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Study of Exotic States of Nuclear Matter using Lobachevsky Geometry. Separation of Narrow Resonances. Data Bases of Experiments with Bubble Chambers

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# Institute for Advanced Studies: Relativistic Nuclear Physics Research

- Educational programs
- Planning of experiments
- Verification of theoretical models
- Search of exotic states of nuclear matter
- Lobachevsky space
- Grassman and Clifford vector algebras
- Image recognition



Figure 3: Three-body weak decay of the heavy, stable, positively charged dibaryon fits possible for two hypotheses:  $H^+ \rightarrow p + \gamma + \Lambda^0$  and  $H^+ \rightarrow p + \pi^0 + \Lambda^0$ .



Figure 4: Two body weak decay  $\Xi^-$  hyperon  $\to \pi^- + \Lambda$ : a)first event;b) second event.



# Data Bases of Experiments with Bubble Chambers

- 4.2 GeV p, d, He, C, Ta +CH;
- 4.2 GeV p, C + Ta;
- 10 GeV p + CH;
- 40 GeV π + CH (Protvino);
- 5.2, 3.8, 2.2, 1.7, 1.4 GeV/c n + p;
- 2.2, 1.7, 1.4 GeV/c d + p.





# Collaboration

- Kazan State University
- Lobachevsky Institute of Mathematics and Mechanics
- University of Bucharest
- Maritime University of Constanta
- Fock Physics Institute, St. Petersburg State University
- Institute of Nuclear Physics, Moscow State University
- Samara State University
- Perm State University
- CERN



**Angle of Parallelism** 

# Proton distribution for two angular intervals in p(10GeV/c)+C



A. A. Baldin, E. G. Baldina, E. N. Kladnitskaya, O. V. Rogachevskii, Phys.Part.Nucl.Lett., vol. 1, no. 4, 7-16 (2004).

# **n+p->***π*- 5GeV



# **n+p->***π*- 5GeV









 $\Delta_{12}^3 = 2\Pi_L \mathbf{\Phi}_3 - \alpha_3$ 







 $defect = \pi - \alpha_1 - \alpha_2 - \alpha_3$ 

Normalized distributions of defects of triangles formed by all combinations of protons and all combinations of mesons registered in p(10GeV/c)+C



 $\pi$  simNote, that the modeladequately reproducesinclusive spectra ofboth protons and  $\pi$  -mesons.The distribution of triosof π-mesons, however,differs noticeably from

experimental data.





## Volumes of tetrahedra



$$\begin{aligned}
\alpha_1 &= \exp(i\overline{AB}); \quad \beta_1 &= \exp(i\overline{BA}) \\
\alpha_2 &= \exp(i\overline{BC}); \quad \beta_2 &= \exp(i\overline{CB}) \\
\alpha_3 &= \exp(i\overline{CD}); \quad \beta_3 &= \exp(i\overline{DC}) \\
\alpha_4 &= \exp(i\overline{DA}); \quad \beta_4 &= \exp(i\overline{AD}),
\end{aligned}$$
(17)

