

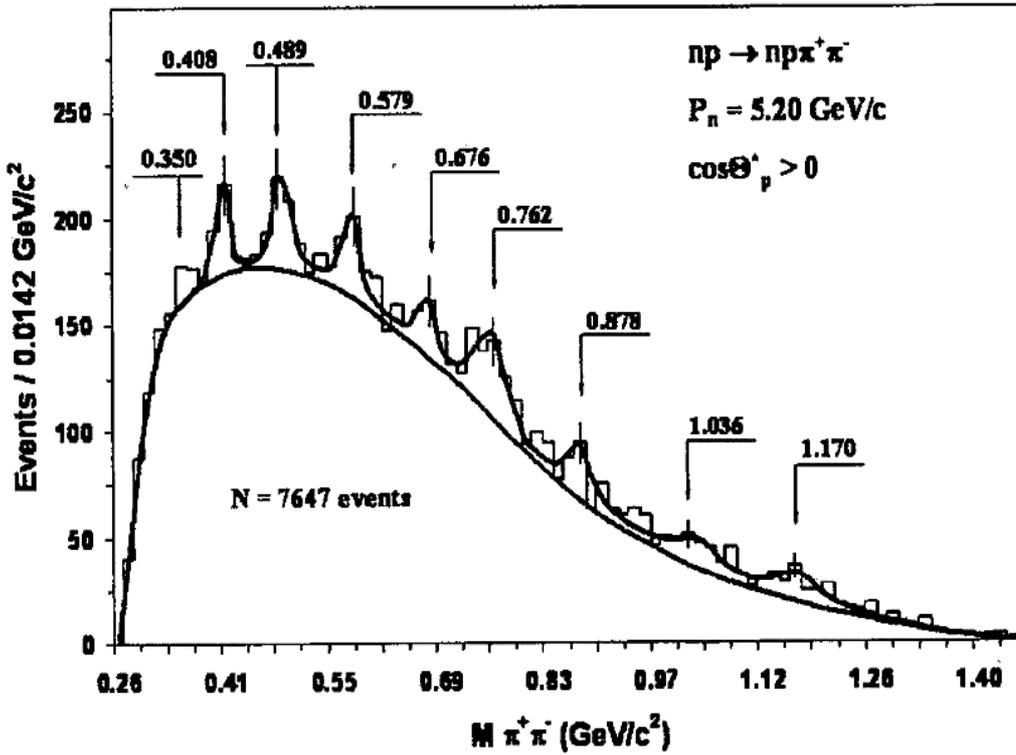
**EXPLANATION OF UNUSUAL  
np  $\rightarrow$  np $\pi^+\pi^-$  AND np  $\rightarrow$  np $K^+K^-$   
REACTIONS AT  $P_n=5.2\text{GeV}/c$   
BY MODEL OF ROTARY  
TWO-NUCLEONS SYSTEM**

**G.M.Amalsky ( PNPI, Gatchina)**

# $np \rightarrow np\pi^+\pi^-$ and $npK^+K^-$ at $P_n=5.2$ GeV/c

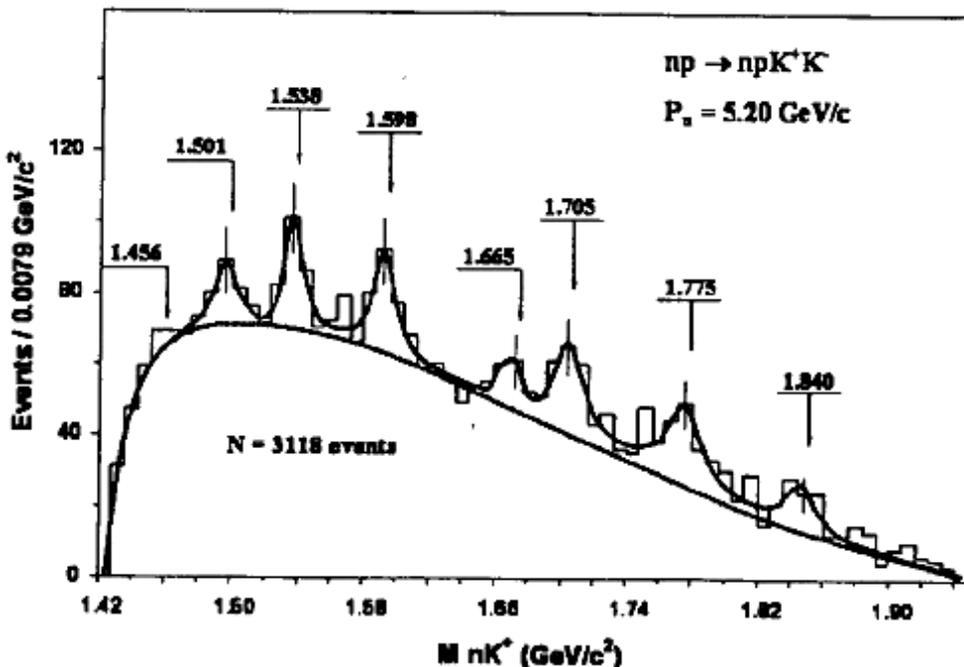
(JINR) Yu.A.Troyan, A.V.Beljaev, A.Yu.Troyan, E.B.Plekhanov, A.P.Jerusalimow, S.G.Arakelian

Proc. XVIII ISHEPP, v.1, p.114, and v.2, p.186



$np\pi^+\pi^-$   
channel

$$\tau = 0$$



$npK^+K^-$   
channel

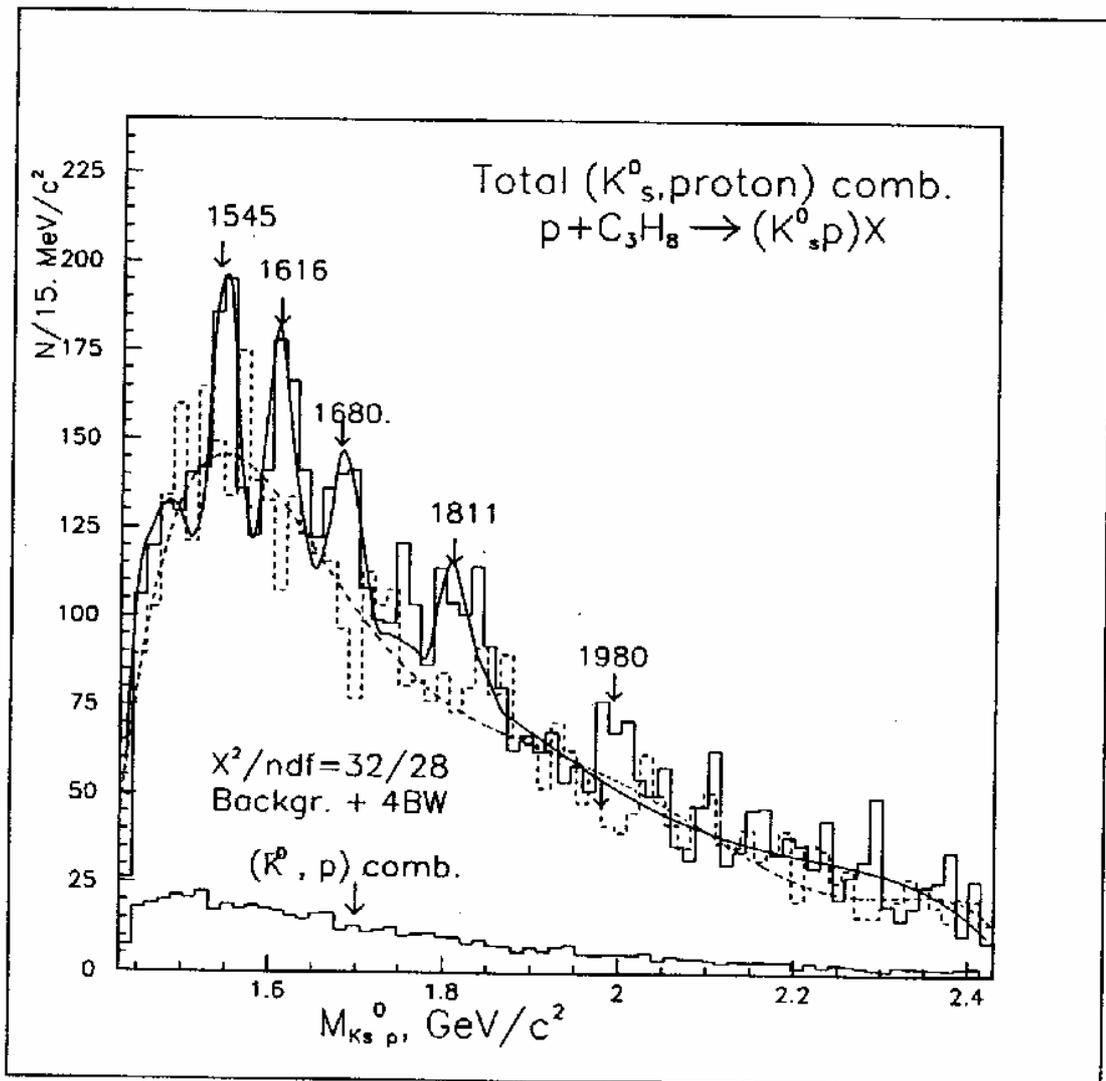
$$\tau = 0$$

$p + C_3H_8 \rightarrow pK^0_s + X$  at  $P_p = 10.0$  GeV/c

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Proc. XVII ISHEPP, 2005, v.II, p. 26.

