1S_0	${}^1S_0^{++}$	${}^1S_0^{--}$	${}^3P_0^e$	${}^3P_0^o$
${}^3S_1 - {}^3D_1$	${}^3S_1^{++}$	${}^3D_1^{++}$	${}^3S_1^{--}$	${}^3D_1^{--}$
	${}^3P_1^e$	${}^3P_1^o$	${}^1P_1^e$	${}^1P_1^o$

$$\Gamma_1(\mathbf{p}) = \frac{1}{\sqrt{16\pi}} \gamma_5 \quad \Gamma_2(\mathbf{p}) = \frac{1}{\sqrt{16\pi}} \gamma_0 \gamma_5$$

$$\Gamma_3(\mathbf{p}) = \frac{1}{\sqrt{16\pi}} \frac{(\mathbf{p}\boldsymbol{\gamma})}{|\mathbf{p}|} \gamma_0 \gamma_5 \quad \Gamma_4(\mathbf{p}) = \frac{1}{\sqrt{16\pi}} \frac{(\mathbf{p}\boldsymbol{\gamma})}{|\mathbf{p}|} \gamma_5$$

Expansion over hyperspherical harmonics

$$Z_{nlm}(\chi, \theta, \phi) = X_{nl}(\chi) Y_{lm}(\theta, \phi),$$

$$X_{nl}(\chi) = \sqrt{\frac{2^{2l+1} (n+1)(n-l)! l!^2}{\pi (n+l+1)!}} \sin^l \chi C_{n-l}^{l+1}(\cos \chi),$$

$$\frac{1}{(p-k)^2 + \mu^2} = 2\pi^2 \sum_{nlm} \frac{1}{n+1} V_n(\tilde{p}, \tilde{k}) Z_{nlm}(\chi_p, \theta_p, \phi_p) Z_{nlm}^*(\chi_k, \theta_k, \phi_k)$$

$$V_n(\tilde{p}, \tilde{k}) = \frac{4}{(\Lambda_+ + \Lambda_-)^2} \left(\frac{\Lambda_+ - \Lambda_-}{\Lambda_+ + \Lambda_-} \right)^n,$$

$$\Lambda_{\pm} = \sqrt{(\tilde{p} \pm \tilde{k})^2 + \mu^2}.$$

$$g_i(p_4, \mathbf{p}) = \sum_{i=1}^{\infty} g_i^j(\tilde{p}) Z_{jlm}(\chi_p, \theta_p, \phi_p)$$

Equations



$$g_i^j(\tilde{p}) = -g^2 \int_0^\infty \frac{d\tilde{k} \tilde{k}^3}{8\pi^2(2j-1)} V_{ij}(\tilde{p}, \tilde{k}) \sum_{n=1}^4 \sum_{m=1}^\infty A_{jm}^{in}(\tilde{k}) g_n^m(\tilde{k}),$$



$$X = g^2 AX,$$

$$X^T = \left([\{g_1^m(\tilde{k}_i)\}_{i=1}^{N_G}]_{m=1}^{M_{\max}}, [\{g_2^m(\tilde{k}_i)\}_{i=1}^{N_G}]_{m=1}^{M_{\max}}, \dots, [\{g_n^m(\tilde{k}_i)\}_{i=1}^{N_G}]_{m=1}^{M_{\max}} \right)$$

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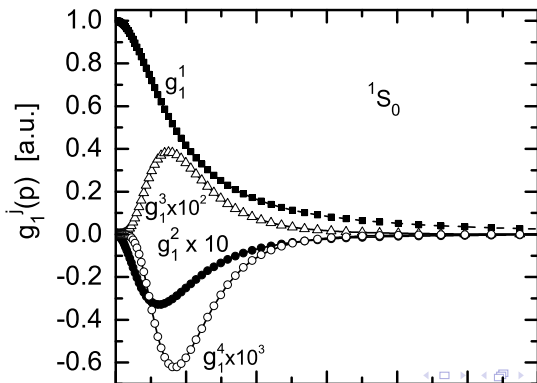
Convergence of the calculated values of mass M

$g^2 = 15$	$\mu = 0.15 \text{ GeV}/c^2$			$\mu = 0.5 \text{ GeV}/c^2$		
M_{\max}	$N_G = 32$	$N_G = 64$	$N_G = 96$	$N_G = 32$	$N_G = 64$	$N_G = 96$
1	1.9399	1.9399	1.9399	1.9984	1.9984	1.9984
2	1.9370	1.9370	1.9370	1.9982	1.9982	1.9982
3	1.9368	1.9368	1.9368	1.9982	1.9982	1.9982
4	1.9368	1.9368	1.9368	1.9982	1.9982	1.9982