we obtain

$$0 = \frac{1}{\alpha_{1}\alpha_{2}\alpha_{3}\alpha_{4}} - \beta_{1}\beta_{2}\beta_{3}\beta_{4} + z^{2}\left(-\frac{\alpha_{1}}{\alpha_{2}\alpha_{3}\alpha_{4}} - \frac{\alpha_{2}}{\alpha_{1}\alpha_{3}\alpha_{4}} - \frac{\alpha_{3}}{\alpha_{1}\alpha_{2}\alpha_{4}} - \frac{\alpha_{4}}{\alpha_{1}\alpha_{2}\alpha_{3}}\right) + \frac{\beta_{1}\beta_{2}\beta_{3}}{\beta_{4}} + \frac{\beta_{1}\beta_{2}\beta_{4}}{\beta_{3}} + \frac{\beta_{1}\beta_{3}\beta_{4}}{\beta_{2}} + \frac{\beta_{2}\beta_{3}\beta_{4}}{\beta_{1}}\right) + z^{4}\left(\frac{\alpha_{1}\alpha_{2}}{\alpha_{3}\alpha_{4}} + \frac{\alpha_{1}\alpha_{3}}{\alpha_{2}\alpha_{4}} + \frac{\alpha_{2}\alpha_{3}}{\alpha_{1}\alpha_{4}} + \frac{\alpha_{2}\alpha_{4}}{\alpha_{2}\alpha_{3}} + \frac{\alpha_{2}\alpha_{4}}{\alpha_{1}\alpha_{3}} + \frac{\alpha_{3}\alpha_{4}}{\alpha_{1}\alpha_{2}} - \frac{\beta_{1}\beta_{2}}{\beta_{3}\beta_{4}} - \frac{\beta_{1}\beta_{3}}{\beta_{2}\beta_{4}} - \frac{\beta_{2}\beta_{3}}{\beta_{1}\beta_{4}} - \frac{\beta_{2}\beta_{4}}{\beta_{2}\beta_{3}} - \frac{\beta_{2}\beta_{4}}{\beta_{1}\beta_{3}\beta_{4}} - \frac{\beta_{3}\beta_{4}}{\beta_{1}\beta_{2}\beta_{3}} - \frac{\beta_{3}\beta_{4}}{\beta_{1}\beta_{2}\beta_{3}} - \frac{\beta_{3}\beta_{4}}{\beta_{1}\beta_{2}\beta_{3}} - \frac{\alpha_{1}\alpha_{2}\alpha_{4}}{\alpha_{4}} - \frac{\alpha_{1}\alpha_{2}\alpha_{4}}{\alpha_{2}} - \frac{\alpha_{1}\alpha_{3}\alpha_{4}}{\alpha_{2}} - \frac{\alpha_{2}\alpha_{3}\alpha_{4}}{\alpha_{1}} - \frac{1}{\beta_{1}\beta_{2}\beta_{3}\beta_{4}}\right).$$
(18)

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#### The Regge symmetry is a scissors congruence in hyperbolic space

Yana Mohanty

$$\overline{AB} = \frac{A + A' + 2B'}{4}; \ \overline{BA} = \frac{2\pi + A - A' + 2C'}{4} 
\overline{BC} = \frac{A + A' - 2B'}{4}; \ \overline{CB} = \frac{2\pi - A + A' - 2C}{4} 
\overline{CD} = \frac{-A - A' - 2B}{4}; \ \overline{DC} = \frac{2\pi - A + A' + 2C}{4} 
\overline{DA} = \frac{-A - A' + 2B}{4}; \ \overline{AD} = \frac{2\pi + A - A' - 2C'}{4}.$$
(21)

$$\begin{split} (T) &= \mathcal{I}(\overline{AB} + \arg z_{-}) + \mathcal{I}(\overline{BA} - \arg z_{-}) + \mathcal{I}(\overline{BC} + \arg z_{-}) \\ &+ \mathcal{I}(\overline{CB} - \arg z_{-}) + \mathcal{I}(\overline{CD} + \arg z_{-}) + \mathcal{I}(\overline{DC} - \arg z_{-}) \\ &+ \mathcal{I}(\overline{DA} + \arg z_{-}) + \mathcal{I}(\overline{AD} - \arg z_{-}) + \frac{1}{2} [\mathcal{I}(\frac{\pi + A - B - C}{2}) \\ &- \mathcal{I}(\frac{\pi + B - A - C}{2}) - \mathcal{I}(\frac{\pi + C - A - B}{2}) + \mathcal{I}(\frac{\pi + B' - A' - C}{2}) \\ &+ \mathcal{I}(\frac{\pi + A + B + C}{2}) + \mathcal{I}(\frac{\pi + C - A' - B'}{2}) + \mathcal{I}(\frac{\pi - A' + B' + C}{2}) \\ &- \mathcal{I}(\frac{\pi + A' + B' + C}{2}) + \mathcal{I}(\frac{\pi + A' - B - C'}{2} - \mathcal{I}(\frac{\pi + A + B' + C'}{2}) \\ &- \mathcal{I}(\frac{\pi + A - B' - C'}{2}) + \mathcal{I}(\frac{\pi + B' - A - C'}{2}) - \mathcal{I}(\frac{\pi - A' - B + C'}{2}) \\ &\mathcal{I}(\frac{\pi + A' - B + C'}{2}) + \mathcal{I}(\frac{\pi + A' + B + C'}{2}) + \mathcal{I}(\frac{\pi - A - B' + C'}{2})], \end{split}$$
(23)

where the quantities with bars are given by (21), and  $z_{-}$  is the solution of the quadratic equation (18) with the negative square root.

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