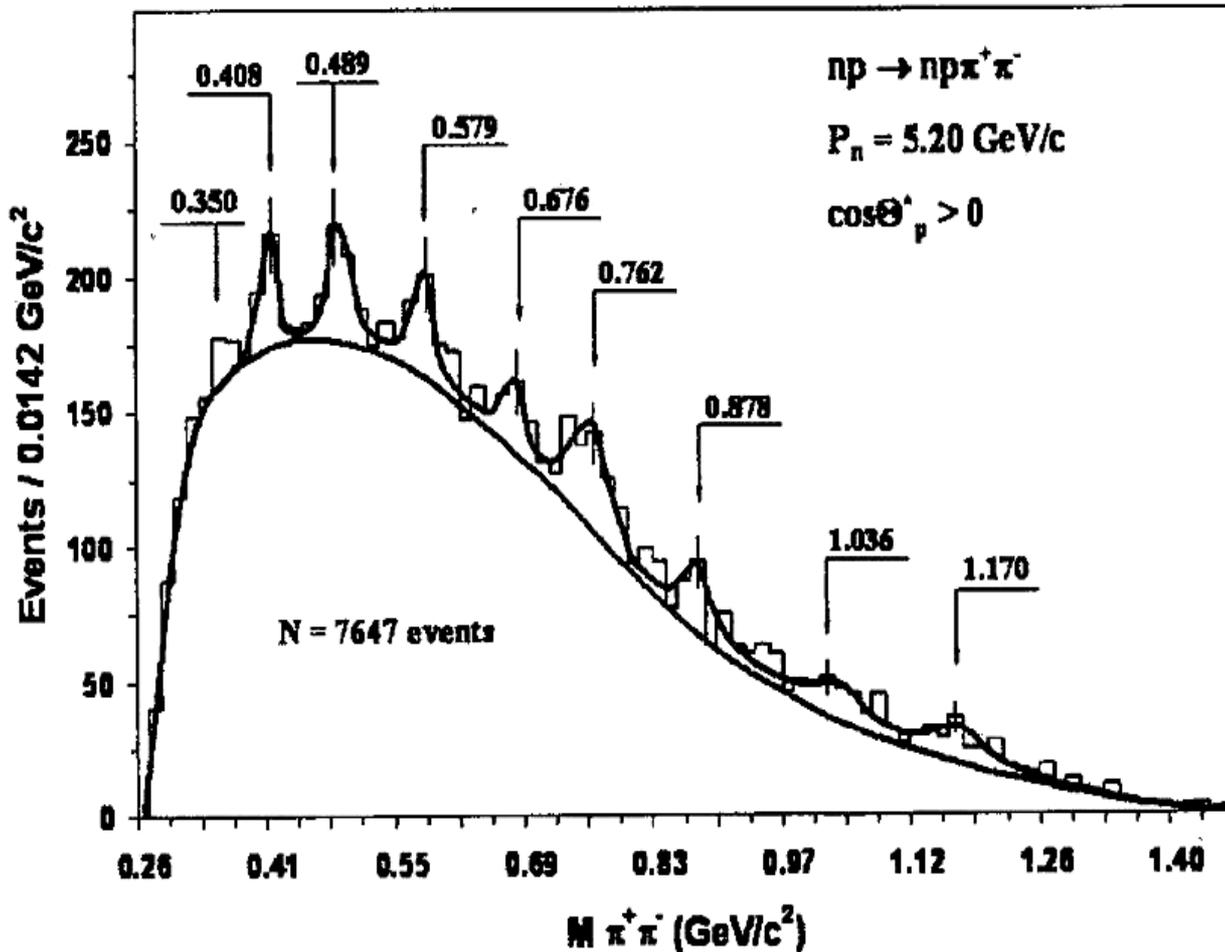


Similar peaks, looks as  $pp \rightarrow pK^0_s + X$   
8 free protons and 18 from 3 nuclei  $C^{12}$ .

$N_p^{(eff)} = 8 + k^{(i)}18$ ; for  $i=1-3$   $k^{(i)} \sim 2/3$

# PROPERTIES OF SELECTED EVENTS

$$np\pi^+\pi^-$$



events  $|P'_{\pi^+} + P'_{\pi^-}| \ll P'_{\pi^+} + P'_{\pi^-}$  are selected.

$\Delta P'_n > 1.5 \text{ GeV}/c$  : isospin (n+p)  $\tau = 0$ .

**Peaks  $\pi^+\pi^- J^{PC} = 0^{++}$  are not resonances**

Data have nothing explanation ~10 years.

# PROPERTIES OF $np\pi^+\pi^-$ EVENTS

Isotropic distribution:  $\pi^+\pi^-$  pairs are not emitted by moving nucleons.

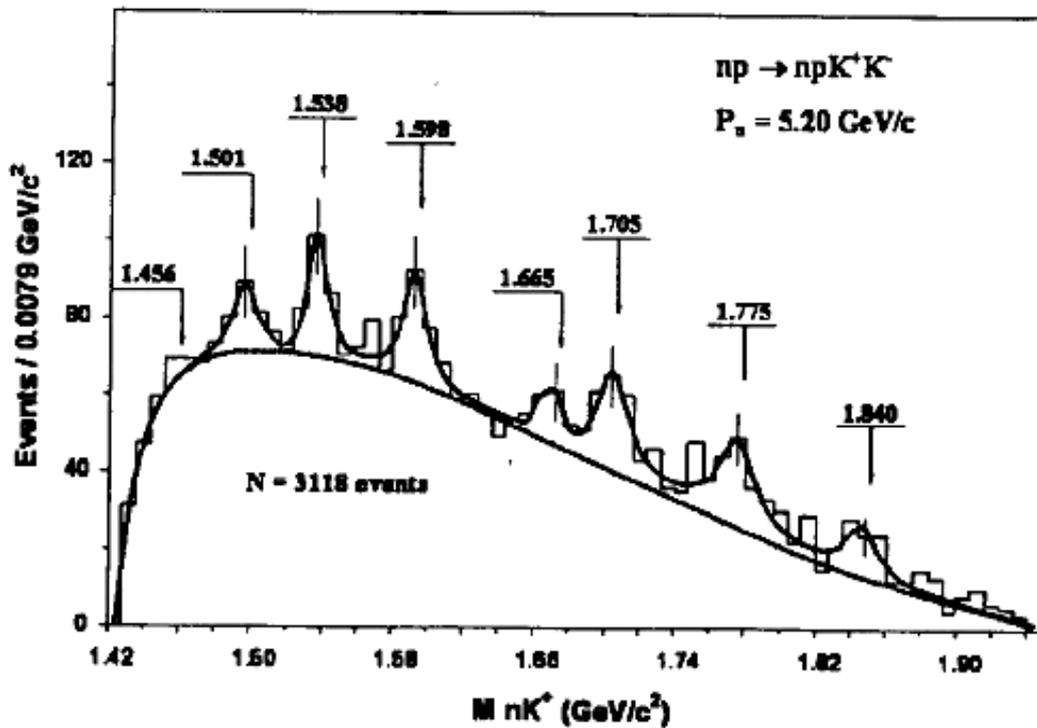
$L_{\pi\pi}=0$  limited sphere  $\Delta r < h/\Delta P_{\pi\pi} \sim 0.2\text{fm}$   
and time  $\Delta t < h/c\Delta P_{\pi\pi} \sim 10^{-25}$  sec  
of formation of pairs with  $\Delta P_{\pi\pi} \sim \text{GeV}/c$ .

Mean cross-section of peaks

$$\langle \sigma_{\pi\pi}^{(i)} \rangle = 124 (18) \text{ mkb} / 8 = 16 (3) \text{ mkb}$$

and  $\sigma_{\pi\pi} \sim 1 \text{ mb}$  of all  $np\pi^+\pi^- 0^{++}$  events  
is more 3 % of all inelastic events  $np$ ,  
despite of limited  $\Delta r$  and  $\Delta t$  of formation.  
 $np\pi^+\pi^-$  with  $L_{\pi\pi} > 0$  are suppressed: **all**  
**inel. event  $np\pi^+\pi^-$  happen in small  $\Delta r$ .**

# $Np \rightarrow npK^+K^-$ at $P_n=5.2 \text{ GeV}/c$



Mean difference of neighbor masses

$$\langle M_{nK}^{(i+1)} - M_{nK}^{(i)} \rangle = 56.5 (1.2) \text{ MeV}/c^2$$

is near to half of mean difference

$$\langle M_{\pi\pi}^{(i+1)} - M_{\pi\pi}^{(i)} \rangle = 109 (2) \text{ MeV}/c^2$$

in first spectrum: the same origin of peaks may be assumed --- quantization of angle momentum  $J=L+S$  of short state of two-nucleons system, formed in collision.

# Short rotary states of (np) system

Peaks  $M_{\pi\pi}^{(i)}$  can be explained  
by quantization of energy

$$\Delta E(J) = (2J+1)V, \text{ emitted}$$

at transition (n+p) system from

$L_1 = J+1, S_z = -1$  to  $L_2 = J-1, S_z = 1,$

and transformed into  $\pi^+\pi^-$  mass

(or divided between  $nK^+$  and  $pK^-$ )

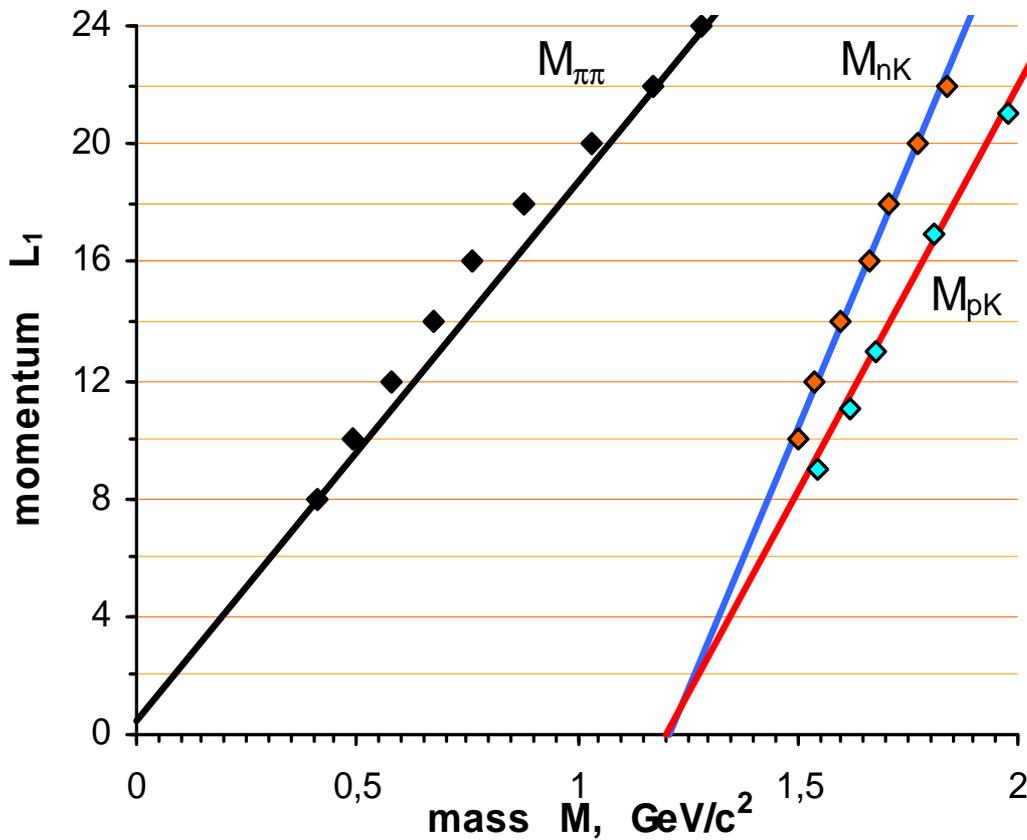
In states with  $S=1$  and  $\tau=0$

moments  $L^{(i)}$  must be even, so

$$L_1^{(i+1)} - L_1^{(i)} = 2.$$

**Short state with moment  $L_1 = bP_n$   
forms in laboratory system.**

# EXP. SPECTRA AND ITS DESCRIPTION



(black line)  $M_{\pi\pi}^{(i)} = 2V (L_1^{(i)} - 1/2)/c^2$  ,

(blue line)  $M_{nK}^{(k)} = V (L_1^{(k)} - 1/2)/c^2 + m + 2m_\pi$

(red line)  $M_{pK}^{(i)} = V' (L_1^{(i)} - 1/2)/c^2 + m + 2m_\pi$

$$V = \frac{h^2}{6mR_0^2}, \quad V' = \frac{h^2}{2mR_0^2} \frac{(m + 4m_\pi)}{(3m + 4m_\pi)}$$