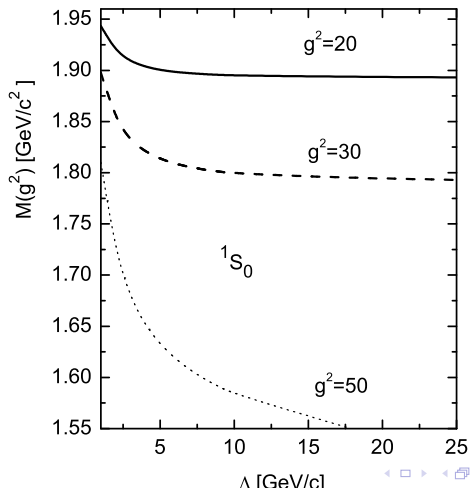
$g^2 = 30$	$\mu = 0.15 \text{ GeV}/c^2$			$\mu = 0.5 \text{ GeV}/c^2$		
M_{\max}	$N_G = 32$	$N_G = 64$	$N_G = 96$	$N_G = 32$	$N_G = 64$	$N_G = 96$
1	1.7932	1.7910	1.7905	1.9167	1.9142	1.9137
2	1.7897	1.7875	1.7871	1.9152	1.9127	1.9122
3	1.7896	1.7874	1.7870	1.9152	1.9127	1.9122
4	1.7896	1.7874	1.7870	1.9152	1.9127	1.9122

Results: components of g_1

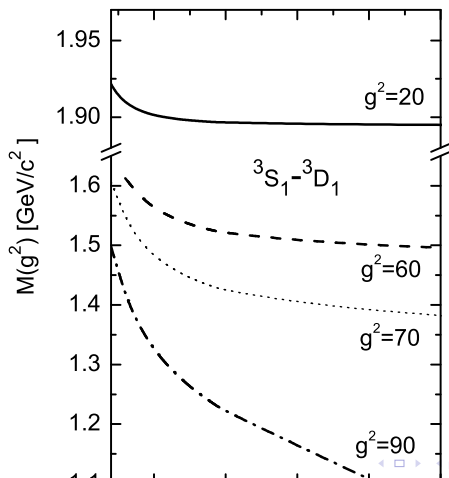
$$g^2 = 15 \quad M = 1.937 \text{ GeV}/c^2$$



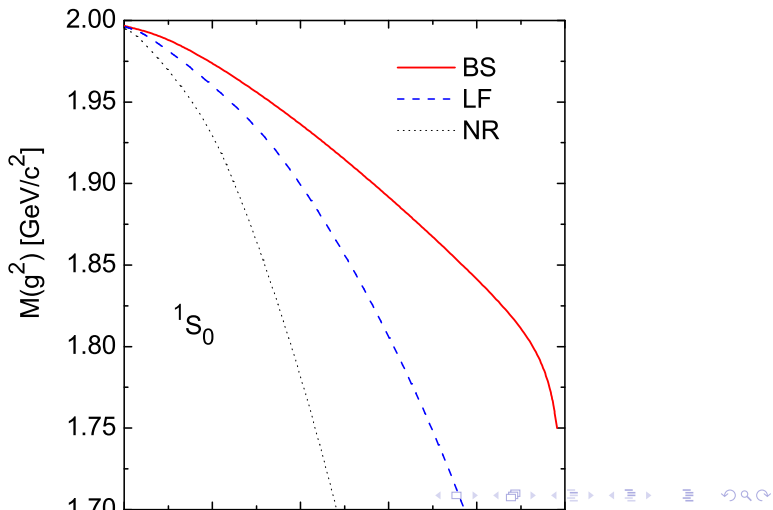
Results: stability of solutions



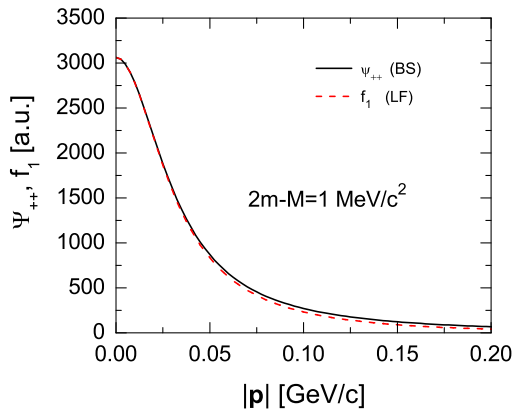
Results: stability of solutions



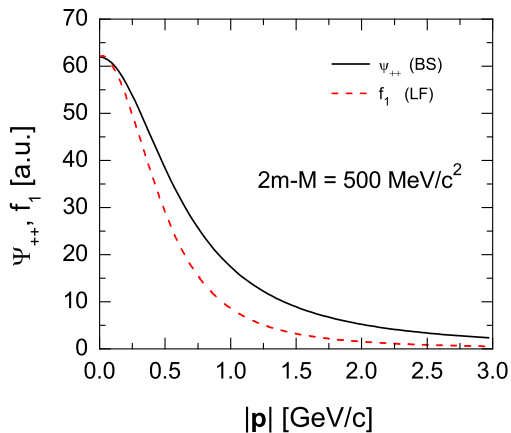
Results: g^2 vs M



Results: comparison to LF calculations



Results: comparison to LF calculations



Conclusions

Novel method of solving Bethe-Salpeter equation is proposed.

Advantages:

- fast convergence
- simple parameterization of solutions
- separable form of a kernel