$R_0 = 0.50 \text{ fm}$ ,  $m = m$  is nucleon mass

# DATA CAN NOT BE DESCRIBED AS COLLISIONS IN S.C.M.

$P_n = 5.2 \text{ GeV}/c$  and  $E = 6.22 \text{ GeV}$  in lab.sys.

correspond to  $E' = 3.42 \text{ GeV}$  and

$P'_n = P'_p = 1.43 \text{ GeV}/c$  in s.c.m. (np)'.

Transitions with  $\Delta L = 2$  and emitted  $\Delta E(J)$  in s.c.m. are impossible.

Spectra confirm to momentum

$L = bP_n$  in fixed laboratory system.

Next sequence of events may be:

1) collision with  $L = bP_n \rightarrow$  2) rotation with

$J = L + S$ , transition  $\Delta L = 2$  and definition of

energy  $\Delta E(J) \rightarrow$  3) decay of rotation and

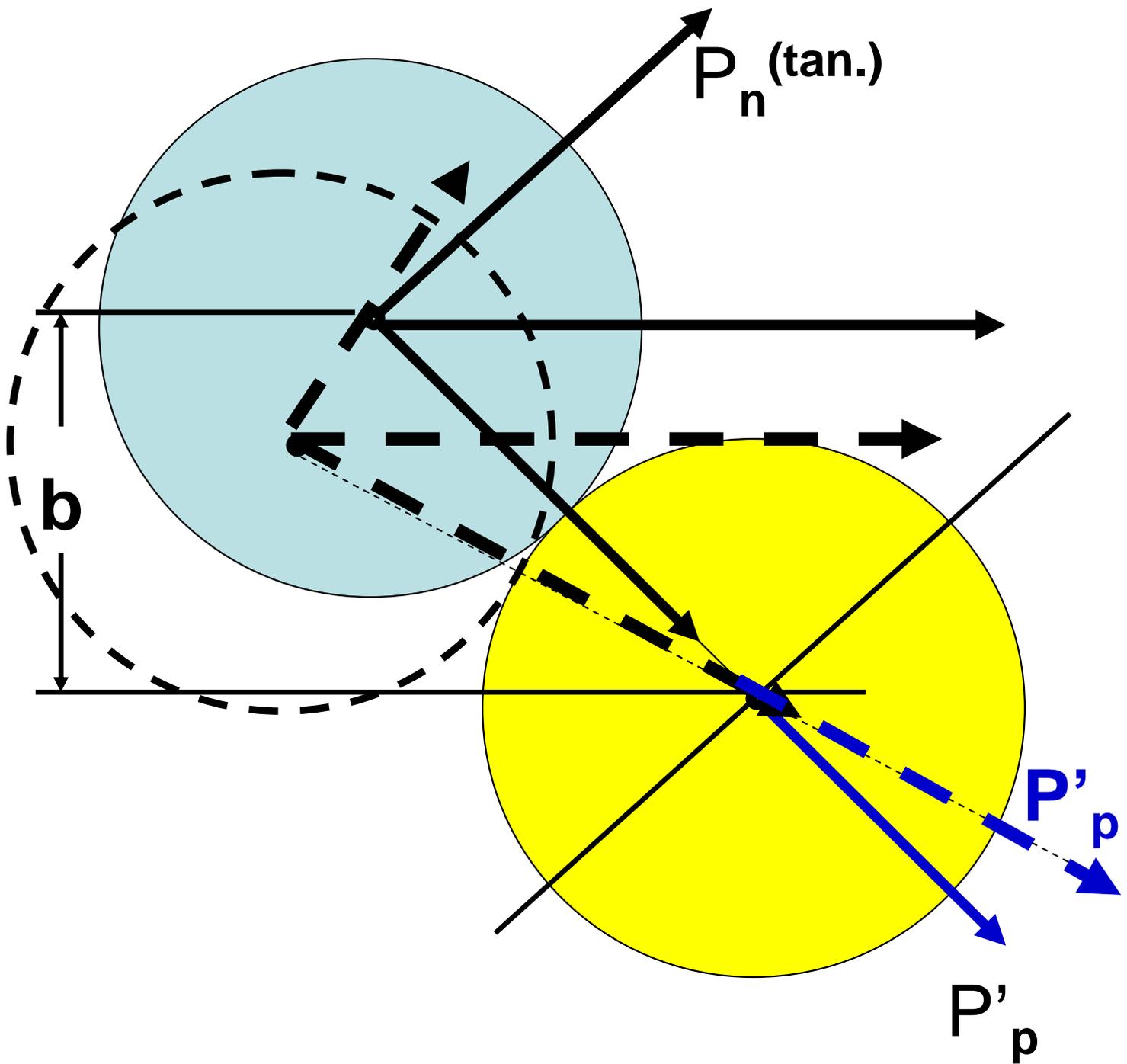
formation of observed finite particles with

$\Sigma L_f = L - 2$  in laboratory system.

Condition  $\Sigma L_f = L - 2$  and energy  $\Delta E(J)$  define

observed correlation of finite particles.

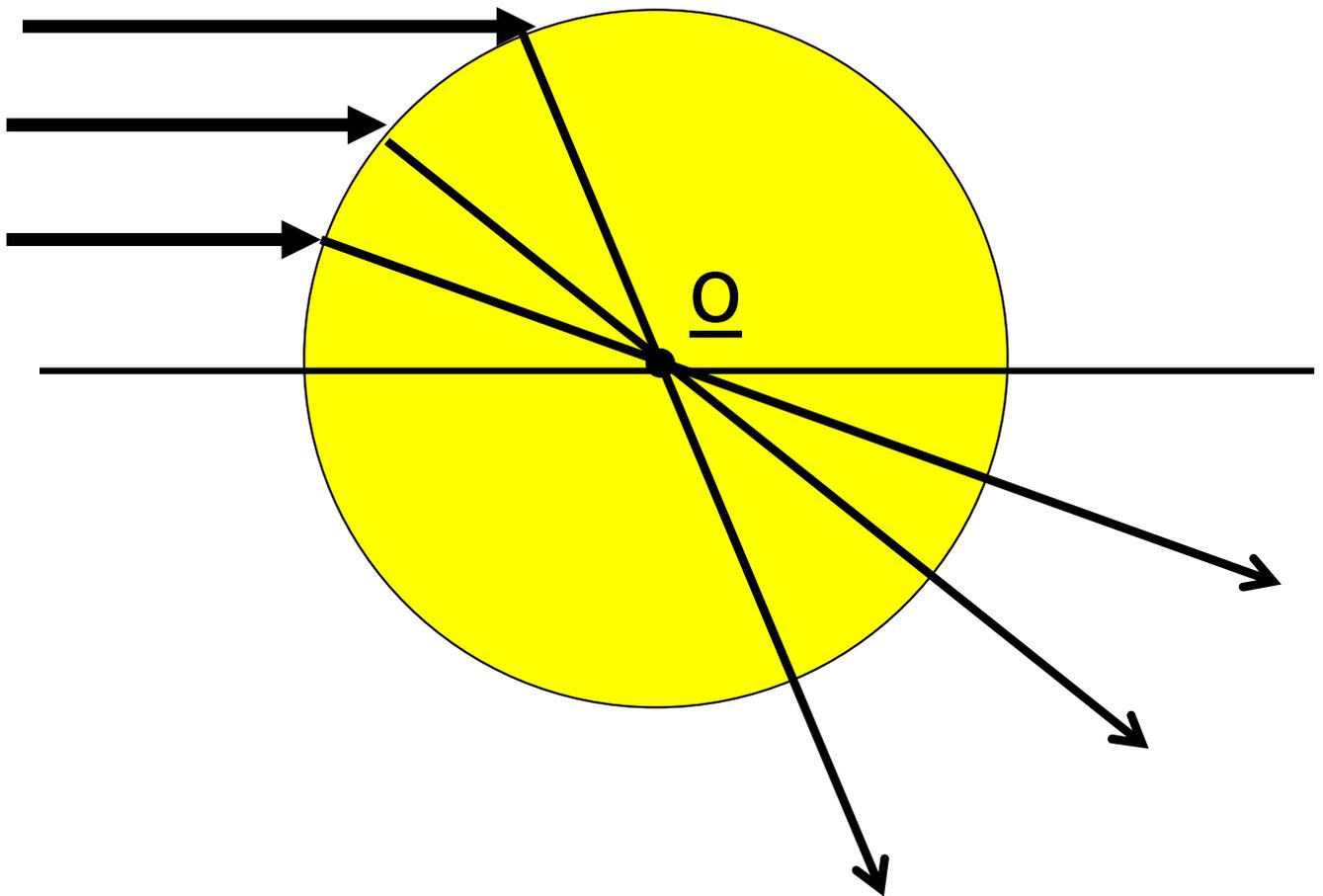
# HARD BALLS SCATTERING IN LAB.S.



In laboratory system:  $L_z = bP_n = 26$  may be,  
 hard balls: explanation  $L'_z = L_z = 2R_0 P_n(\text{tan.})$ ,  
 observed in spectra.

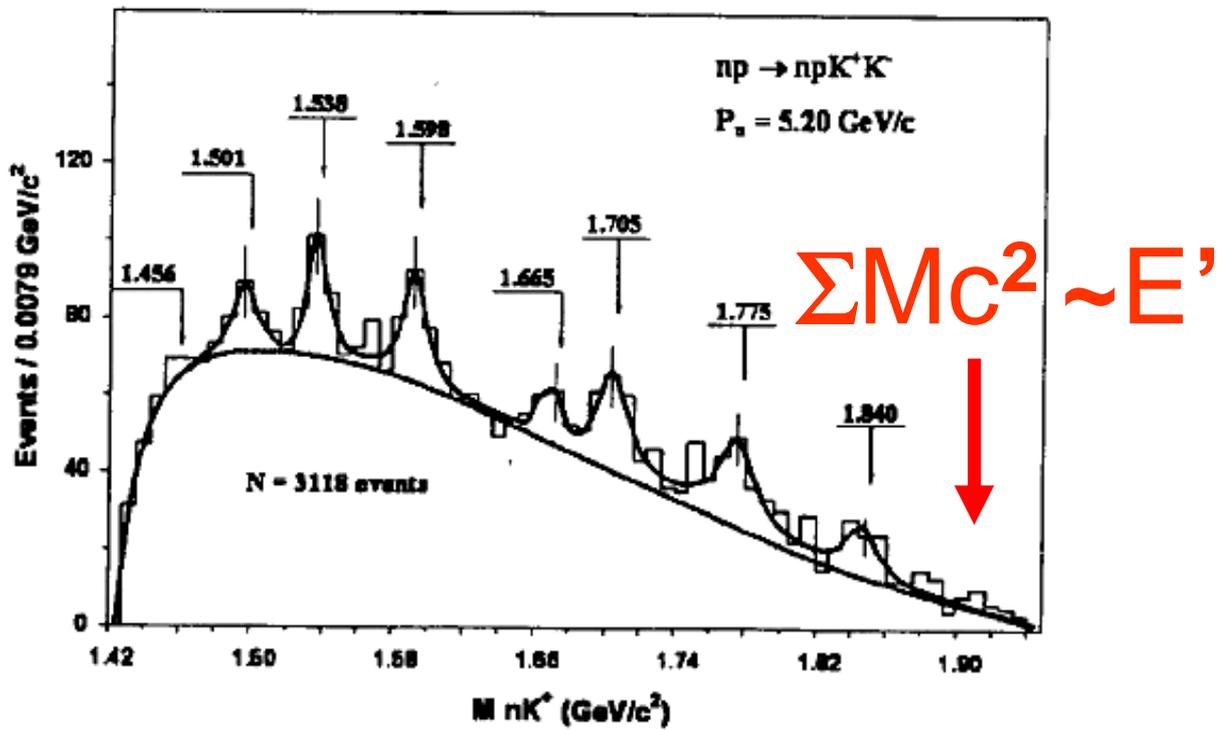
in lab.system: moment  $J = L + S$  can be kept

# QUANTIZATION OF $P'_p$ DIRECTION (IN LAB. SYSTEM) EXPLAINS $L'=L$



Radial movement of center of masses gives zero contribution in momentum of system  $L$  relative to point  $O$  in lab. sys., so  $L$  keeps as moment of rotation of two-nucleons system  $L'=L$  in its s.c.m..

# $Np \rightarrow npK^+K^-$ at $P_n=5.2 \text{ GeV}/c$



**Data are incompatible with invariance:**

in events  $c^2(M_{nK} + m_p + m_K) \sim E' = 3.42 \text{ GeV}$

momentum in s.c.m.  $L' \sim 0$  and  $L \sim 0$  in

such events of oncoming n and p



with energy  $E = E'$  in laboratory system.

# Remarks of **invariance**

Events  $np\pi^+\pi^-$  with large  $L > 20$  and

$$c^2(M_{\pi\pi} + m_p + m_n) \sim E' = 3.42 \text{ GeV}$$

are impossible in collisions of oncoming  $n$  and  $p$  with total energy  $E'$  too because of law of conservation of momentum  $L$ .

Absence of **invariance** of s.c.m. and of laboratory system means: **relativity**, which claims existence of invariance, is mistaken.

It may be proved straight.

**QUANTIZATION  $J=L+S$  AND  
PROPOSED DESCRIPTION OF DATA  
ARE INCOMPATIBLE WITH  
RELATIVITY TOO**

Relativistic Lorentz ratio

$$J'_z = J_z(1 - v^2/c^2)^{1/2}$$

of transverse moments in moving  
and in fixed systems means

**nonconservation of momentum**

**$L_z = bP_n$  and forbids quantization**

**of angle momentum in general.**

**Useful for description of data**

**model with quantized  $L'=L$  must**

**be incompatible with relativity.**

**(Relativity is incompatible**

**and with other laws of nature.)**

# Further plan

## 1. PHENOMENOLOGICAL ANALYSIS:

hypothesis of nonequilibrium interaction  
+ “hard ball” assumption of nucleons,  
**definition of empirical parameters,**  
opposition to well-known conceptions.

## 2. SOLID-BODY ROTATION MODEL:

gives ratio  $V_{sL} = (I_s / I_L) h^2 / (I_s + I_L)$  for  
potential energy  $\Delta U^{(s-b)}_{12} = V_{sL} (2L_1 - 1)$ ,  
emitted at transition  $L_1 \rightarrow L_1 - 2$ .

## 3. IMPERATIV SLOWING-DOWN OF NONEQUILIBRIUM ROTATION:

explanation of 100% transition probability  
and preliminary transition  $S_0=0$  into  $S_z=-1$

## 4. CONCLUSIONS:

conservation of moment  $J$  in lab. system defines  
impulses of final particles and mesons formation.

in free time: other grounds for model assumptions

# NUCLEONS AS “BLACK BALLS”

Near equidistant peaks show that momentum of inertia  $I$  is independent of momentum  $L$

It treats: in short state with momentum  $J=L+S$  nucleons interact as hard balls with radius  $R_0$ , independent of  $L$  and  $b$ .

**Nucleon is stable distribution of probabilities of possible events with any particles, which form it.**

# **ABSENCE OF RETARDING EFFECT**

Nucleon is probability distribution of some unknown possible inner events in its volume.

**Probability** of possible event **are non-material abstract possibility** to detect it with using other events.

Therefore all **probabilities of possible (in future!) events** may be changed **simultaneously** in agree with changing conditions, which can be defined in point of contact of distributions (nucleons).

# DEFINITION OF EMPIRICAL PARAMETERS OF SPECTRA

Ratio  $c^2 M_{\pi\pi}^{(i)} = 2V_{\pi\pi}(L_1^{(i)} - 1/2)$  and  $\tau = 0$

define values  $L_1^{(i)}$  and  $L_1^{(\max)}$  by using

$$V_{\pi\pi} = c^2(M_{\pi\pi}^{(8)} - M_{\pi\pi}^{(1)})/28 = 27.2(0.5) \text{ MeV}$$

Values  $V_{nK} = 28.2(0.6) \text{ MeV}$  in emp.

ratio  $c^2 M_{nK}^{(k)} = V_{nK}(L_1^{(k)} - 1/2) + M_0$

defined as  $V_{nK} = c^2(M_{\pi\pi}^{(7)} - M_{\pi\pi}^{(1)})/12$ .

Using of  $V_{nK}$ ,  $M_{nK}^{(\max)}$  and  $L_1^{(\max)}$

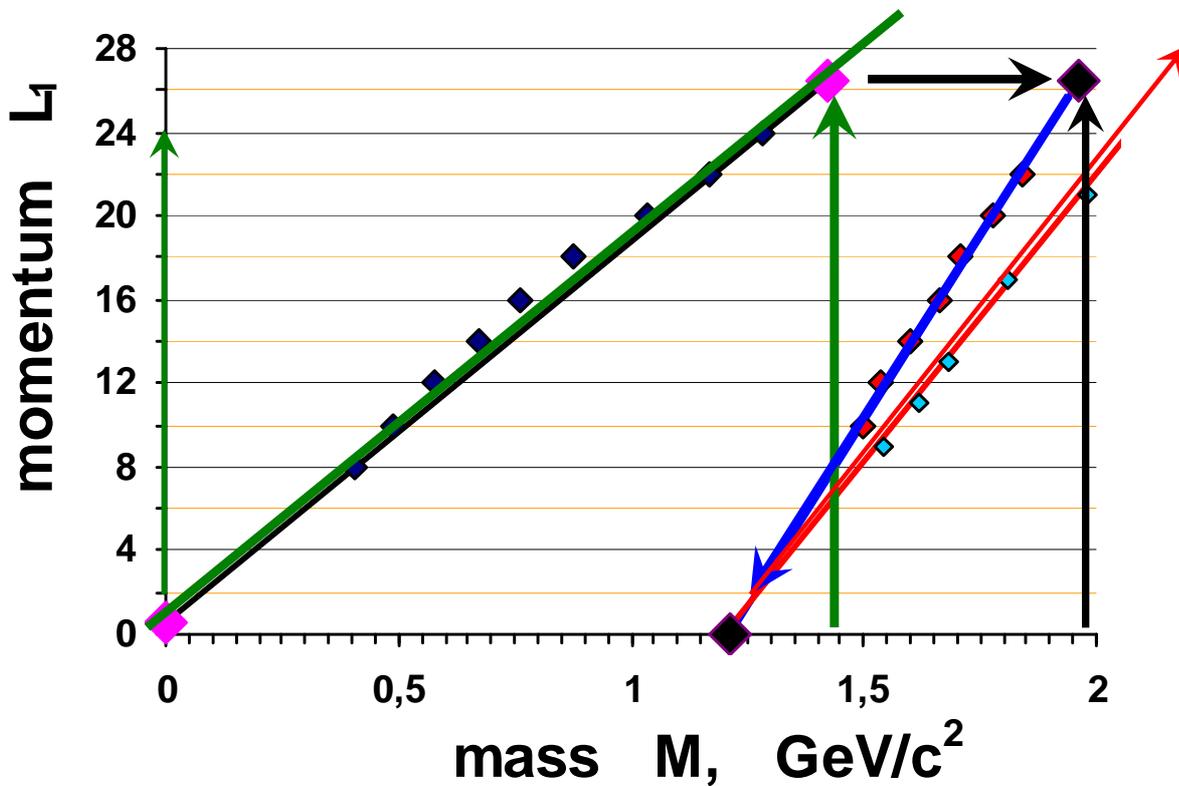
gives values of  $L_1^{(k)}$  and

$$M_0 = M_{nK}^{(\max)} - V_{nK}(L_1^{(\max)} - 1/2)/c^2 =$$

$$= 1230(20) \text{ MeV}/c^2 \cong m + 2m_\pi .$$

Data define all parameters in ratios.

# SCHEME OF DEFINITION OF EMPIRICAL PARAMETERS



- $M_{\pi\pi}^{(8)}, \dots, M_{\pi\pi}^{(1)}$  fix green line, slope  $V_{\pi\pi}$
- Its crossing with green line  $M_{\pi\pi}^{(\max)}$  gives  $L_1^{(\max)}=26.4$  and fixes blue line with slope  $V_{nK}$  in point  $(L_1^{(\max)}, M_{nK}^{(\max)})$ .
- Crossing of blue line with axis  $M$  gives  $M_0=1230$  (20)  $\text{MeV}/c^2 = m+2m_\pi$ .

Here red line  $M_{PK}$  begins with slope  $V_{PK0}$ .

# DIFFERENT INITIAL MOMENTS IN $npK^+K^-$ AND $np\pi^+\pi^-$ CHANNELS

In  $npK^+K^-$  events initial moments:

$$\text{odd } L_0^{(i)} = L_1^{(i)} - 1 \text{ and } S_0 = 0.$$

Preliminary transition happens

$$L_0^{(i)} \rightarrow L_1^{(i)}, S_0 \rightarrow S_1 = 1$$

with observed in data increase of masses:  $m_n \rightarrow m_n + 2m_\pi$  .

At second transition

$$L_1^{(i)} \rightarrow L_2^{(i)}, S_1 \rightarrow S_2$$

reactions must happen

$$(n + 2\pi) + (p + 2\pi) \rightarrow (n + K^+) + (p + K^-).$$

with transformation of

2 pairs  $\pi^+\pi^-$  into  $K^+K^-$  .

# PROBABILITIES OF PEAKS

Mean cross-section of peaks

$$\langle \sigma_{nK} \rangle = 63 \text{ (9) mkb} / 7 = 9 \text{ (1.5) mkb}$$

is  $\sim 1/2$  of  $\langle \sigma_{\pi\pi} \rangle = 16 \text{ (2) mkb}$  :

it agrees with ratio of statistical

weighs 
$$\frac{g(S=0)g(S_z=-1)}{g(S_z=-1)} = \frac{1/4 \times 1/2}{1/4} = \frac{1}{2}$$

Reduced probability of peak

$$W^{(i)} = \langle \sigma^{(i)} \rangle / (\sigma_{(\text{inel.})} g(S)g(L^{(i)})) \sim \mathbf{0.06},$$

with  $\sigma_{(\text{inel.})} = 30 \text{ mb}$ ,  $(L^{(i)}) = 1/L^{(\text{max})} \sim 0.04$ .

# NONEQUILIBRIUM NN INTERACTION

With momentum  $L \sim 20$  and moment of inertia  $I_L = 2mR_0^2$  rotation energy

$$E = \frac{h^2 L(L+1)}{2I_L} \sim 20 \text{ GeV} \gg E_0$$

For conservation energy during

$$\Delta t \sim \frac{h}{\Delta E} \sim 0.01 \text{ fm}/c \sim 3 \cdot 10^{-26} \text{ sec}$$

potential energy of nucleons interaction

must be decreased according to  $\Delta E$  :

$$\Delta U^{(neq)} = - \Delta E .$$

Here system (np) with momentum  $L$  at once is created in such states with large probability  $\sim 100\%$  of suitable (np) events.

Short existence  $\Delta t$  of rotary system can explain small  $\Delta r$  of formation  $\pi^+ \pi^-$  .

# Interaction is nonequilibrium even in stationary states

In stationary state with energy  $E$  of particle in equilibrium potential  $U(r)$  nonequilibrium interaction  $U(r)^{(neq)}$  arises always when

$$E - U(r) < 0:$$

$E - U(r)^{(neq)} = T(r) > 0$  keeps positive sign of kinetic energy.

This is shown in quantum theory (without wave properties of particles) “**Statistical physics of undistinguishable events (formalism and examples of use)**”

preprint **PNPI-2005 2628**, 2005, Gatchina

# WHAT'S NECESSARY TO EXPLAIN

It is desirable to explain:

1) how energy  $\Delta E(J) = M_{\pi\pi}c^2$  reveals of accident kinetic energies of nucleons ;

2) what properties of nucleons define empirical parameter

$$V(\text{emp}) = (V_{\pi\pi} + V_{nk})/2 = 27.7 (0.5) \text{ MeV};$$

3) what mechanism can give ~100% probability of  $\Delta L=2$  transition at short time  $\Delta t \sim 10^{-26}$  sec;

4) why preliminary transition to  $S=1$  in  $npK^+K^-$  events happens with ~100% probability and with formation of 2 pairs  $\pi\pi$  .

# TRANSITION TO SOLID-BODY ROTATION

Formed in moment of collision  
orbital movement with momentum

$$L, \quad \Omega = Lh/I_L, \quad E = \Omega^2 I_L / 2$$

is unstable and at once reduces to  
solid-body rotation with momentum

$$\text{of inertia } I_{s-b} = I_L + I_S, \quad \Omega_{s-b} = Lh/(I_L + I_S)$$

$$\text{and } E_{s-b} = \Omega_{s-b}^2 (I_L + I_S) / 2.$$

$$E - E_{s-b} = E I_S / (I_L + I_S) = U_{s-b}$$

is potential energy of interaction of

$$hL_{s-b} = \Omega_{s-b} I_L \quad \text{and} \quad hS_{s-b} = \Omega_{s-b} I_S$$

at the time  $\Delta t$  of solid-body rotation

$$\text{with momentum } L = L_{s-b} + S_{s-b}.$$

# DIFFERENCE OF POTENTIAL ENERGIES IN STATES WITH $\Delta L=2$

At transition  $L_1^{(i)} \rightarrow L_2^{(i)}$ ,  $S_1 \rightarrow S_2$

$$\begin{aligned}\Delta U^{(h-b)}_{12} &= U_{h-b}(L_1) - U_{h-b}(L_2) = \\ &= \frac{(2L_1-1) h^2 I_S}{I_L (I_L + I_S)} = (2L_1-1) V_{SL}.\end{aligned}$$

With parameter  $I_S = mR^2$

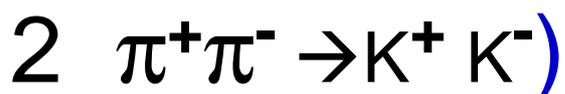
$$V_{SL} = \frac{h^2 I_S}{I_L (I_L + I_S)} = 27.6 \text{ MeV}$$

coincides with

$$V^{(\text{emp})} = (V_{\pi\pi} + V_{nk})/2 = 27.7 (0.5) \text{ MeV}$$

# ENERGY $\Delta U^{(s-b)}_{12}$ LIBERATION

At decay of rotary system excess energy  $\Delta U^{(s-b)}_{12}$  keeps as inner energy of nucleons, which end interaction, and can be emitted in this moment in form of  $\pi^+\pi^-$  energy (or in form energy of reaction



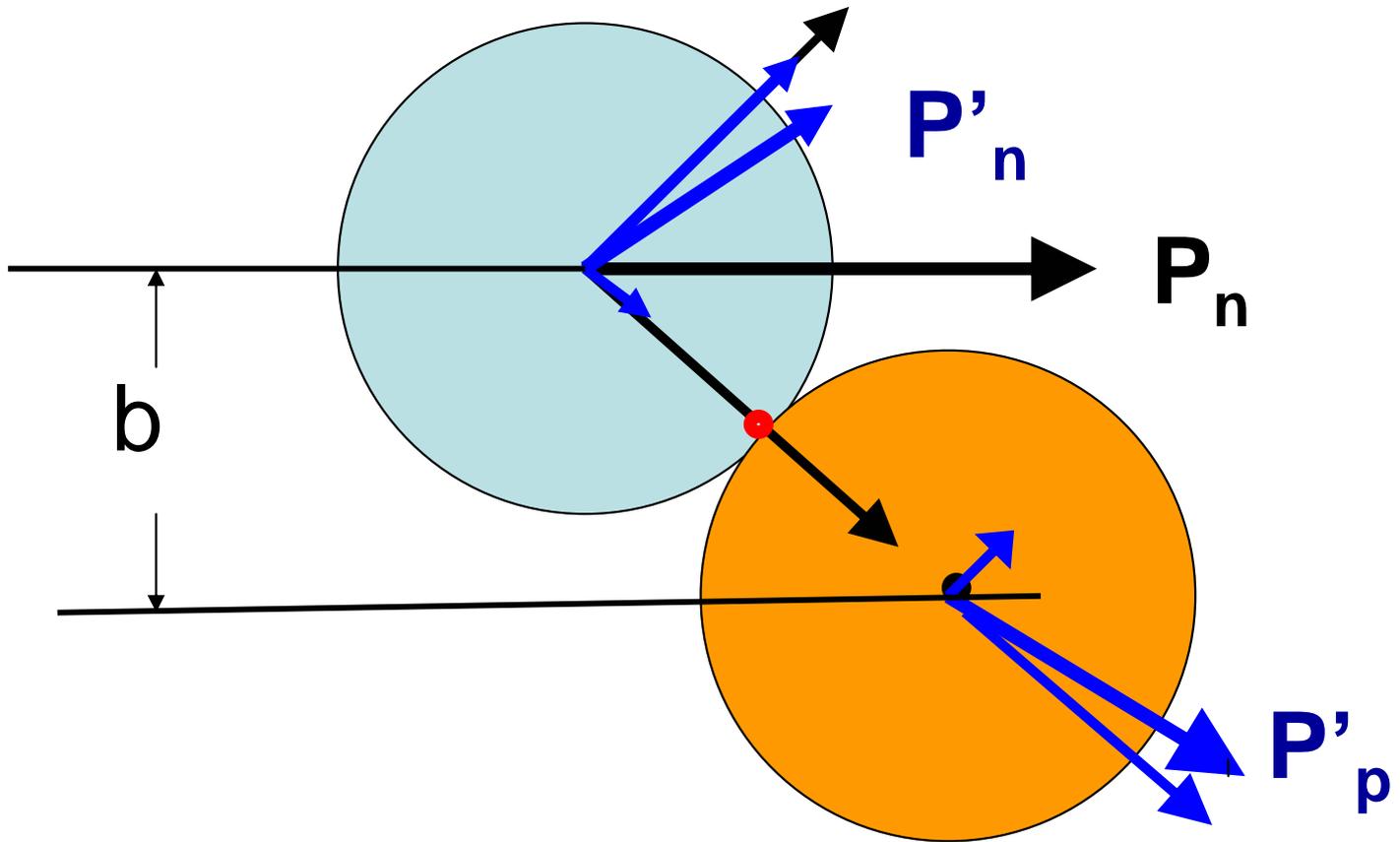
# “FORSED” STRONG DELAY OF NONEQUILIBRIUM ROTATION

Law  $\Delta \mathbf{B}=0$  means: “black balls” must safe for description of nucleons.

Together with exchange of radial impulses of nucleons exchange of tangential impulses is going on, so orbital moment must decrease.

Therefore only reducing rotation may be exist  $\Delta t$ , where contribution of spin  $S=1$  in moment  $J=L+S$  replaces a decrease of momentum  $L=\Omega(I_L+I_S)$  of rotation with decreasing of  $\Omega$ .

# DECREASE OF ORBITAL MOMENTUM



At nonequilibrium rotation:

$L = P_n b$  as in lab.sys.: transfer of radial impulse --- as hard balls.

$L_1 = 2R_0 P_n^{(\text{tan.})}$  can not be kept

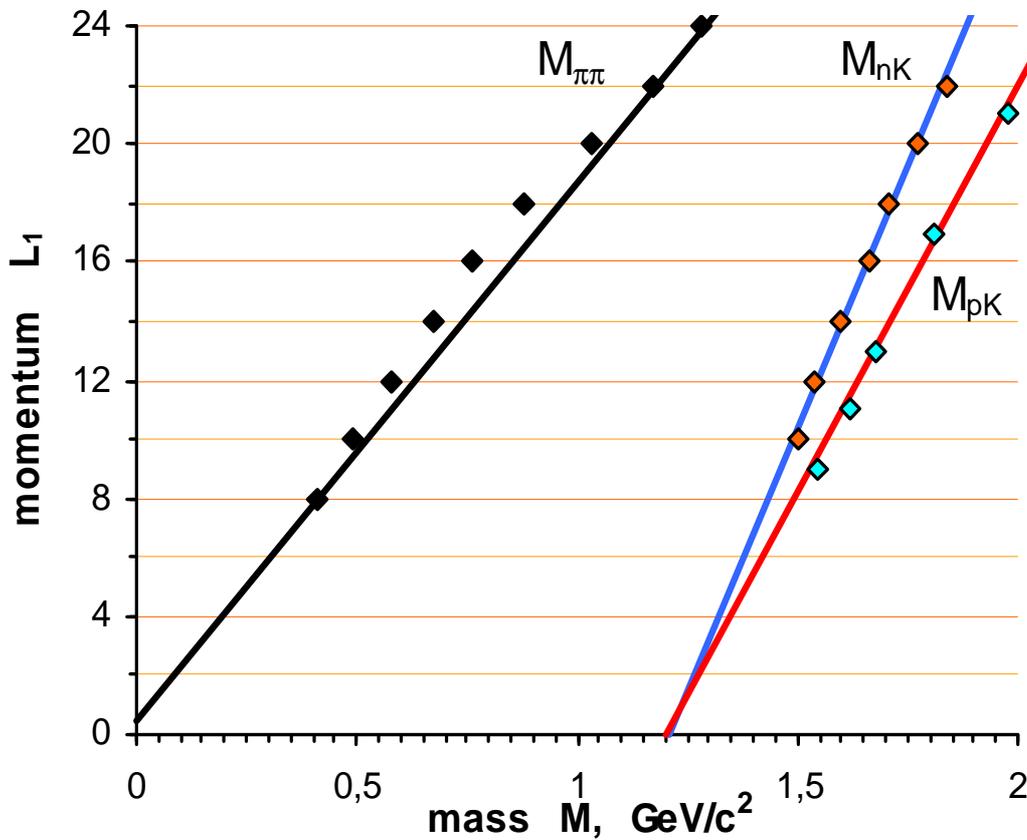
$L_2 = 2R_0 (P'_n^{(\text{tan})} - P'_p^{(\text{tan})}) = L_1 - 2$

# CAUSAL EVENTS IN NONEQUILIBRIUM NN SYSTEM

**Law  $\Delta B=0$**  for nucleon (impenetrable distributions --- “black balls”),  
**law of conservation  $J + L + S$**  and  
**law of conservation energy** lead to  
formation of short delay rotation  
state with **causal evolution and**  
**replacing of  $\Delta L=2$  by  $\Delta S=2$**  .

For short  $\Delta t$  probabilities of any  
accidental events is negligible small.  
Energy  $\Delta U^{(h-b)}_{12}$  is additional inner  
energy of nucleons in moment of  
disintegration of two-nucleon system

# EXP. SPECTRA AND ITS DESCRIPTION



(black line)  $M_{\pi\pi}^{(i)} = 2V (L_1^{(i)} - 1/2)/c^2$  ,

(blue line)  $M_{nK}^{(k)} = V (L_1^{(k)} - 1/2)/c^2 + m + 2m_\pi$

(red line)  $M_{pK}^{(i)} = V' (L_1^{(i)} - 1/2)/c^2 + m + 2m_\pi$

$$V = \frac{h^2}{6mR_0^2}, \quad V' = \frac{h^2}{2mR_0^2} \frac{(m + 4m_\pi)}{(3m + 4m_\pi)}$$

$R_0 = 0.50 \text{ fm}$ ,  $m = m$  is nucleon mass

**DATA OF  $p + C_3H_8 \rightarrow pK^0_S + X$   
at  $P_p = 10.0 \text{ GeV}/c$**

Spectrum  $M_{pK^0}$  may be described  
the same ratio as  $M_{nK^+}$  spectrum;  
increasing of value  $V'$  can be  
explained by contribution of 4  $\pi^-$   
mesons in inner momentum inertia  $I_S$

Value  $\sigma^{(i)} \sim 90 \text{ mkb}$  for one peak  
corresponds the same reduced  
probability  $W^{(i)} \sim 0.06$

at effective number  $N_p = 20$  in  $C_3H_8$ ,  
or weight  $2/3$  of nuclear protons

# COMMENTARY

These data can not be described by present quantum theory, which made for calculation of probabilities of possible accident events from them distributions.

Crude model of solid-body rotation of two hard balls system with laws of conservation energy and momentum well describes data and explains main properties of unusual causal events with whole nucleons.

# SOME PREDICTIONS OF MODEL

Momentum  $S'_p$  of inner rotation of protons with energy of this rotation must be transmitted to neutrons at decay of rotary system since can not be kept at free proton formation.

Free movement of finite proton gives zero contribution in momentum  $\Sigma L_f$  relative to point O ---center of proton.

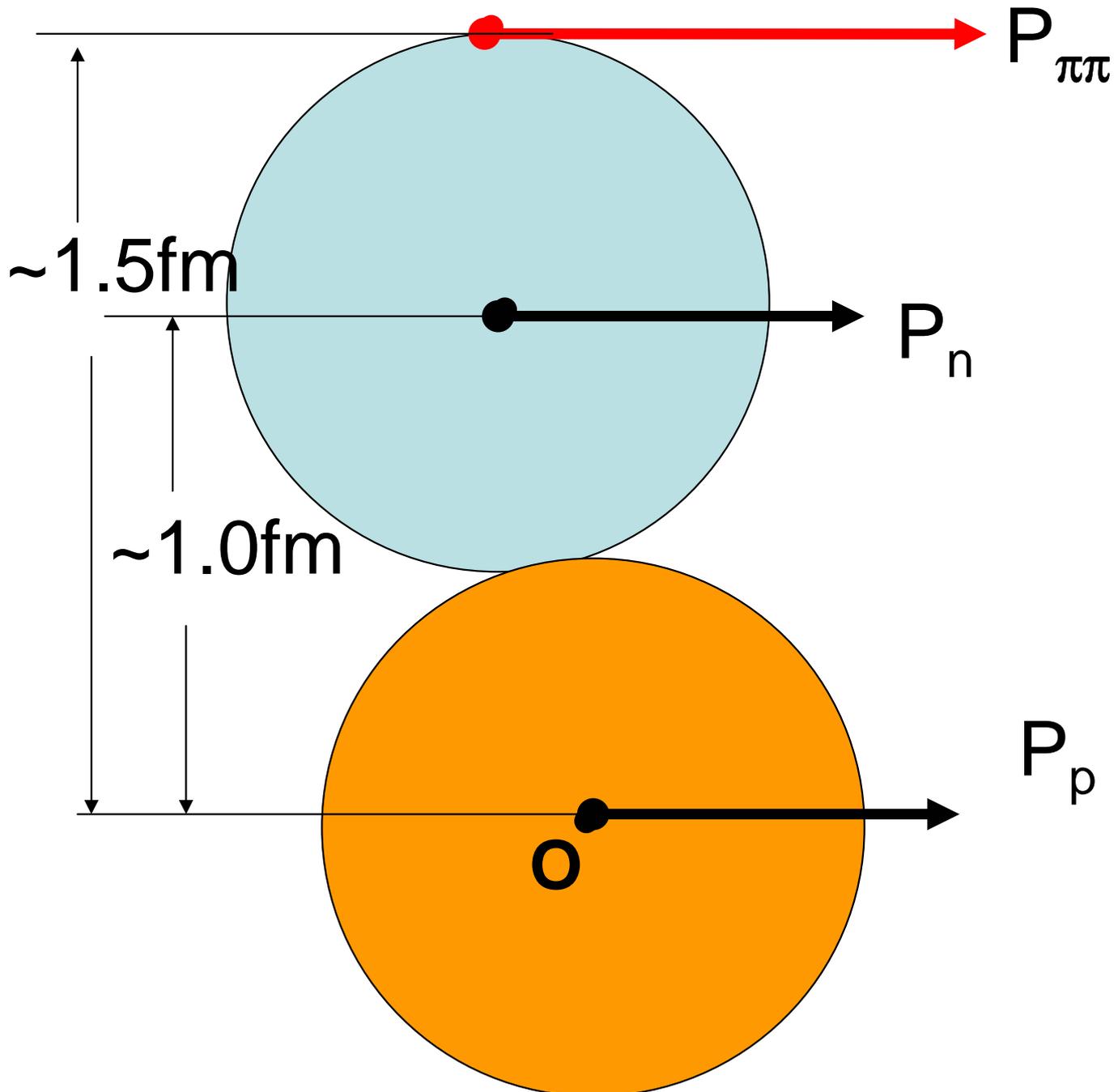
Therefore pairs  $\pi^+\pi^-$  or  $K^+K^-$  mesons form at decay of excited state of neutron.

In the end of distribution  $M_{\pi\pi} \sim 1.4 \text{ GeV}/c^2$  and kinetic energy of relative movement of n,p and pair mesons in s.c.m. is small:  
 $E' - (m_n - m_p - M_{\pi\pi})c^2 < 100 \text{ MeV}$ .

Value  $L=26$  and condition  $\Sigma L_f = L-2$  lead to next conclusions:

# SOME PREDICTIONS OF MODEL

Point of formation  $\pi^+\pi^-$



$$P_n = 1.5 \text{ GeV}/c, \quad L_n = 7.5$$

$$P_{\pi\pi} = 2.2 \text{ GeV}/c, \quad L_{\pi\pi} = 16.5$$

$$L_n + L_{\pi\pi} = 24 = 26 - 2$$

## PREDICTIONS OF MODEL FOR $npK^+K^-$

The end of distribution  $M_{nK} \sim 1.98 \text{ GeV}/c^2$  corresponds to small kinetic energy of relative movement of  $nK^+$  and  $p, K^-$  mesons:  $E' - (m_n - m_p - M_{\pi\pi})c^2 \sim 0$ .

Impulse  $P'_n = P'_{K^+} = 0.66 \text{ GeV}/c$  of  $n$  and  $K^+$  in s.c.m. and their energies  $E'_n = 1.15 \text{ GeV}$ ,  $E'_{K^+} = 0.83 \text{ GeV}$ , moved with s.c.m., give impulses  $P_n = 1.75 \text{ GeV}/c$ ,  $P_{K^+} = 1.26 \text{ GeV}/c$ . Contribution of proton  $L_p = 0$  here too.

With impulse of  $K^-$   $P_{K^-} = 0.76 \text{ GeV}/c$  summary impulse of mesons  $2.02 \text{ GeV}/c$  can create moment  $L_{K^+K^-} = 15$ , if  $K^+K^-$  form as pairs  $\pi^+\pi^-$ .

Together with moment  $L_n = 9$  of neutron it gives model value  $\Sigma L_f = 24 = L^{(\text{max})} - 2$ .

Some prediction for the end of distribution

$M_{pK_0}$  of reaction  $ppK^0_S K^0_L$  at  $P_p=10\text{GeV}/c$

With  $R_0=0.50\text{ fm}$   $L^{(\text{max})}=2R_0P_p/h = 50$  and

$$M_{pK_0}^{(\text{max})}=M_0+V'_{SL}L^{(\text{max})}/c^2=3.06\text{ GeV}/c^2.$$

It confirm with maximum contribution of

$pK^0_S$  in  $E'=(E_0^2-P_0^2c^2)^{1/2}=4.54\text{ GeV}$  :

$$M_{pK_0}^{(\text{max})}=E'/c^2-m_p-m_K=3.10\text{ GeV}/c^2.$$

If impulses  $P'_p=P'_{K_0}=1.36\text{GeV}/c$  relative movement  $p$  and  $K^0_S$  in s.c.m. are trans-

verse, energies  $E'_p$  and  $E'_{K_0}$  in s.c.m. give longitudinal impulses  $P_p=3.63\text{ GeV}/c$  and

$P_{K_0}=3.19\text{ GeV}/c$  in lab. system.

With contribution of proton  $L_p=2R_0P_p/h=18$  summary moment of two K-mesons

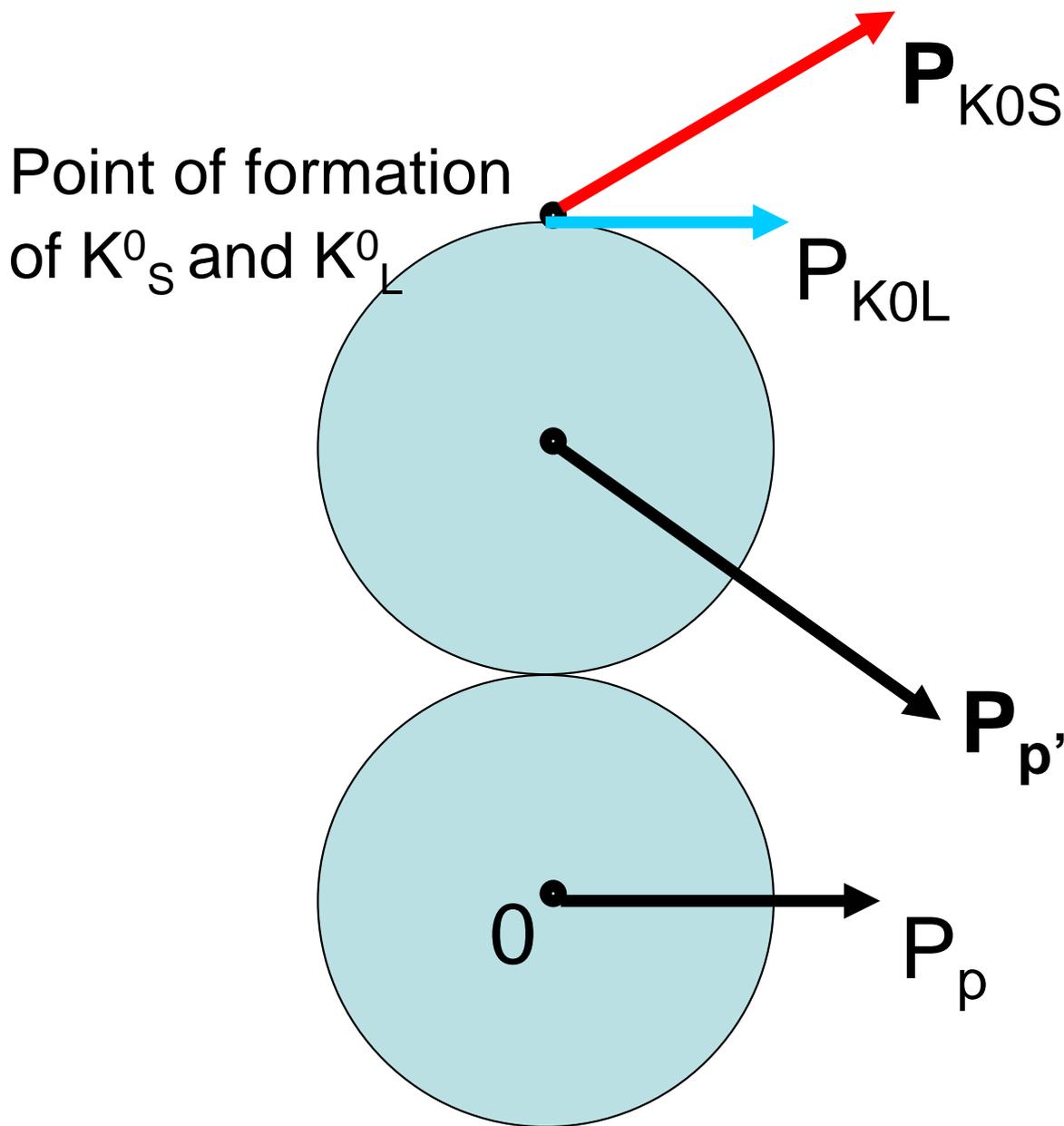
$$L_{KK}=3R_0(3.19+1.10)\text{GeV}/(hc)=32\text{ gives}$$

$$L_p+L_{KK}=50=L^{(\text{max})},\text{ if pair of K-mesons}$$

forms in the same point as  $\pi^+\pi^-$  mesons.

It confirms that model is consistent.

Probable positions of final particles in events  $ppK^0_S K^0_L$  with maximum  $L_f=50$



Model of hard balls and keeping of moment  $L$  (relative point 0 in lab. system)  $L_f=L_{p'}+L_{K^0_S}+L_{K^0_L}$  (here for  $L_f^{(max)}=50$ ) define spatial positions of formation of final particles and them impulses

# COMMON CONCLUSIONS

Inelastic NN interaction is like as collision of black balls with radius  $R_0 = 0.50$  fm.

With large probability (possibly, 100%) short rotation states form in laboratory system with quantized momentum  $J=L+S$ .

These facts are beyond question and must be taken into account.

Data prove absence of “invariance”. Short strong interaction can explain origin all events, which are treated as observation of small “partons”.

# SEPARATELY OF PAIRS $\pi\pi$ ON SURFISE OF NUCLEONS

These data give unique chance: here all characteristics of two-nucleonic system are defined and it is known, that formation of pairs  $\pi\pi$  with 100 % probability happens in defined point in result of nonequilibrium interaction with known value.

But without right electrodynamics it is impossible to use this to gain a better insight of observed events.

**THANK YOU**

**FOR ATTENTION !**

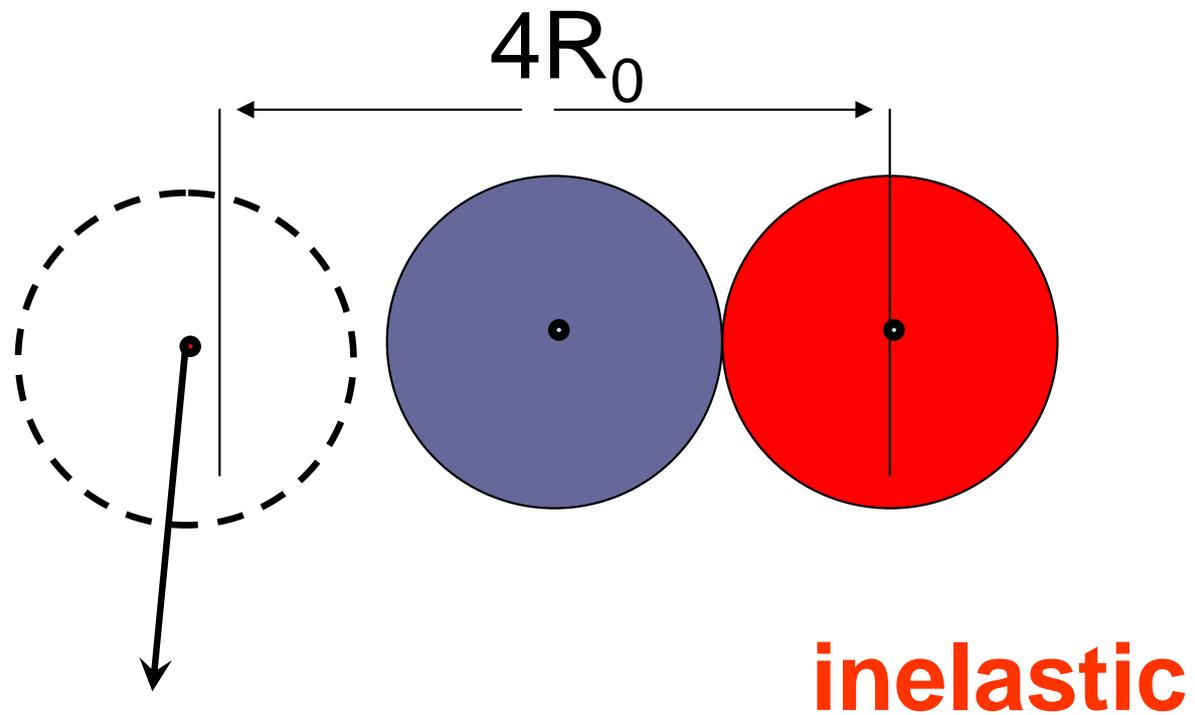
# CONFIRMATION OF “BLACK BALL” MODEL BY NN SCATTERING DATA

Independent of energy in wide interval 10-1000 GeV cross-sections of elastic and inelastic NN scattering  $\sigma_{el} \cong 8$  mb and  $\sigma_{inel} \cong 32$  mb have geometrical sense  $\sigma_{inel} \cong \pi(R_0)^2$ ,  $\sigma_{el} \cong \pi(2R_0)^2$  and correspond to collisions of black balls with radius  $R_0 \cong 0.50$  fm.

It explains ratio  $\sigma_{inel} / \sigma_{el} \cong 4$ , which is fulfilled at energy  $E > 10^6$  GeV.

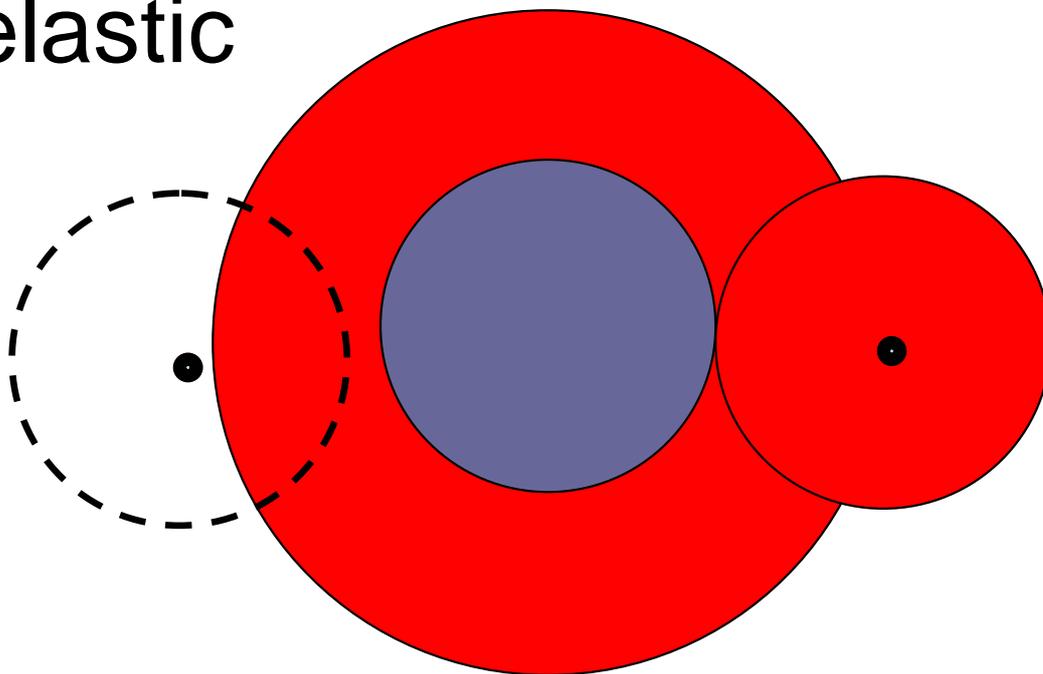
(Generally accepted theory of scattering of plane wave  $\psi = e^{ikz}$  is useless for description of black balls scattering.)

# GEOMETRICAL TREATMENT OF NN SCATTERING CROSS-SECTIONS



elastic

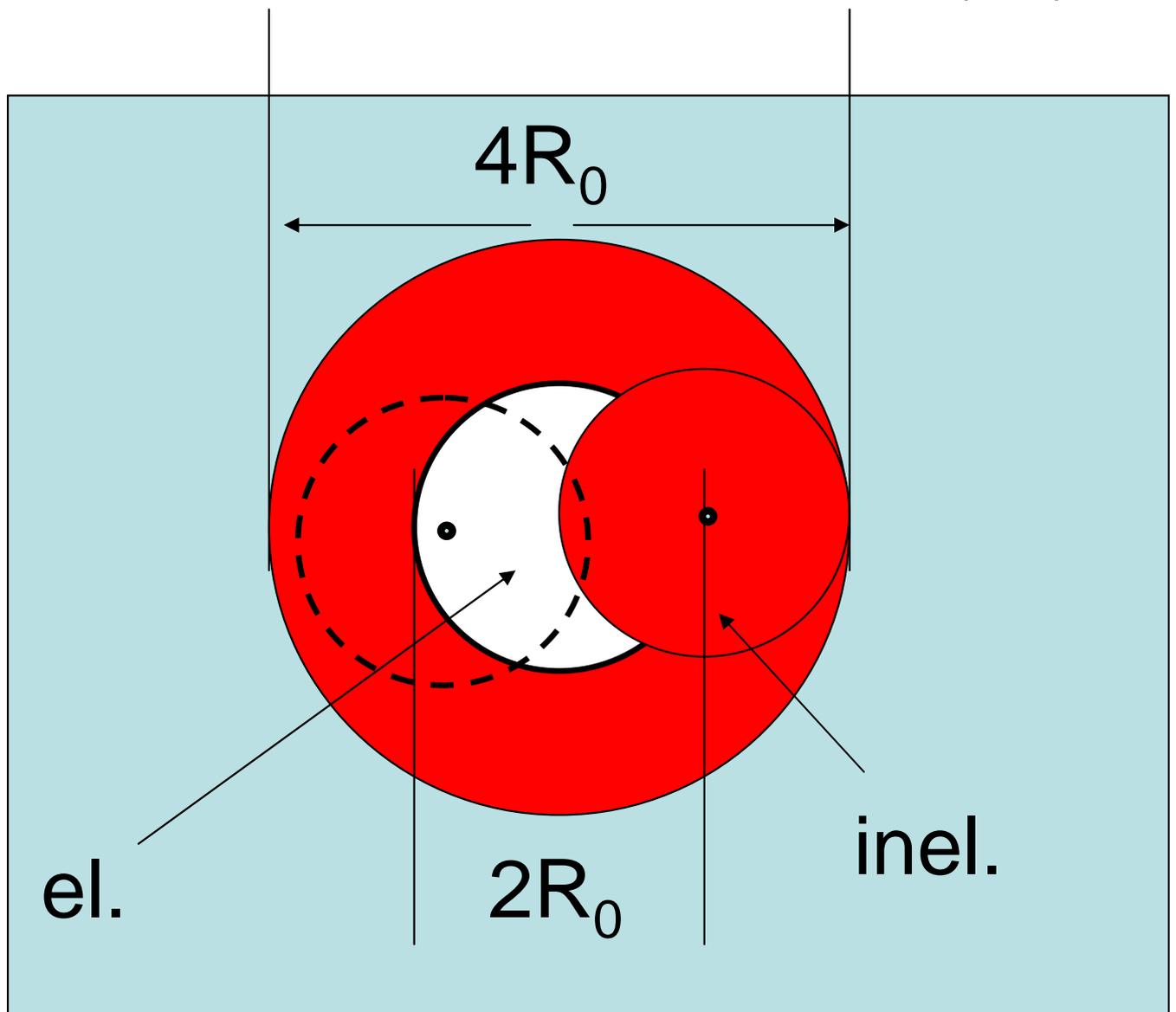
inelastic



$$\sigma_{\text{(inel.)}} = \pi (2R_0)^2$$

# RATIO $\sigma_{(inel.)} / \sigma_{(el.)} = 4$ FOR BLACK BALLS SCATTERING

hole  $4R_0$  ,  $S = \sigma_{(inel.)}$



$$\sigma_{(el.)} = \pi R_0^2 = 1/4 \sigma_{(inel.)}$$

# RELATIVITY IS INCOMPATIBLE WITH MAXWELL EQUATIONS

$\text{div } \mathbf{E} = \rho$  and relativistic ratio  $c\mathbf{H} = [\mathbf{v}\mathbf{E}]$   
for magnetic field of moving point  
charge  $e$  with field  $\mathbf{E}(\mathbf{r}-\mathbf{v}t)$  gives

$$\text{crot } \mathbf{H} = \text{div } \mathbf{E} - (\mathbf{v} \text{grad}) \mathbf{E} = \mathbf{v}e = \mathbf{j},$$

where only conduction current  $\mathbf{j}$   
is, instead of Maxwell equation

$$\text{crot } \mathbf{H} = \mathbf{j} + \partial \mathbf{E} / \partial t$$

with displacement current  $\partial \mathbf{E} / \partial t$ .

Right equation  $\mathbf{P} = (\mathbf{P}^2 + m^2 c^2)^{1/2} \mathbf{v} / c$  for  
impulse of moving object is empirical  
law, independent of any theories.

# RIGHT SOLUTION OF MAXWELL EQUATIONS

$\partial \mathbf{E}(\mathbf{r}-\mathbf{v}t)/\partial t = \partial \mathbf{E}(\mathbf{r}-\mathbf{v}t_0)/\partial r_v$  with  $dr_v = -v dt$ .

$d\mathbf{r} = d_r r_v - d_{(\text{perp})} r_v$  and  $\partial \mathbf{r} / \partial r_v = 1$ ;

**( $t_0$  is essential !)** The placing of next ratio

$\partial \mathbf{E} / \partial t = \partial \mathbf{E}(\mathbf{r}-\mathbf{v}t_0) / \partial \mathbf{r} (\partial \mathbf{r} / \partial r_v) \partial r_v / \partial t =$

$= -\mathbf{v} \partial \mathbf{E}(\mathbf{r}-\mathbf{v}t_0) / \partial \mathbf{r} = -\mathbf{v} \text{div } \mathbf{E} = -\mathbf{v}e = -\mathbf{j}$

in Maxwell equations

$c \text{rot } \mathbf{E} = -\partial \mathbf{H} / \partial t, \quad \text{div } \mathbf{H} = 0, \quad (1,2)$

$c \text{rot } \mathbf{H} = \partial \mathbf{E} / \partial t + \mathbf{j}, \quad \text{div } \mathbf{E} = \rho. \quad (3,4)$

gives equation  $c \text{rot } \mathbf{H} = \mathbf{j} - \mathbf{j} = 0$  for free

point charge with velocity  $\mathbf{v}$  and solution

$\mathbf{H} = 0$ , which is incompatible with relativity.